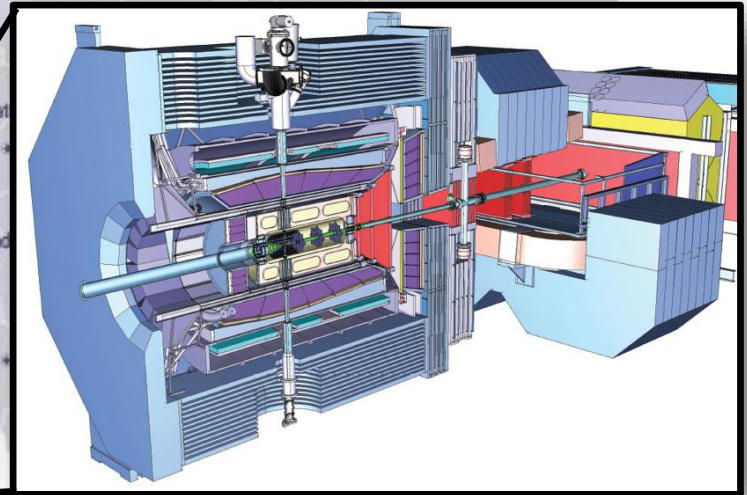
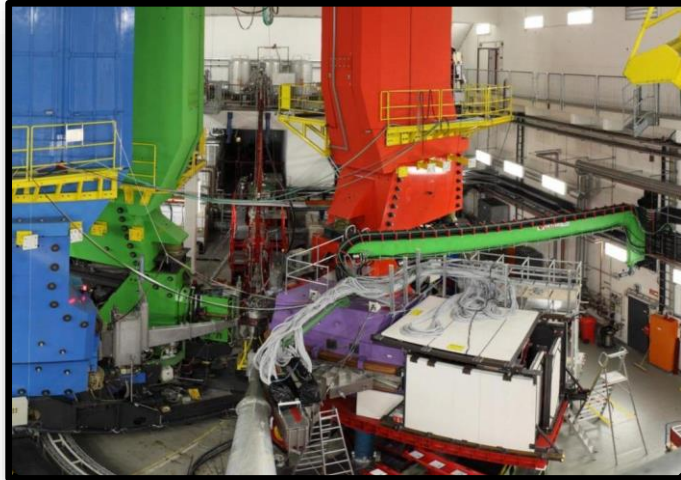


Solving puzzles in hypernuclear physics at MAMI and PANDA

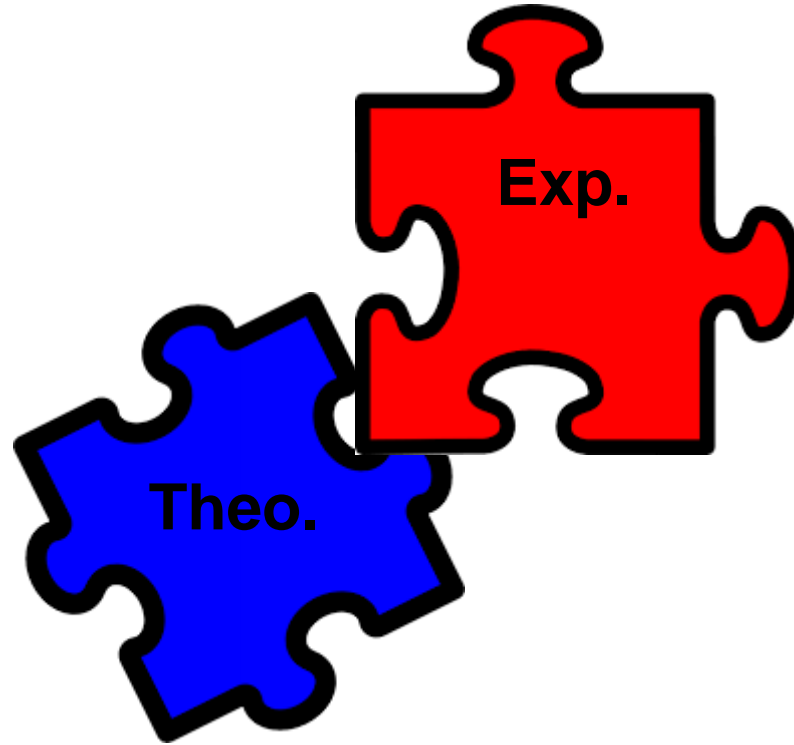


Patrick Achenbach

U Mainz

October 2016

Light Λ hypernuclei and the puzzle of charge symmetry breaking



The world as we know it relies on charge symmetry breaking

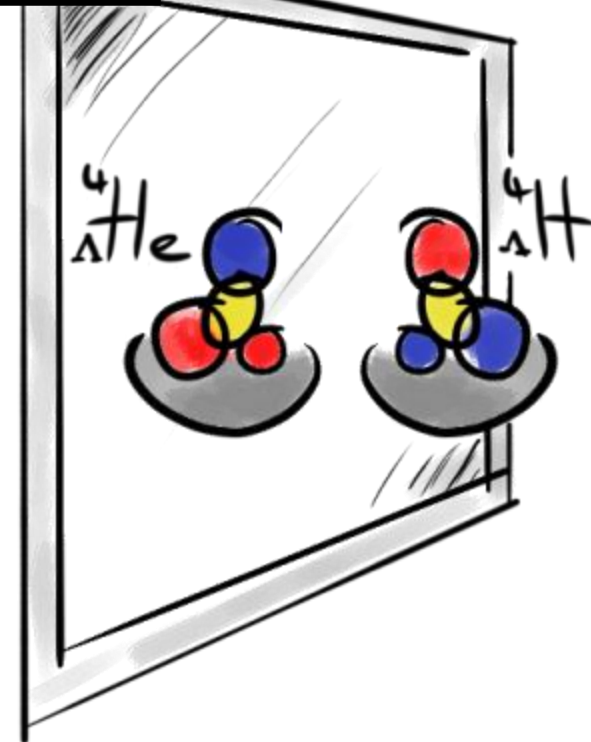
Charge independence: strong force
independent of nucleon isospin state

$$(F_{p-p} = F_{n-n} = F_{p-n})$$

Charge symmetry: strong force
independent of nucleon isospin exchange

$$(F_{p-p} = F_{n-n})$$

Coulomb interactions are rather ineffective in
breaking the isospin symmetries in nuclei



If charge symmetry would be satisfied in nuclear two-body force...

... mirror hypernuclei would have (nearly) identical mass

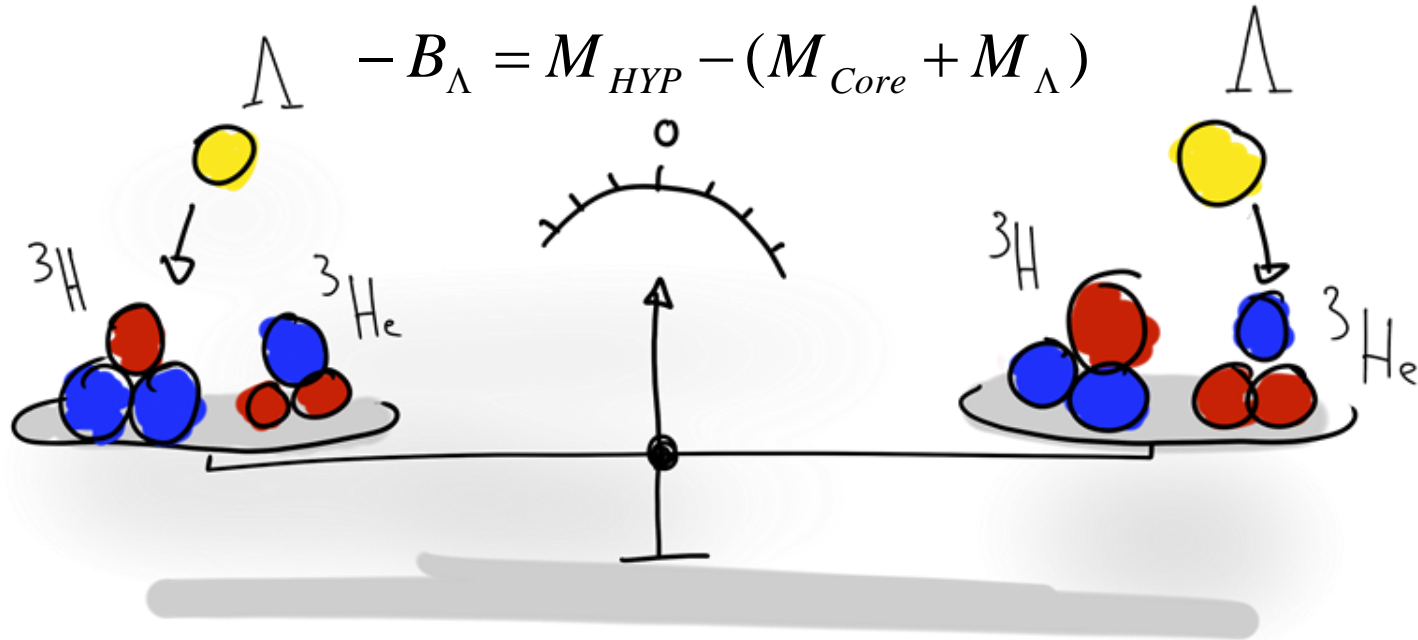
... protons would be heavier than neutrons because of electrostatic repulsion

... no free protons would have survived the primordial nucleosynthesis

... the Sun and the stars would have no slow-burning fuel

Charge symmetry in light hypernuclei

Opportunity to study strong force symmetries with Λ as neutral probe



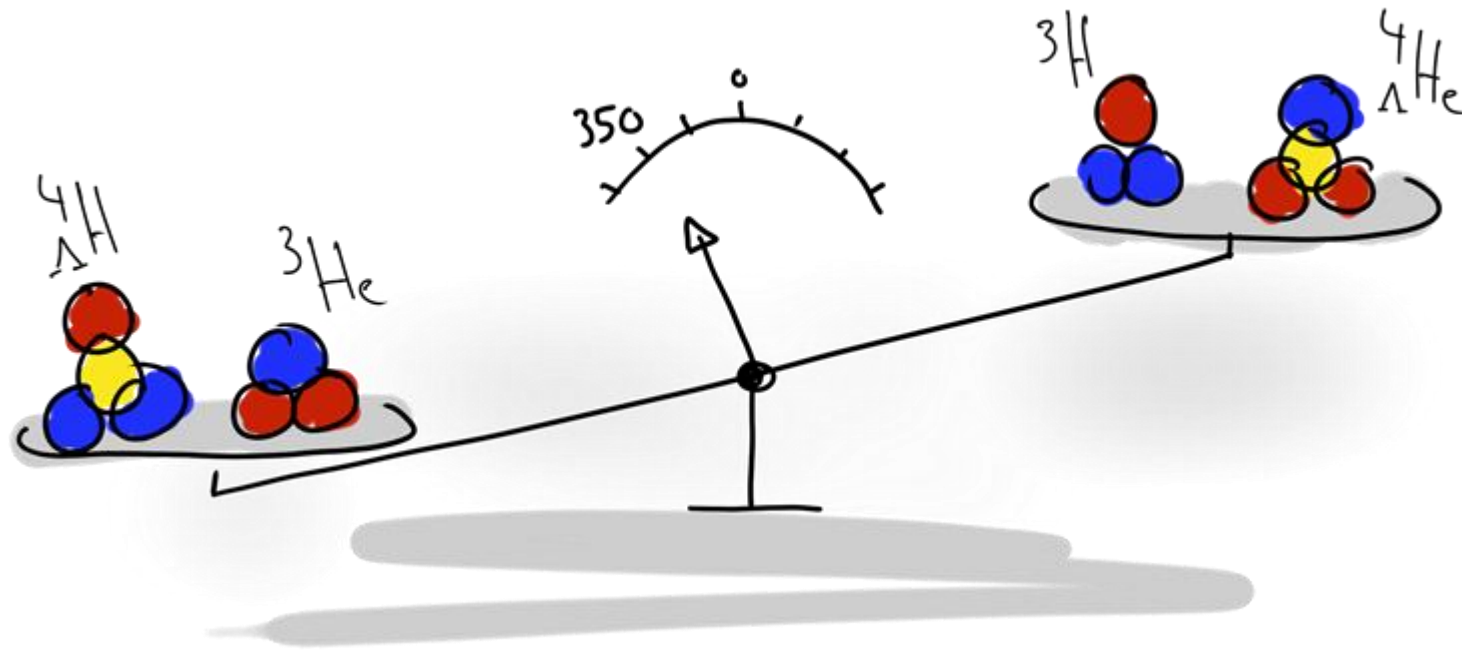
Λ hyperon has no isospin and no charge \rightarrow
 Λ binding in mirror hypernuclei directly tests charge symmetry

$$F_{\Lambda\text{-p}} = F_{\Lambda\text{-n}} \rightarrow B_{\Lambda}({}_{\Lambda}^AZ) = B_{\Lambda}({}_{\Lambda}^{AZ+1})$$

Large charge symmetry breaking in $A = 4$

${}^4_{\Lambda}\text{H} - {}^4_{\Lambda}\text{He}$ binding energy difference exceptionally large > 300 keV

[M. Juric et al. NP B52 (1973)]



- Charge symmetry breaking 5 times larger than in ${}^3\text{H} - {}^3\text{He}$ system
- Λp interaction stronger than Λn interaction

Observations not consistently reproduced by theory before 2015

[A. Nogga, NP A 914,140 (2013)]

