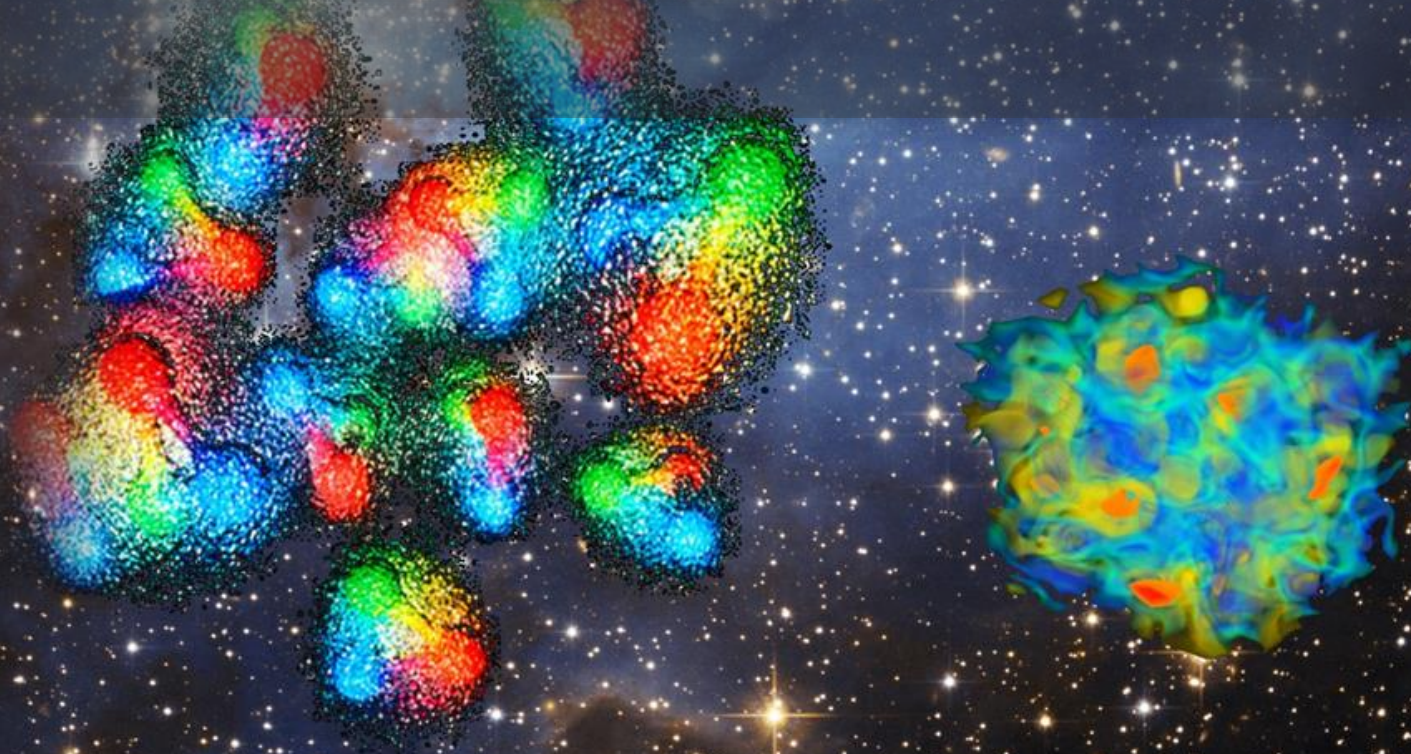


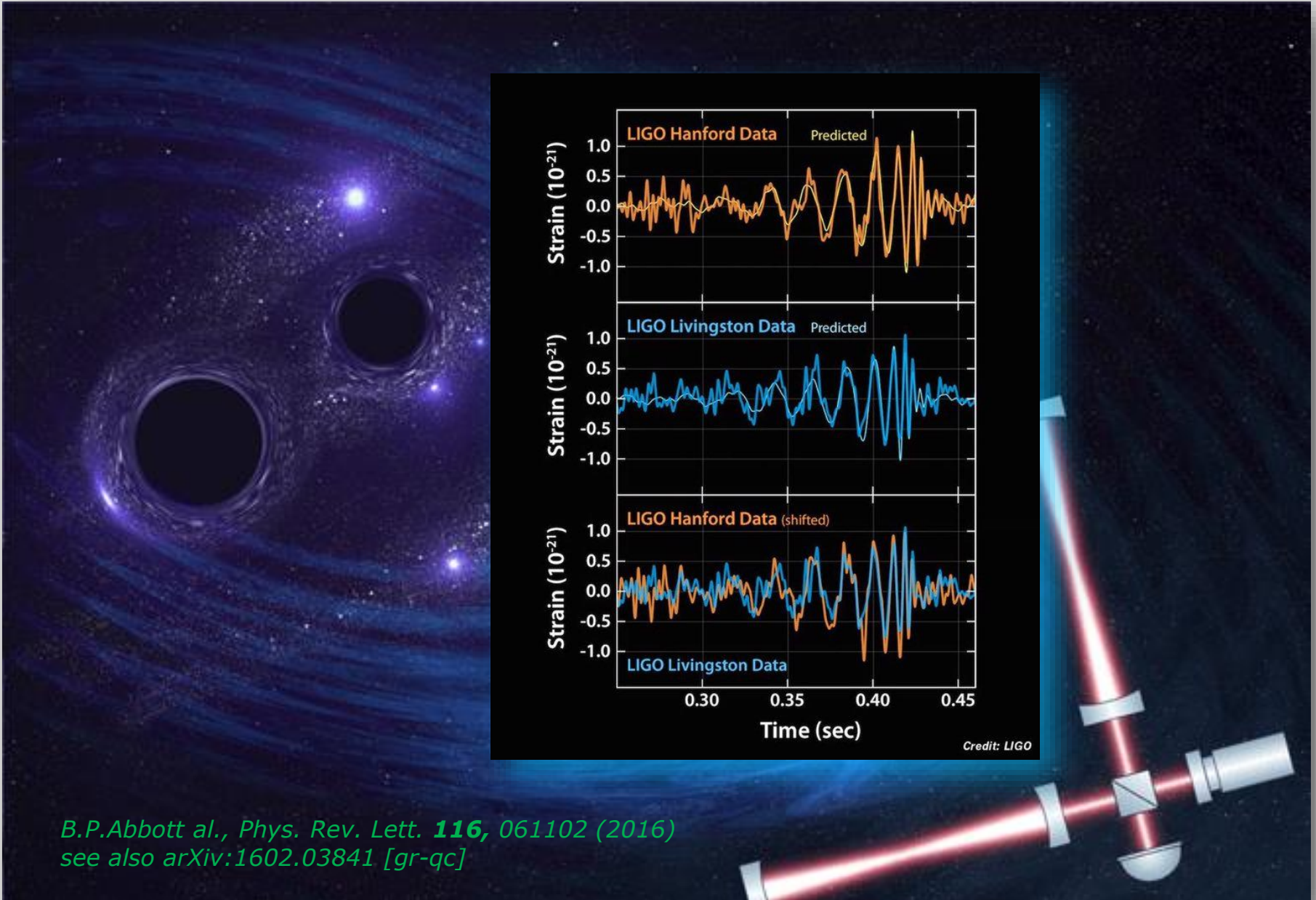
It's a strange world where QCD meets Gravity



Josef Pochodzalla

Helmholtz-Institut Mainz





B.P.Abbott et al., *Phys. Rev. Lett.* **116**, 061102 (2016)
see also [arXiv:1602.03841 \[gr-qc\]](https://arxiv.org/abs/1602.03841)

A visualization of gravitational waves from a neutron star merger. A large, textured sphere representing the merged neutron star is at the bottom center. A smaller, white sphere representing the original neutron star is at the top left. Concentric blue circles represent the expanding gravitational wave front. A yellow waveform, representing the detected signal, is overlaid on the scene, showing a series of peaks that decrease in amplitude as they move away from the source.
$$M(\text{PSR J1614-2230}) = 1.97 \pm 0.04 M_{\odot}$$

$$M(\text{PSR J0348+0432}) = 2