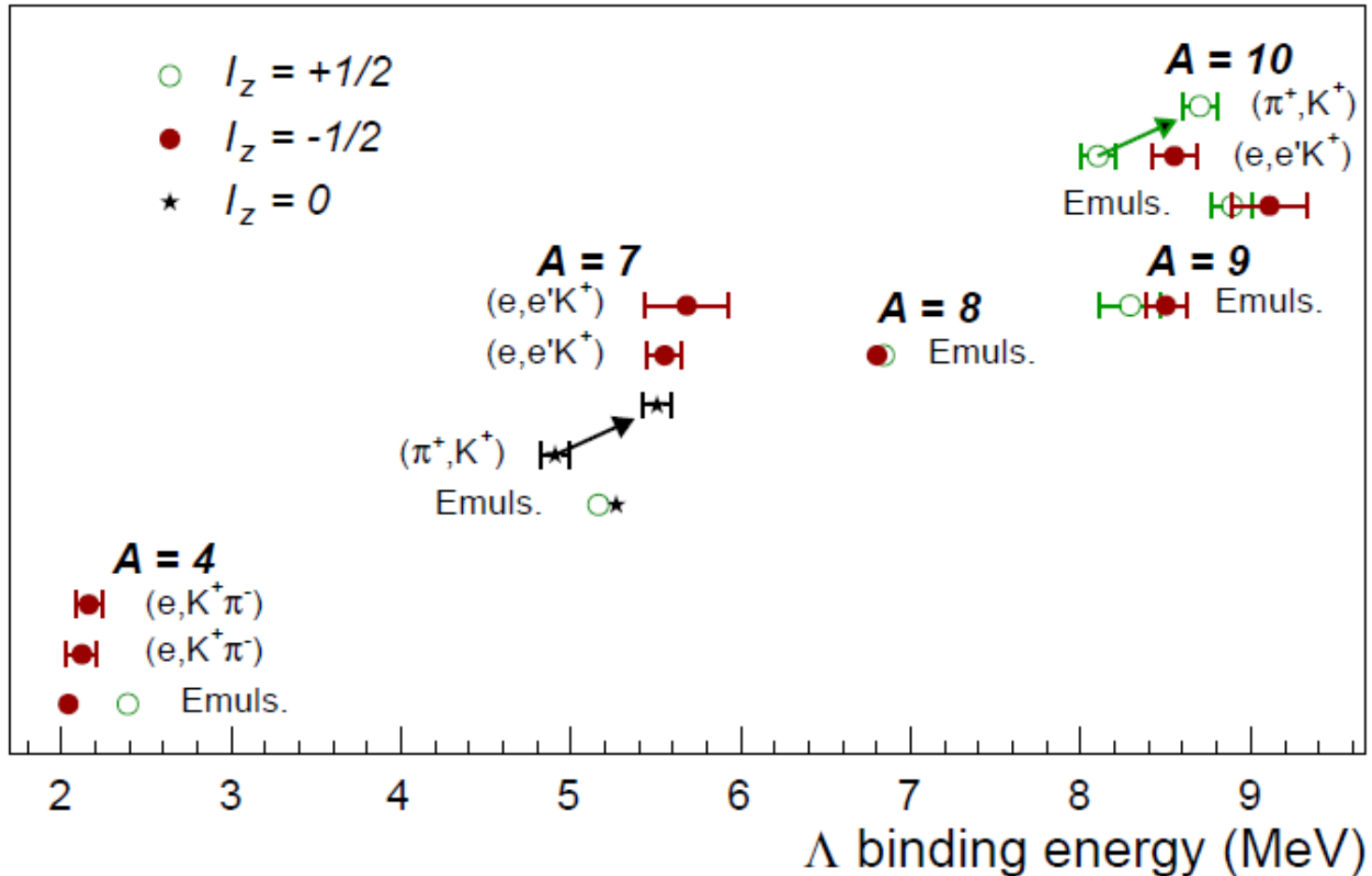
The $A = 4$ CSB puzzle

Calculation	Interaction	$B_{\Lambda}(^4_{\Lambda}\text{H}_{\text{gs}})$	$B_{\Lambda}(^4_{\Lambda}\text{He}_{\text{gs}})$	ΔB_{Λ} ($^4_{\Lambda}\text{He}-^4_{\Lambda}\text{H}$)
A. Nogga, H. Kamada and W. Gloeckle, PRL 88, 172501 (2002)	SC97e	1.47	1.54	0.07
	SC89	2.14	1.80	0.34
H. Nemura, Y. Akaishi and Y. Suzuki, PRL 89, 142504 (2002)	SC97d	1.67	1.62	-0.05
	SC97e	2.06	2.02	-0.04
	SC97f	2.16	2.11	-0.05
	SC89	2.55	2.47	-0.08
E. Hiyama, M. Kamimura, T. Motoba, T. Yamada and Y. Yama PRC 65, 011301 (R) (2001)	AV8	2.33	2.28	-0.05
world data average		2.04 ± 0.04	2.39 ± 0.03	0.35 ± 0.06

- calculations since decades fail to explain large ΔB_{Λ}
- coupled channel calculation using SC89 fails to bind excited state

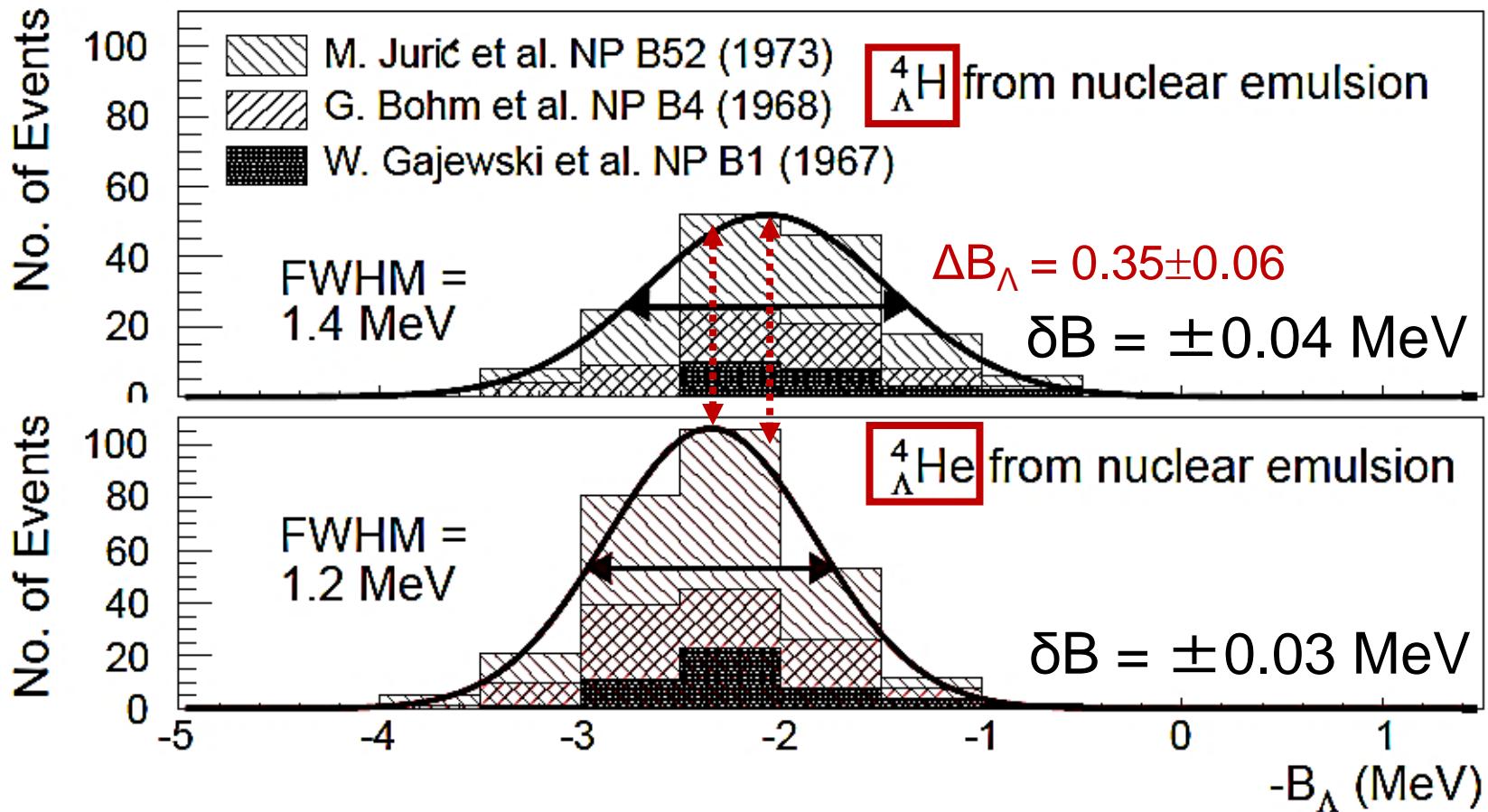
World data on $A \leq 10$ systems

Hypernuclear isomultiplets



- clearest signature of charge symmetry breaking in $A = 4$ system
- weak indications of charge symmetry breaking in $A \neq 4$ systems

Emulsion results on ${}^4_{\Lambda}\text{H}$ and ${}^4_{\Lambda}\text{He}$

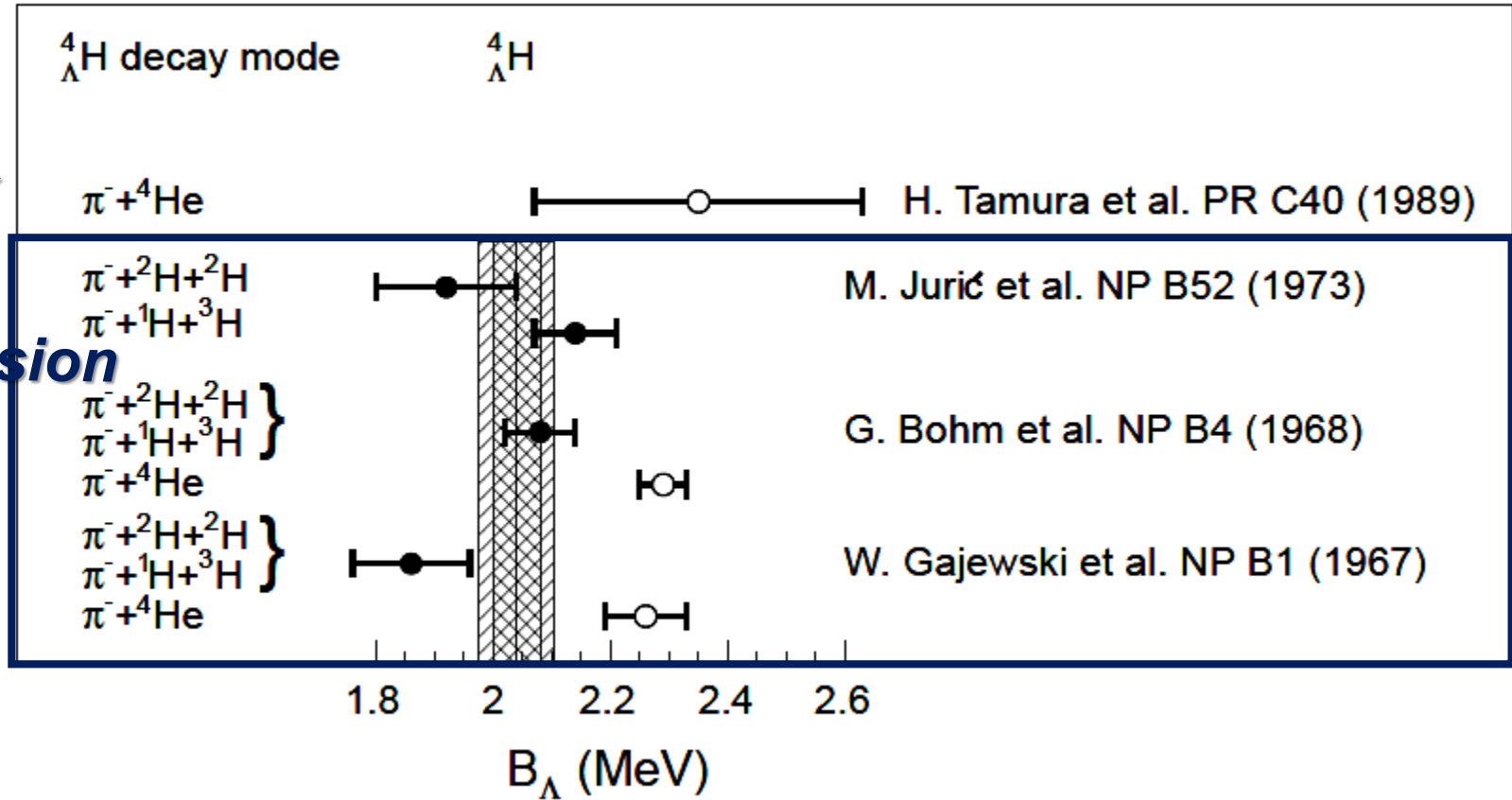


- only three-body decay modes used for hyperhydrogen
- 155 events for hyperhydrogen, 279 events for hyperhelium

World data on ${}_{\Lambda}^4\text{H}$

KEK

emulsion



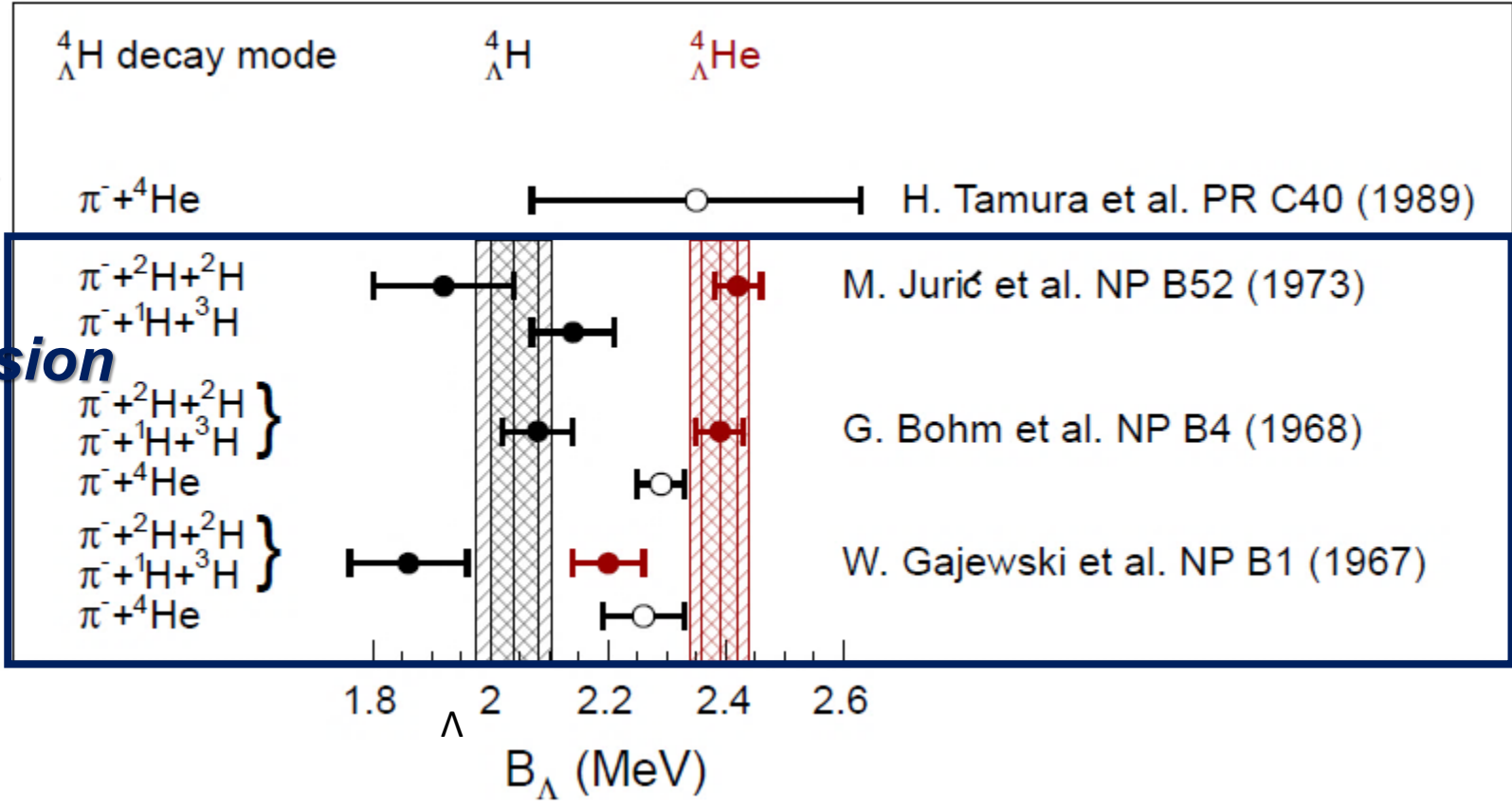
$$\left. \begin{array}{l}
 {}_{\Lambda}^4\text{H} \xrightarrow{\text{decay}} \pi^- + {}^1\text{H} + {}^3\text{H}: B = 2.14 \pm 0.07 \text{ MeV} \\
 {}_{\Lambda}^4\text{H} \xrightarrow{\text{decay}} \pi^- + {}^2\text{H} + {}^2\text{H}: B = 1.92 \pm 0.12 \text{ MeV}
 \end{array} \right\} 0.22 \text{ MeV difference}$$

Total: $B = 2.08 \pm 0.06 \text{ MeV}$ [M. Juric et al. NP B52 (1973)]

World data on $A = 4$ system

KEK

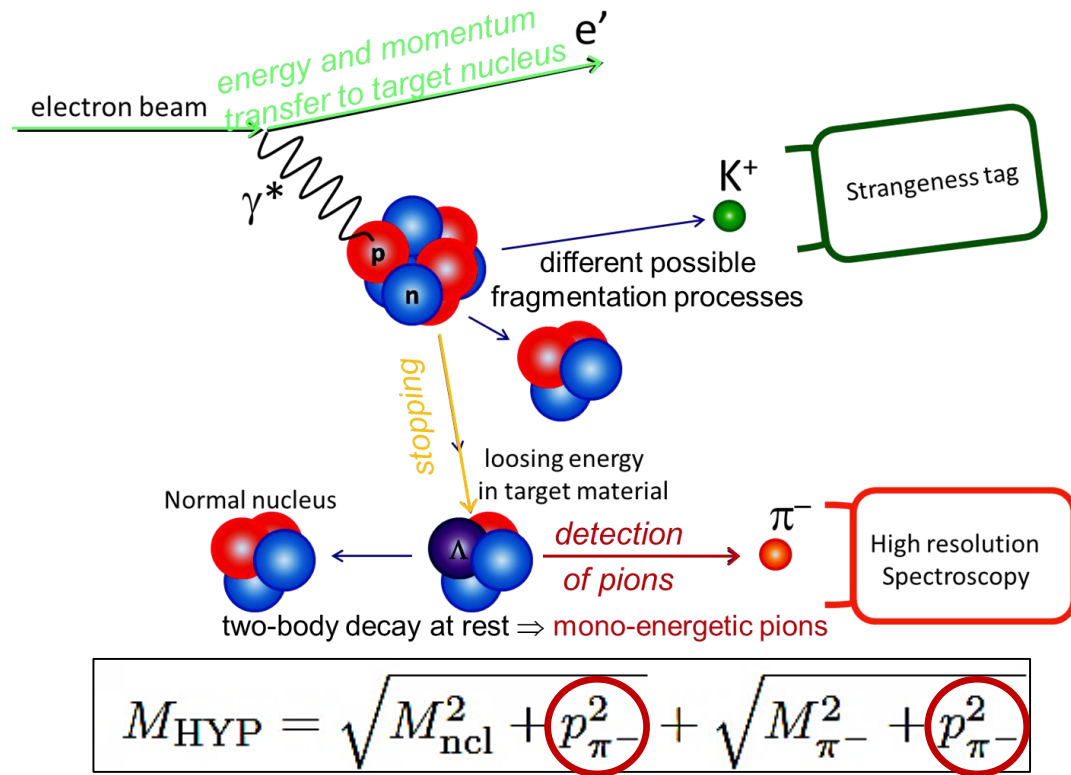
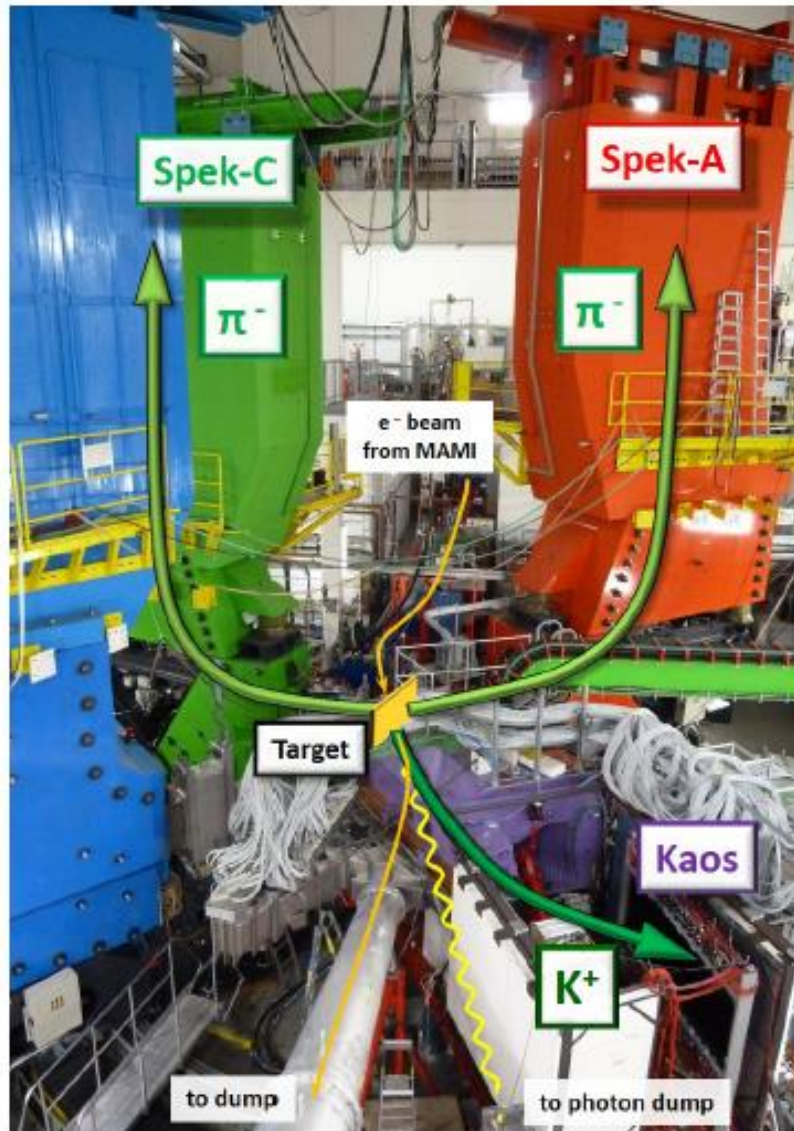
emulsion



$$\left. \begin{aligned}
 &{}^4_{\Lambda}\text{He} \xrightarrow{\text{decay}} \pi^- + {}^1\text{H} + {}^3\text{He}: B = 2.42 \pm 0.05 \text{ MeV} \\
 &{}^4_{\Lambda}\text{He} \xrightarrow{\text{decay}} \pi^- + 2{}^1\text{H} + {}^2\text{H}: B = 2.44 \pm 0.09 \text{ MeV}
 \end{aligned} \right\} 0.02 \text{ MeV difference}$$

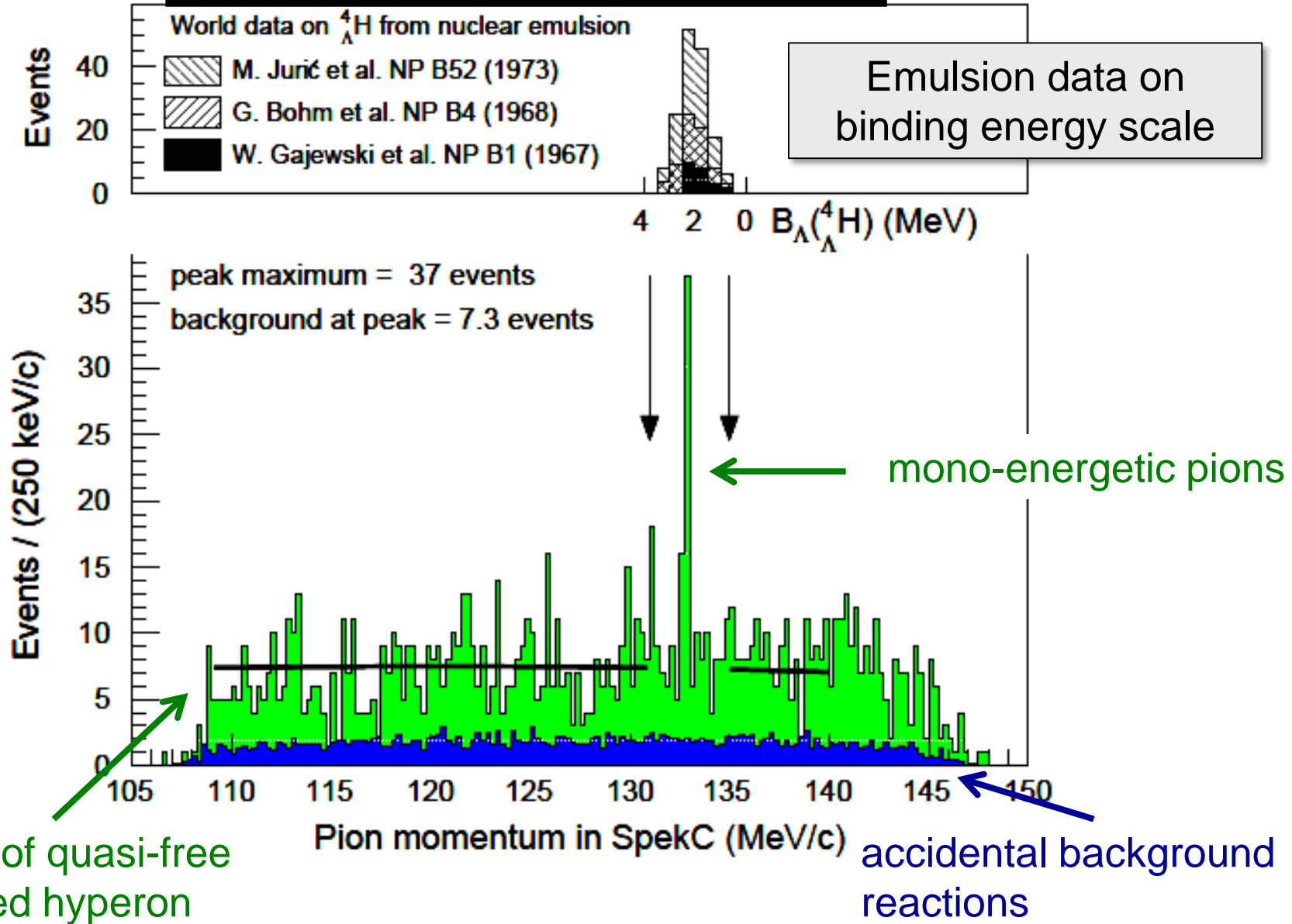
Total: $B = 2.42 \pm 0.04 \text{ MeV}$ [M. Juric et al. NP B52 (1973)]

Hyperfragment decay-pion spectroscopy with electron beams

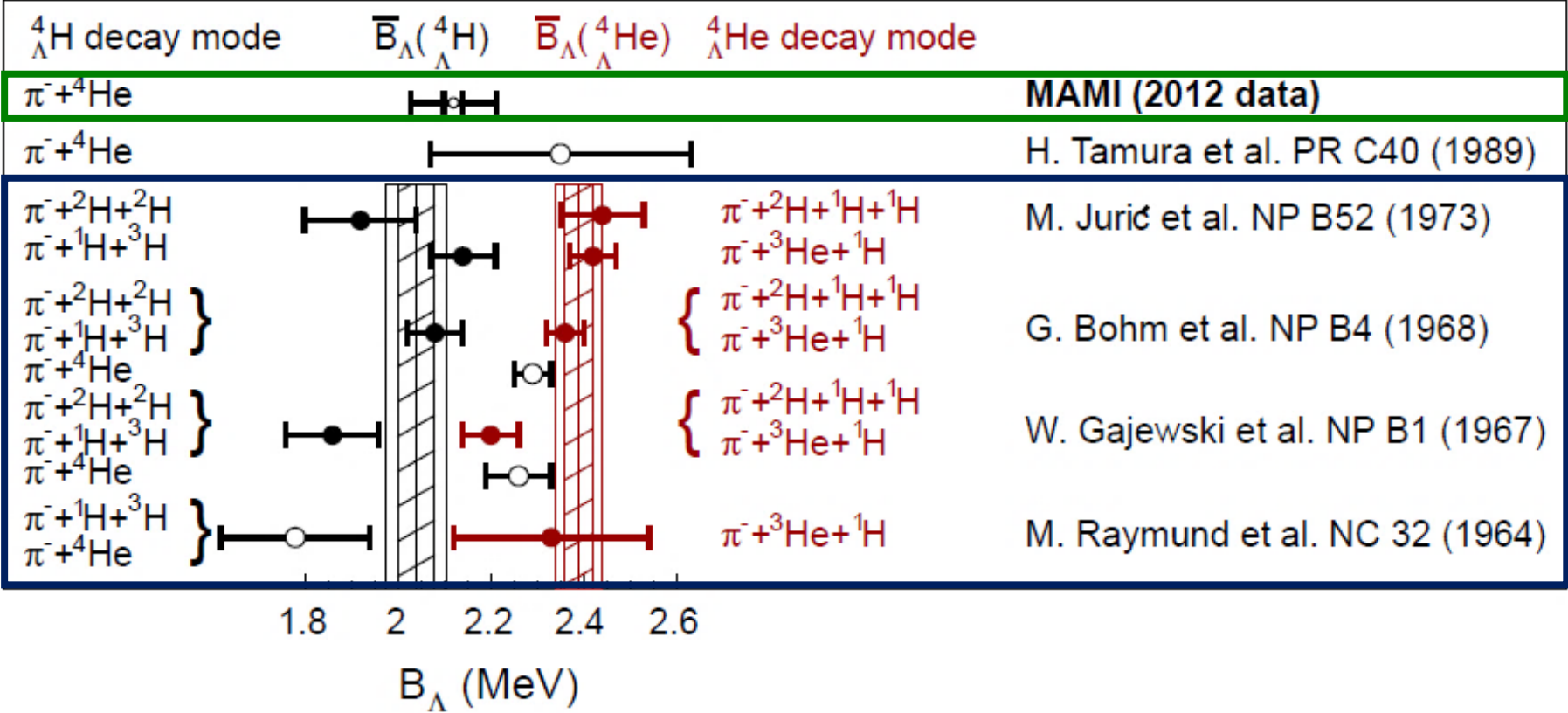


Hypernuclear experiments at MAMI done in Collaboration with Tohoku University (S.N.N. Nakamura *et al.*)

Decay-pion spectrum



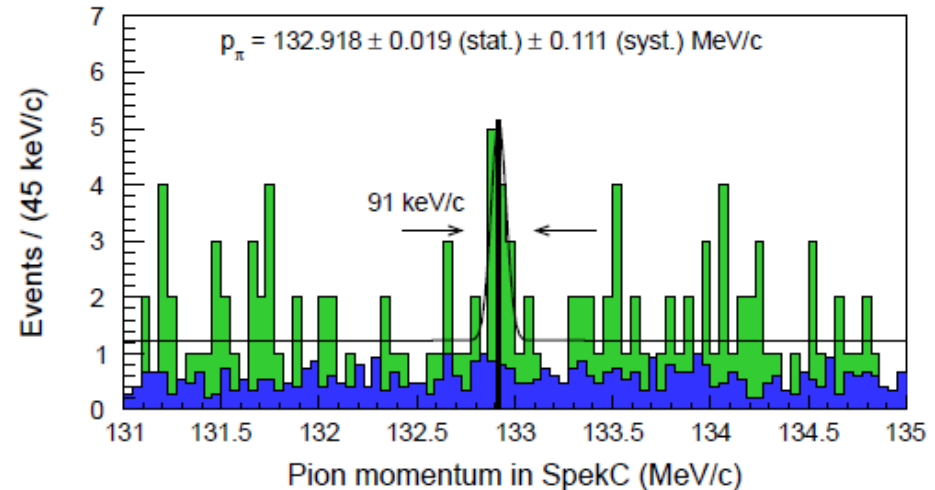
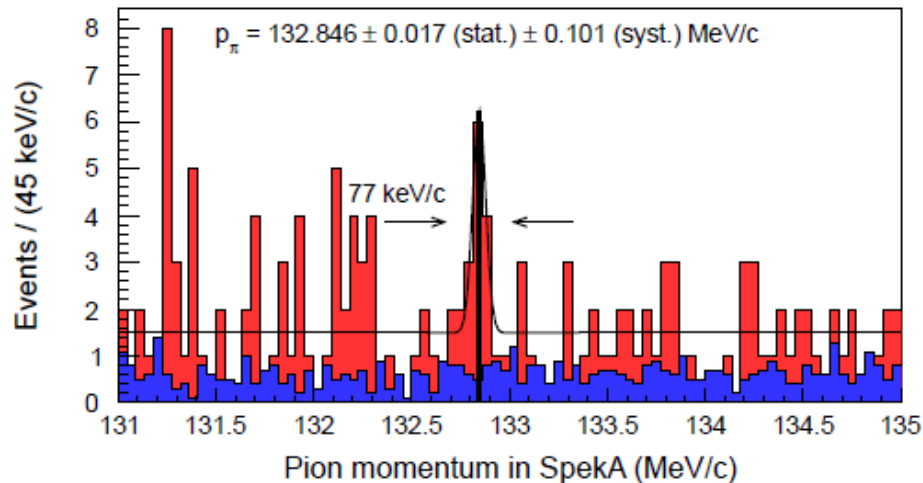
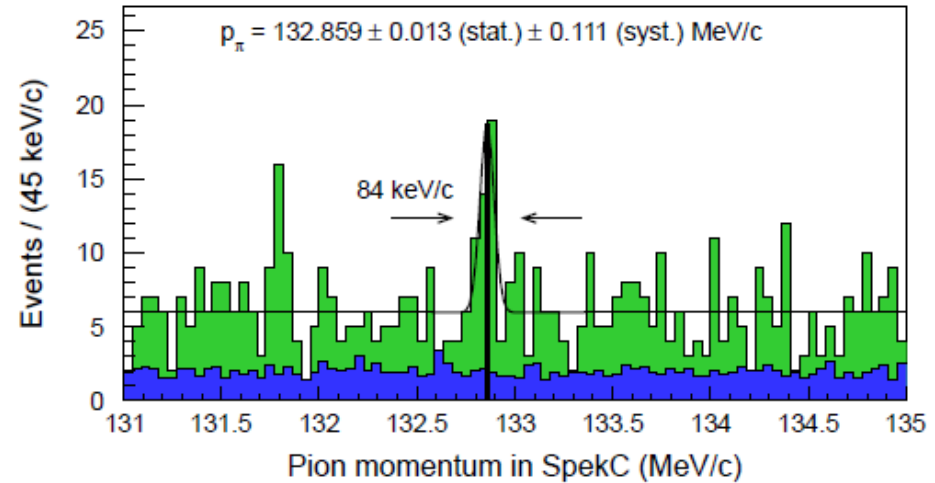
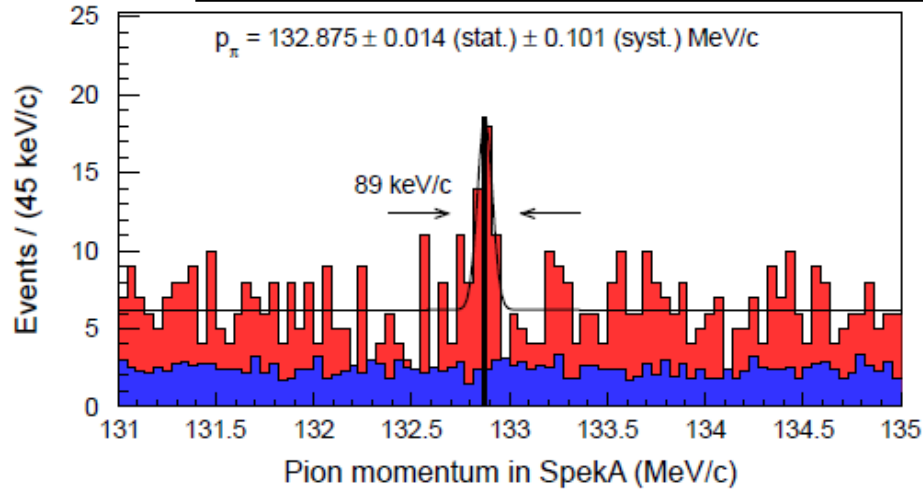
World data on $A = 4$ system



MAMI 2012 experiment Λ binding energy of ${}^4_\Lambda\text{H}$:
 $B_\Lambda = 2.12 \pm 0.01$ (stat.) ± 0.09 (syst.) MeV

A. Esser et al. (A1 Collaboration), PRL 114, 232501 (2015)

Systematic studies of the decay-pion line

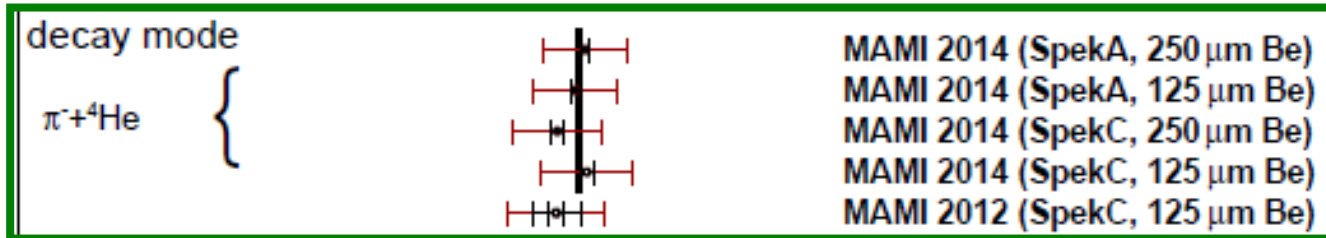


- consistent result for $B_{\Lambda}(^4_{\Lambda}\text{H})$ from MAMI 2012 and MAMI 2014 independent measurement in two specs, two targets, two beam-times

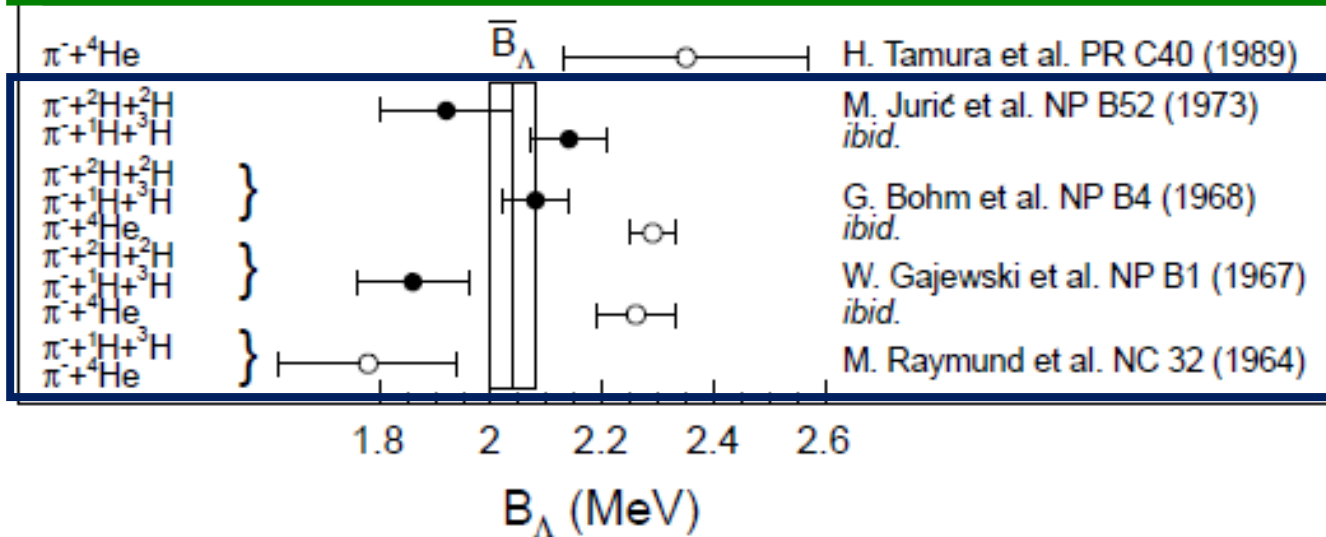
World data on ${}^4_{\Lambda}\text{H}$ mass

outer error bars correlated from calibration

MAMI

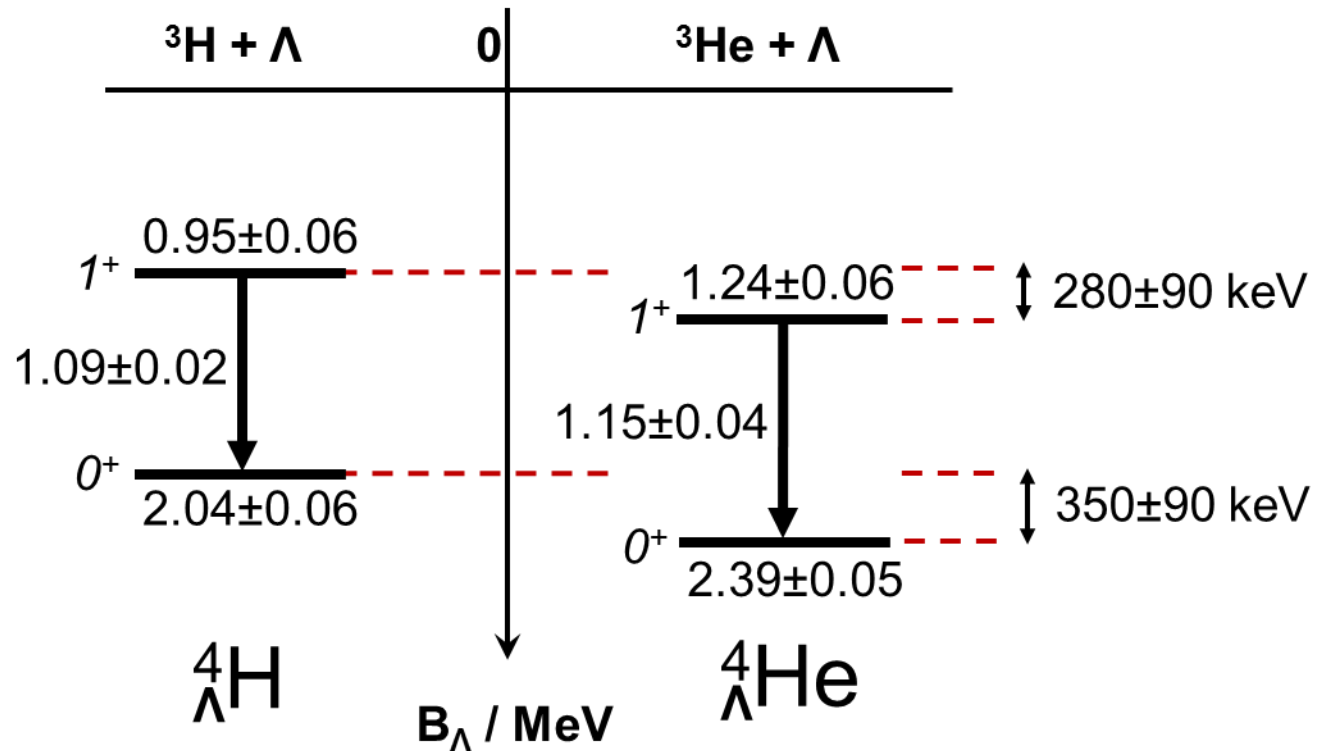


emulsion



	$B_{\Lambda}({}^4_{\Lambda}\text{H})$ (stat.)	(syst.)	
emulsion:	2.04	± 0.04	± 0.05 MeV [M. Juric et al. NP B52 (1973)]
MAMI 2012:	2.12	± 0.01	$\pm 0.08 \pm 0.03$ MeV [A. Esser et al., PRL 114 (2015)]
MAMI 2012:	2.12	± 0.01	$\pm 0.08 \pm 0.03$ MeV [P. Achenbach et al., JPS (2016)]
MAMI 2014:	2.16	± 0.01	± 0.08 MeV [F. Schulz et al., NPA 954 (2016)]

The $A = 4$ level schemes (before 2015)

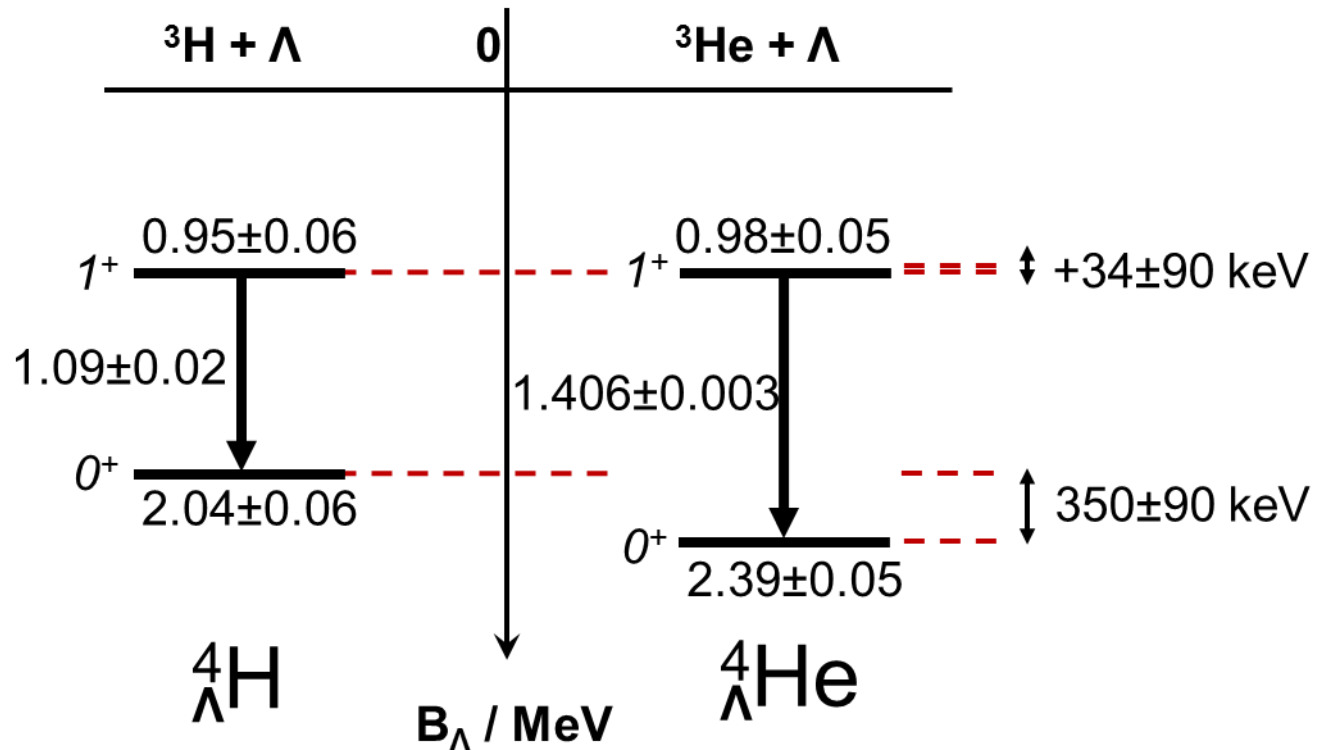


$$\Delta B_{\Lambda} (1+ \text{ ex. st.}) = 280 \pm 90 \text{ keV} \quad (\textit{emulsion} + \textit{old } \gamma\text{-ray data})$$

$$\Delta B_{\Lambda} (0+ \text{ gr. st.}) = 350 \pm 90 \text{ keV} \quad (\textit{emulsion data})$$

- charge symmetry breaking only if emulsion data is correct
- spin-independent charge symmetry breaking likely

The $A = 4$ level schemes (as of 2015)



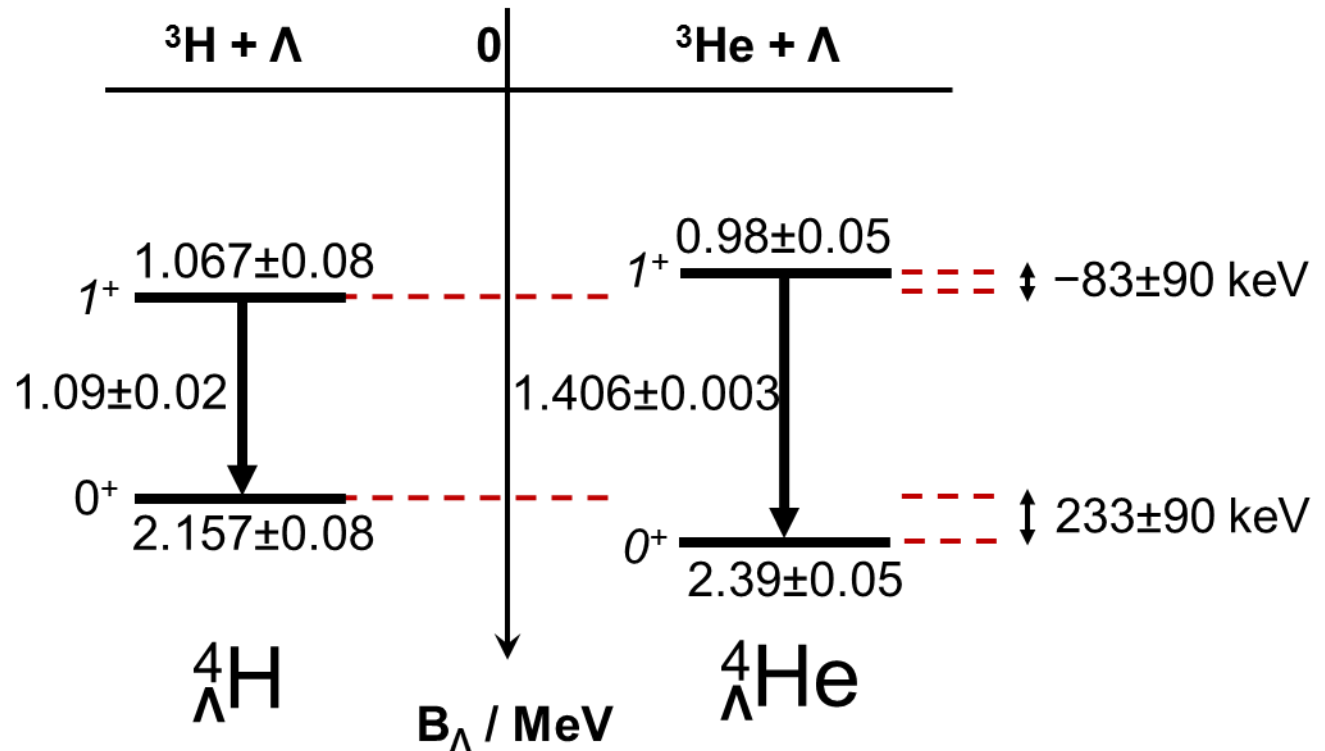
$$\Delta B_{\Lambda} (1+ \text{ ex. st.}) = +34 \pm 90 \text{ keV} \quad (\textit{emulsion} + \textit{new } \gamma\text{-ray data})$$

[T. O. Yamamoto, PRL 115 (2015)]

$$\Delta B_{\Lambda} (0+ \text{ gr. st.}) = 350 \pm 90 \text{ keV} \quad (\textit{emulsion data})$$

- observation of charge symmetry breaking in γ -ray spectroscopy
- spin-dependent charge symmetry breaking

Current knowledge on CSB in the $A = 4$ system

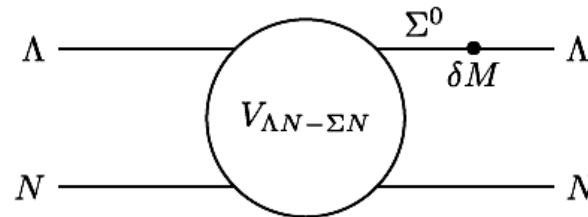


- $\Delta B_{\Lambda} (1+ \text{ ex. st.}) = -83 \pm 90 \text{ keV}$ (*MAMI + emulsion + new γ -ray data*)
 [F. Schulz et al., NPA 954 (2016)]
- $\Delta B_{\Lambda} (0+ \text{ gr. st.}) = 233 \pm 90 \text{ keV}$ (*MAMI + emulsion data*)

- CSB considerably stronger than in other hyper- or ordinary nuclei
- CSB possibly changing sign between ground and excited states

Consistent experimental data with latest theory

latest chiral effective models with central force ΛN - ΣN coupling
→ mixing of $l=0$ and $l=1$ hyperons leads to long-range pion exchanges



based on OBE Nijmegen NSC97 model YN potential

$$\Delta B_{\Lambda} (0+ \text{ gr. st.}) = +226/266 \text{ keV}$$

$$\Delta B_{\Lambda} (1+ \text{ ex. st.}) = +30/39 \text{ keV}$$

[A. Gal, PLB 744, 352 (2015)]

based on LO chiral EFT Bonn-Jülich YN potential

$$\Delta B_{\Lambda} (0+ \text{ gr. st.}) = +180 \pm 130 \text{ keV}$$

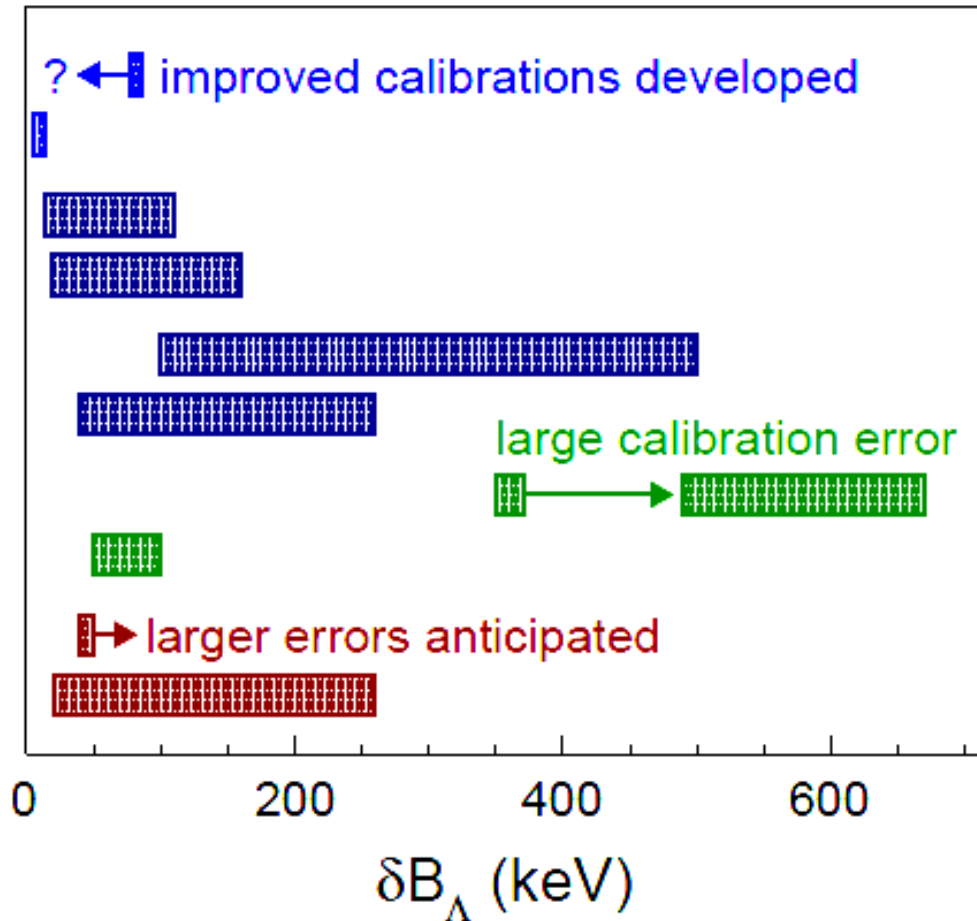
$$\Delta B_{\Lambda} (1+ \text{ ex. st.}) = -200 \pm 30 \text{ keV}$$

[D. Gazda & A. Gal, PRL 116, 122501 (2016)]

[D. Gazda & A. Gal, NPA 954, 161 (2016)]



Errors on binding energy by method



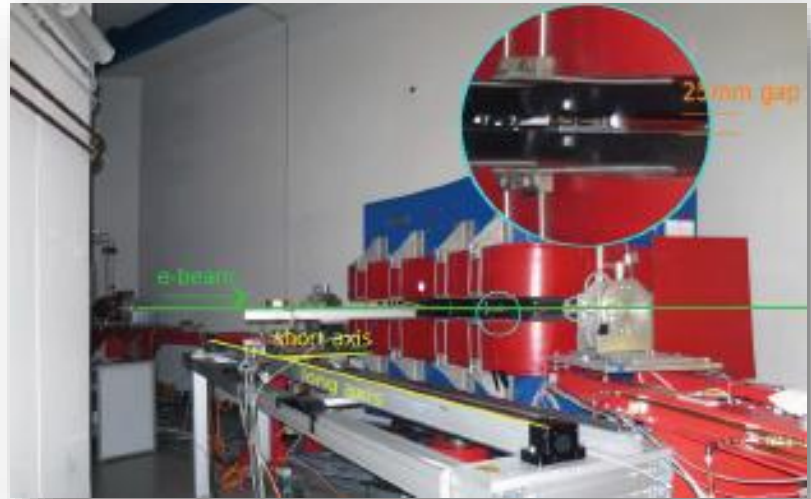
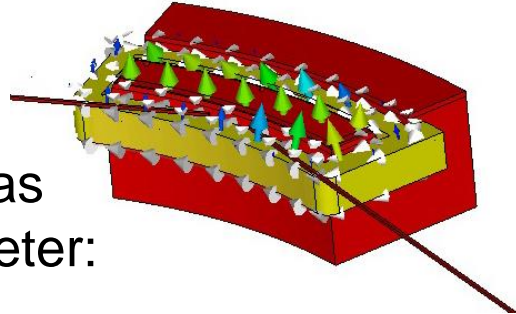
$(e, K^+ \pi^-)$ MAMI syst. error
 $(e, K^+ \pi^-)$ MAMI stat. error
 $(e, e' K^+)$ JLab syst. error
 $(e, e' K^+)$ JLab stat. error
 $(K_{\text{stop}}^-, \pi^-)$ FINUDA syst. error
 $(K_{\text{stop}}^-, \pi^-)$ FINUDA stat. error
 (π^+, K^+) KEK-SKS syst. error
 (π^+, K^+) KEK-SKS stat. error
 emulsion syst. error
 emulsion stat. error

- goal of calibrations is syst. error comparable to stat. error < 20 keV
- decay-pion spectroscopy will be the most precise method of all

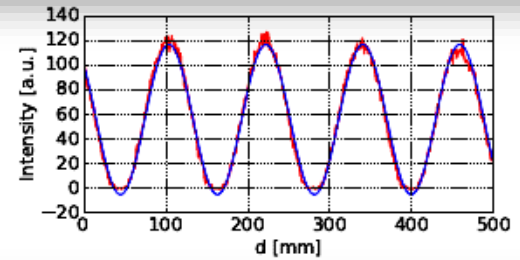
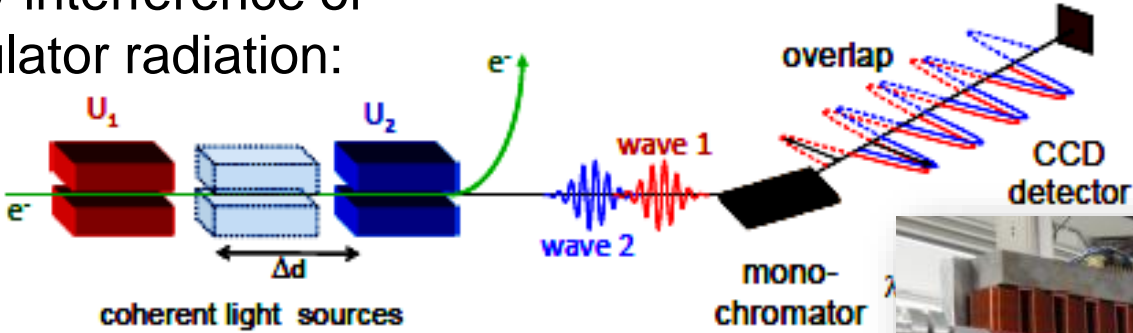
Reducing the systematic error

MAMI energy measurement with $O(10^{-4})$ precision

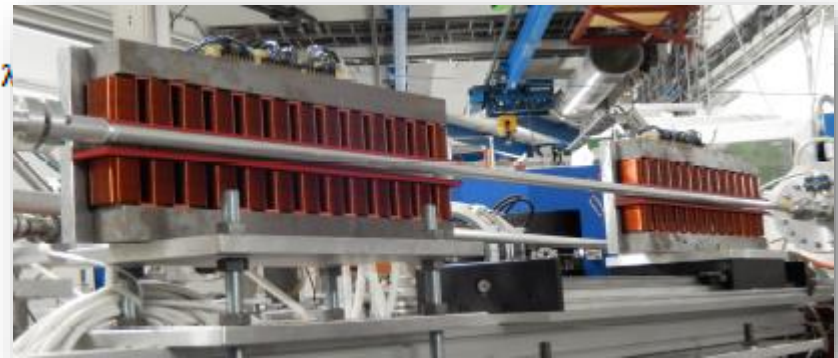
1) by a dipole used as beam-line spectrometer:



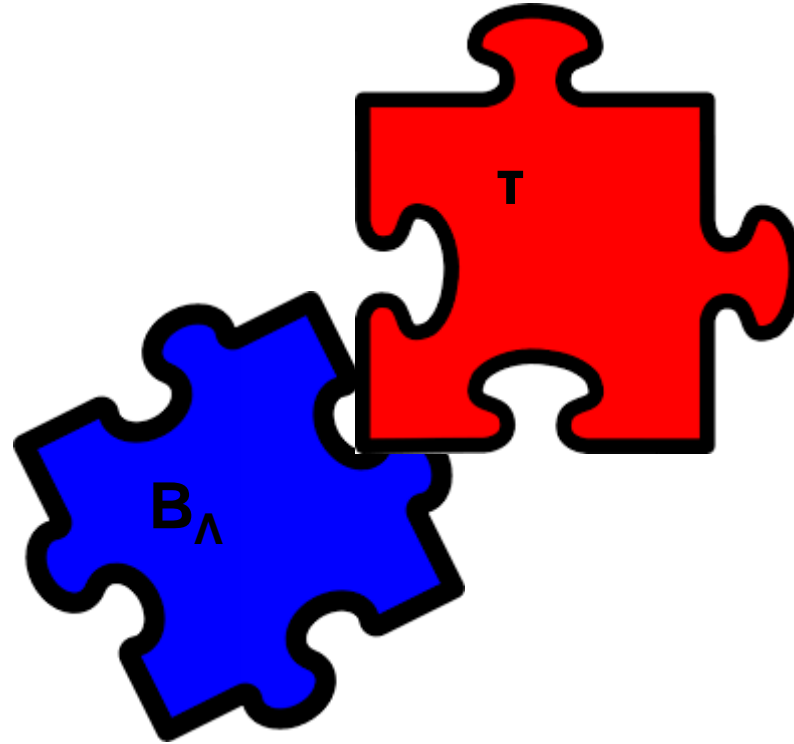
2) by interference of undulator radiation:



$$\gamma_{beam}^2 = \frac{1}{2} \frac{\lambda_{osc}}{\lambda_{rad}}$$

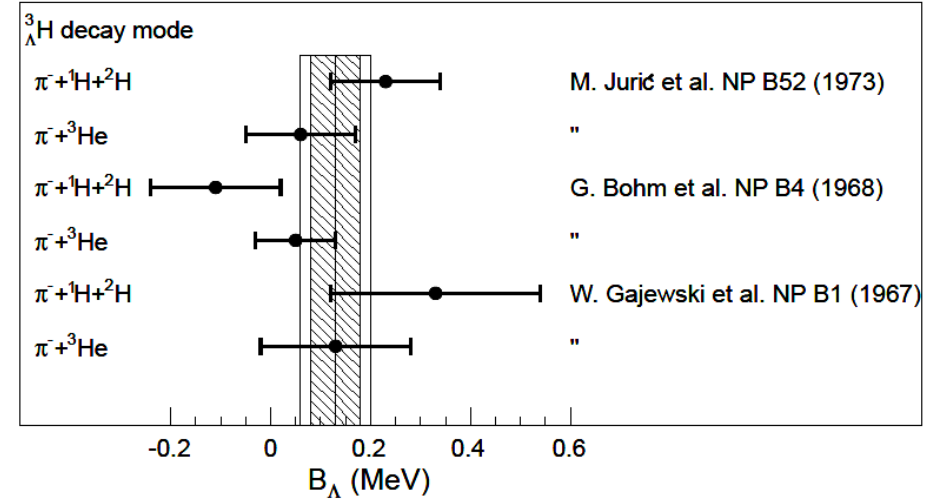
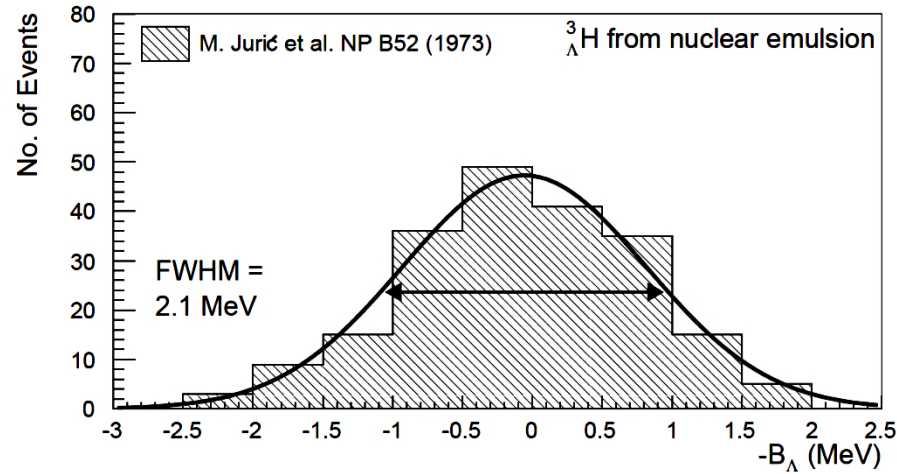


The lightest Λ hypernucleus and the puzzle of lifetime



The quest for ${}^3_{\Lambda}\text{H}$

only about 200 analyzed events from emulsion:



these data (from 2 decay modes): only source of binding energy information

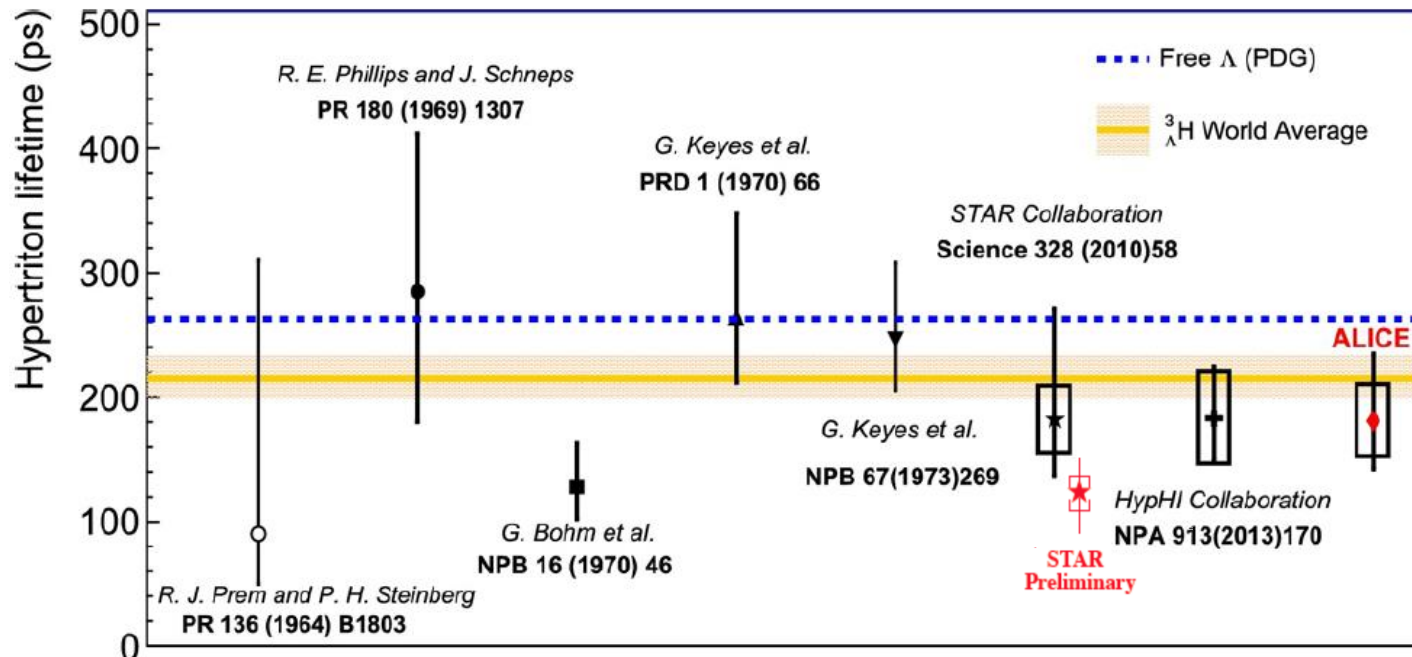
$$\left. \begin{array}{l}
 {}^3_{\Lambda}\text{H} \xrightarrow{\text{decay}} \pi^- + {}^3\text{He}: \quad B = 0.07 \pm 0.06 \text{ MeV} \\
 {}^3_{\Lambda}\text{H} \xrightarrow{\text{decay}} \pi^- + {}^1\text{H} + {}^2\text{H}: \quad B = 0.12 \pm 0.08 \text{ MeV}
 \end{array} \right\} 0.05 \text{ MeV difference}$$

Total: $B = 0.13 \pm 0.05 \text{ MeV}$ [M. Juric et al. NP B52 (1973)]

$^3_\Lambda\text{H}$ lifetime

Small Λ binding energy implies extended wave function \rightarrow hypernucleus lifetime should be comparable to free Λ lifetime

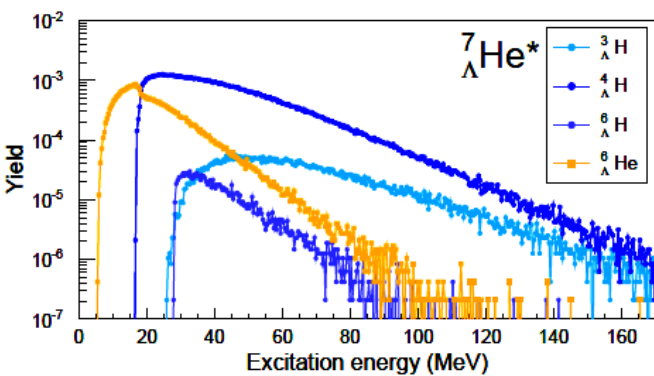
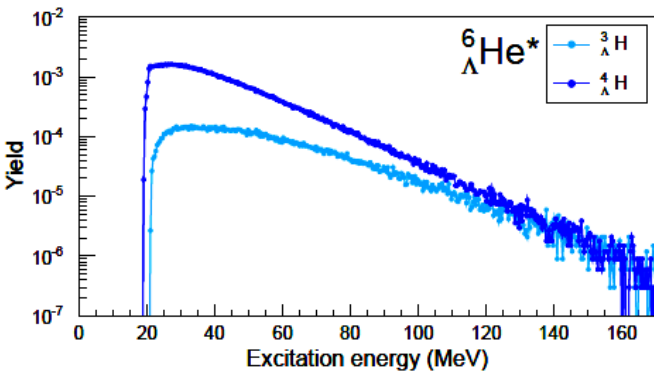
all calculations of lifetime predict no more than 10% deviation



Values from heavy ion experiments are surprisingly small

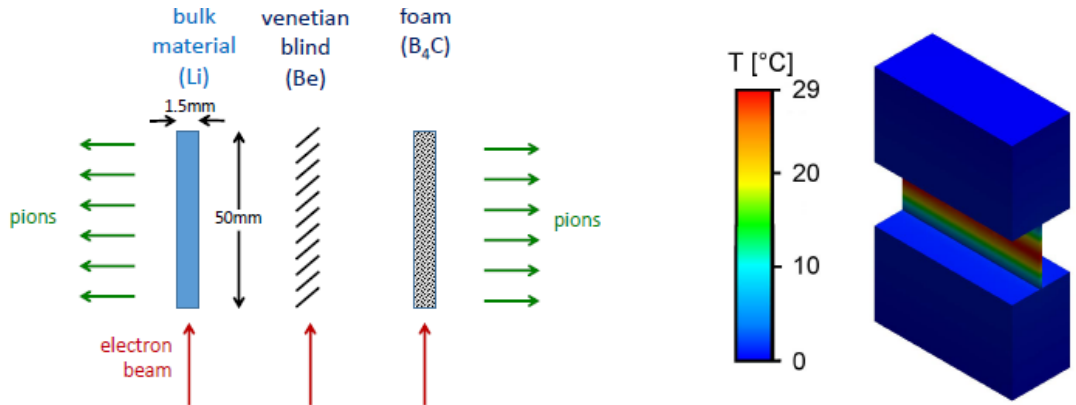
Li: an ideal target for ${}^3_{\Lambda}\text{H}$ production

Statistical decay calculations



D_C and $P(x)$ in units of 10^{-4} , x in units of mm

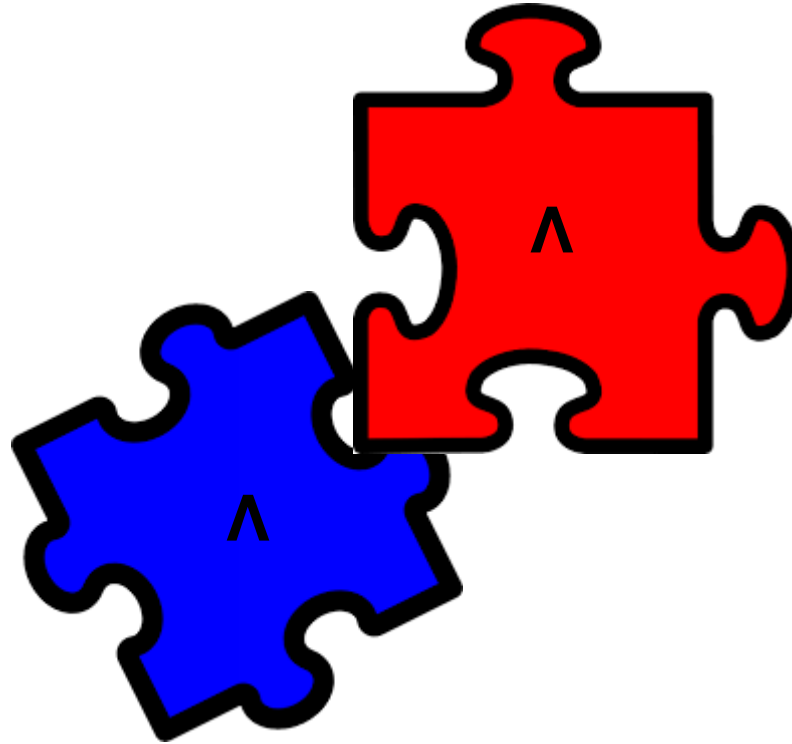
	$E_x = 20 \text{ MeV}$			$E_x = 40 \text{ MeV}$			$E_x = 60 \text{ MeV}$		
${}^6_{\Lambda}\text{He}^*$	D_C	$P(0.25)$	$P(1.5)$	D_C	$P(0.25)$	$P(1.5)$	D_C	$P(0.25)$	$P(1.5)$
${}^3_{\Lambda}\text{H}$	—	—	—	27.	1.4	3.7	17.	1.0	2.1
${}^4_{\Lambda}\text{H}$	38.	8.1	17.	56.	11.	22.	20.	3.8	7.5
${}^7_{\Lambda}\text{He}^*$	D_C	$P(0.25)$	$P(1.5)$	D_C	$P(0.25)$	$P(1.5)$	D_C	$P(0.25)$	$P(1.5)$
${}^3_{\Lambda}\text{H}$	—	—	—	9.2	0.4	1.3	8.5	0.5	1.1
${}^4_{\Lambda}\text{H}$	55.	11.	23.	48.	9.3	18.	22.	4.1	8.3
${}^6_{\Lambda}\text{H}$	—	—	—	1.3	0.2	0.1	0.3	< 0.1	< 0.1
${}^6_{\Lambda}\text{He}$	36.	4.8	5.6	6.9	0.9	1.1	1.1	0.1	0.2



expected excitation energy

- convert proton into $\Lambda \Rightarrow$ proton hole state $\sim 20 \text{ MeV}$
- kinetic energy of captured Λ : $p_F^2 / 2M_{\Lambda} \sim 20 \text{ MeV}$
- Binding energy of $\Lambda < 10 \text{ MeV}$

Double Λ hypernuclei and the puzzle of $\Lambda\Lambda$ interaction



S=-2 systems

missing mass (K^-, K^+) reactions \Rightarrow Ξ bound state

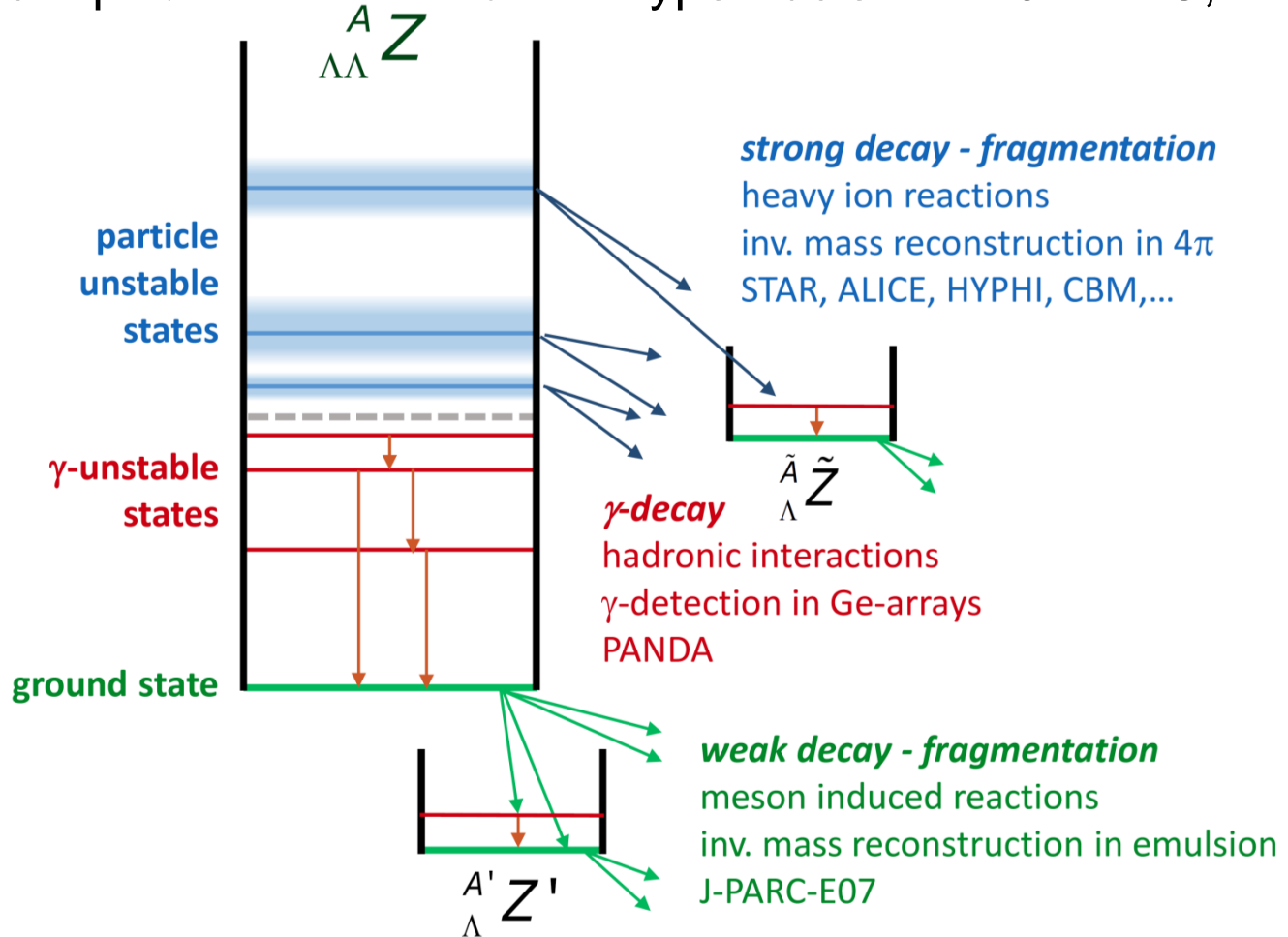
J-PARC

Ξ capture \Rightarrow Ξ atoms

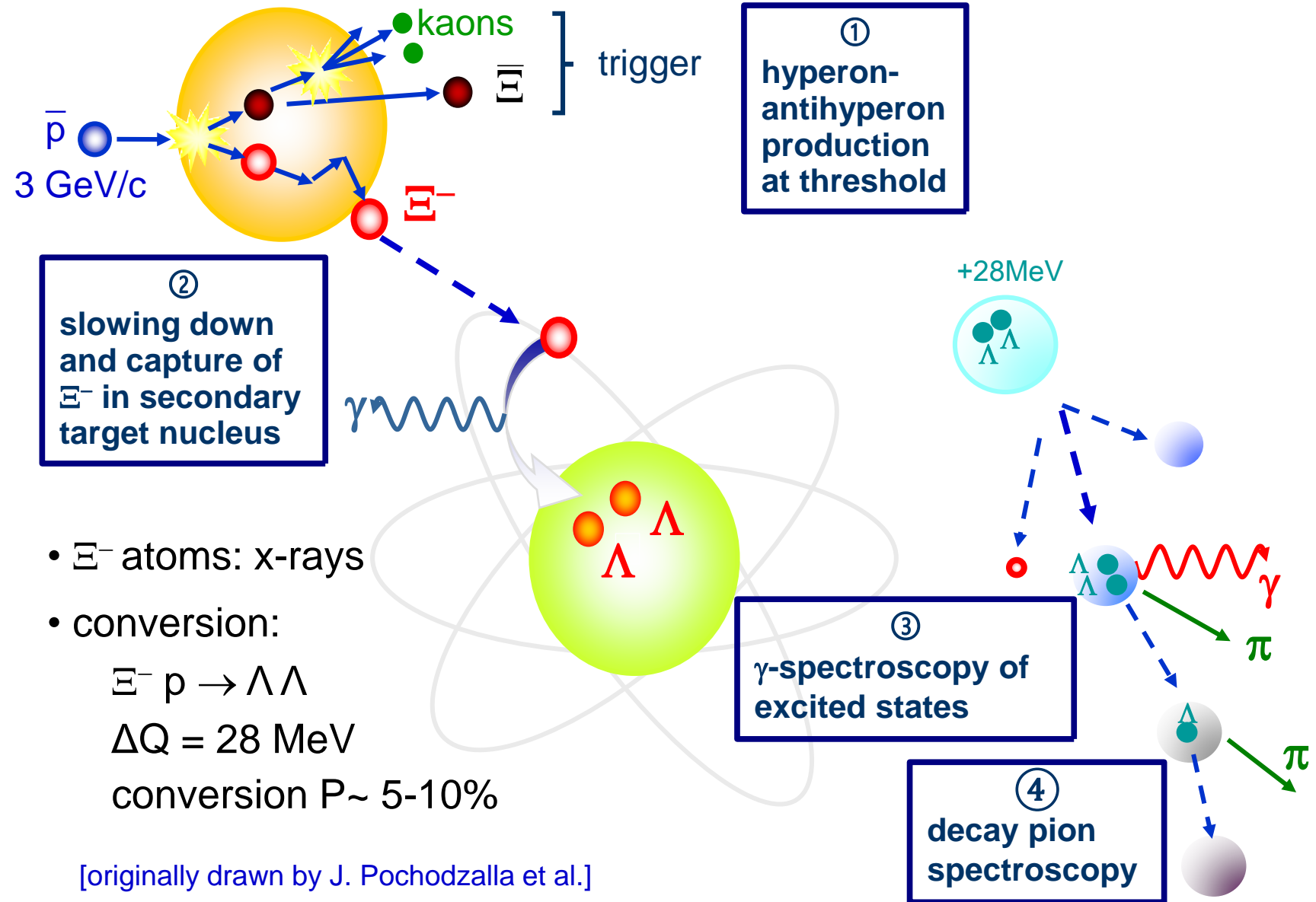
J-PARC, PANDA

Ξ capture and $\Xi p \rightarrow \Lambda\Lambda$ \Rightarrow $\Lambda\Lambda$ hypernuclei

J-PARC, PANDA, HI

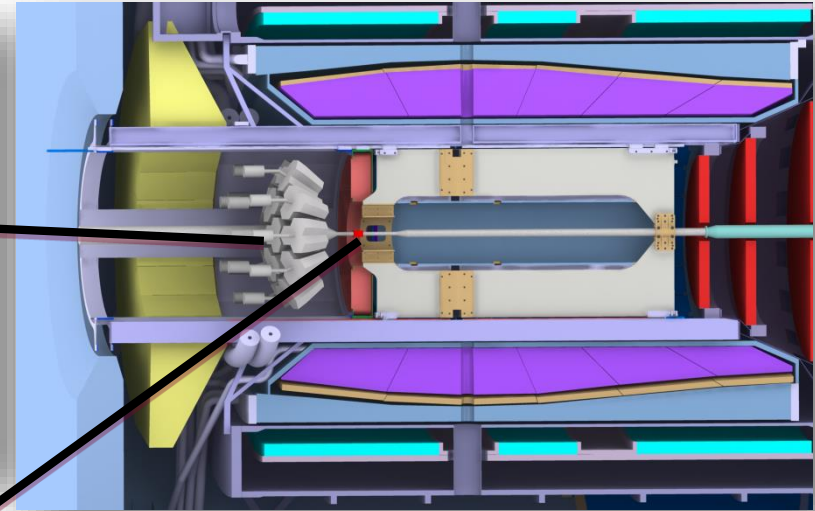
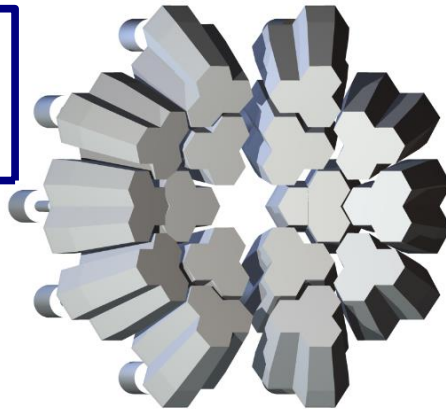


Production mechanism and detection strategy at PANDA

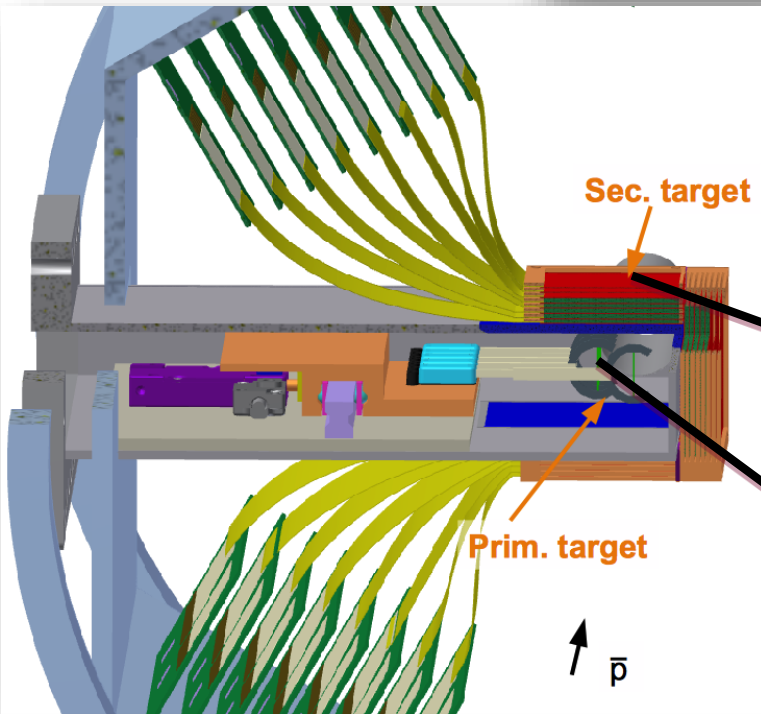


Instrumentation for hypernuclear physics at PANDA

Germanium detector array for hypernuclei spectroscopy

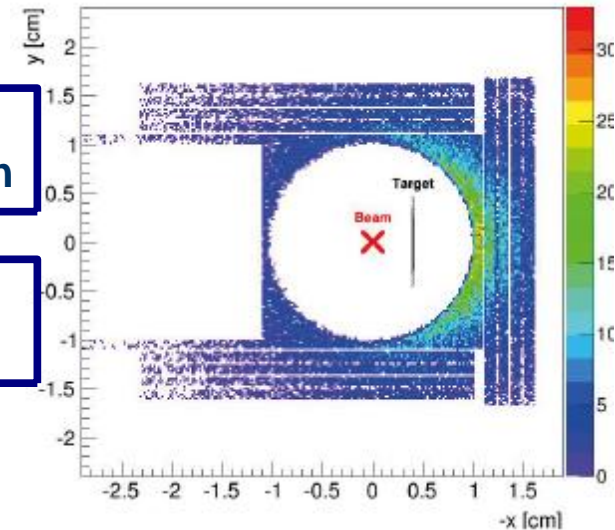


the Panda detector for forward angle particle identification



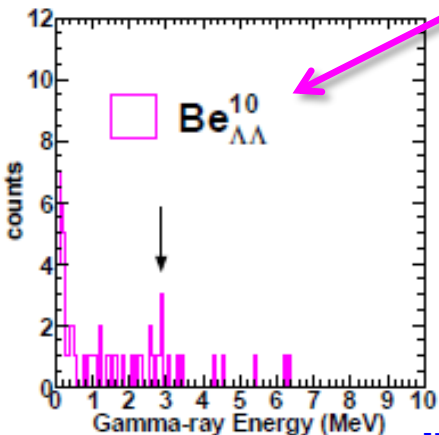
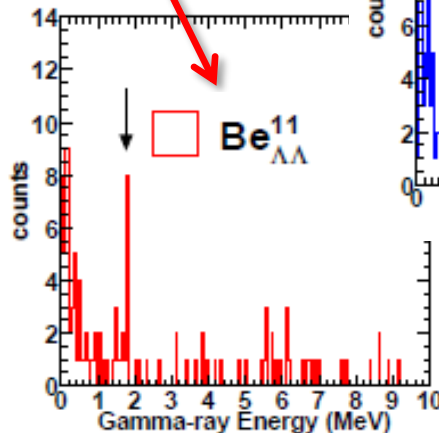
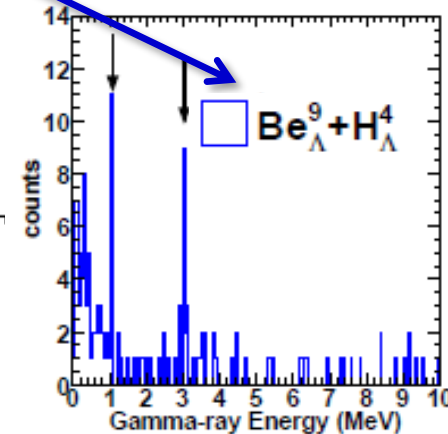
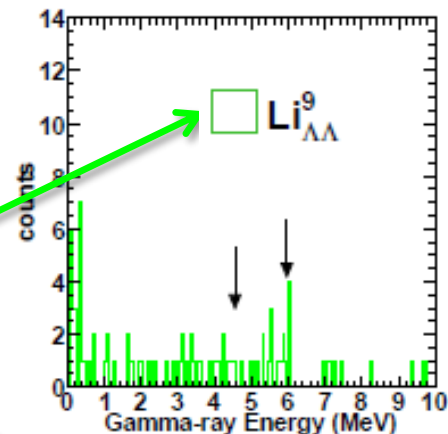
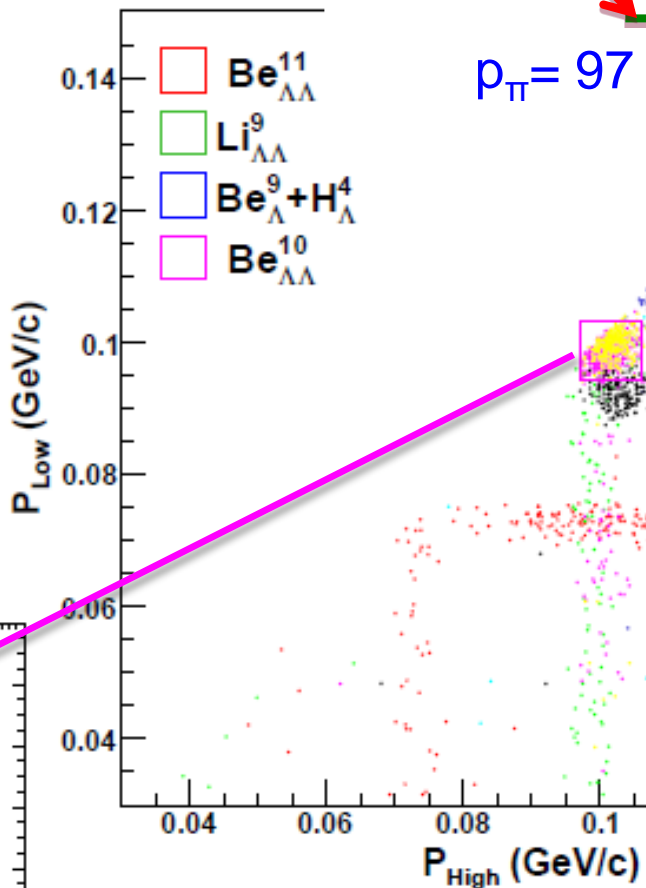
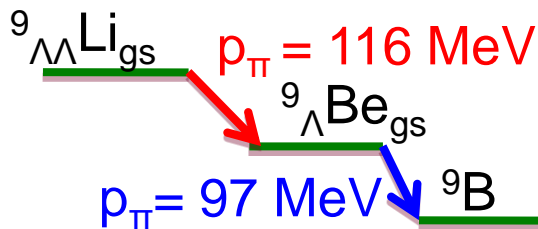
secondary target for hypernuclei formation

primary target for Ξ production



Background suppression by decay pion correlation

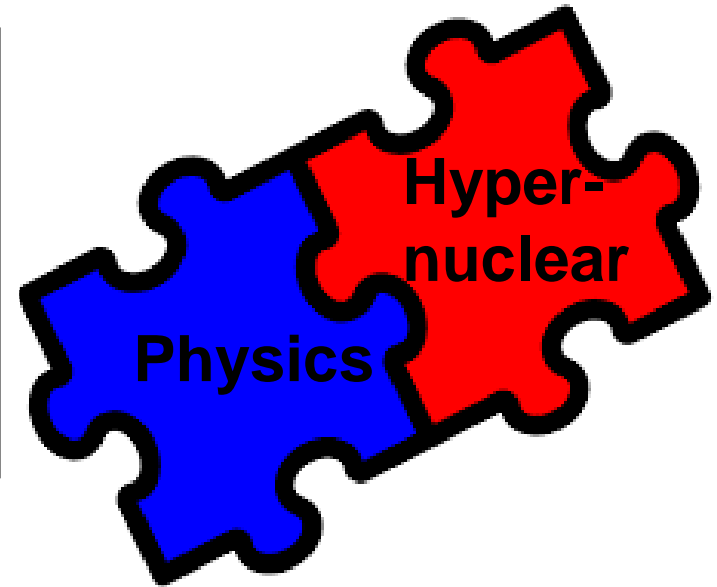
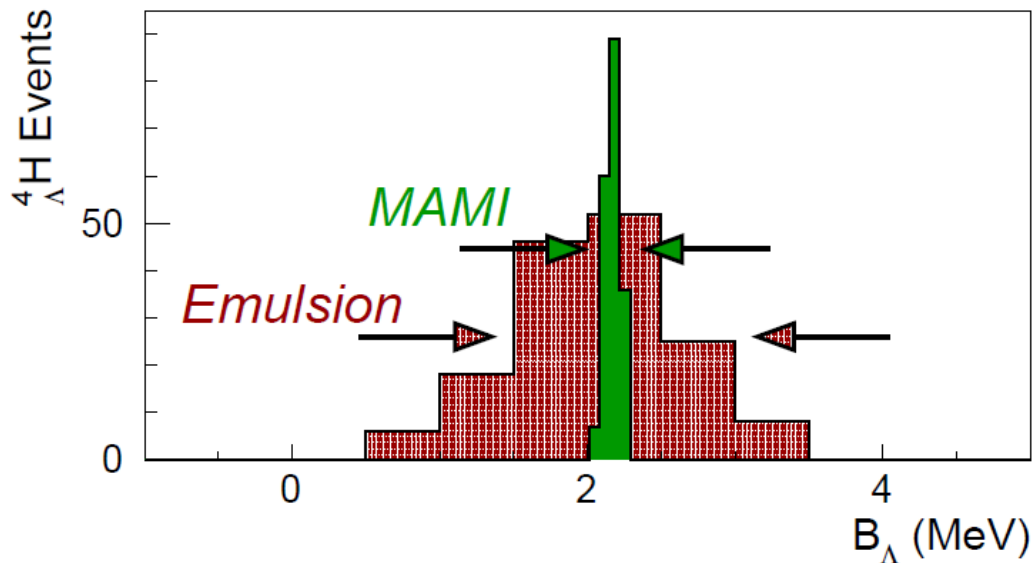
Data from simulations:



[PANDA Physics Performance Report]

New precision era in hypernuclear physics

1. High-resolution pion spectroscopy has replaced emulsion technique
2. High-purity germanium spectroscopy has replaced NaJ detectors



Hypernuclear physics at MAMI and PANDA is complementary to J-PARC activities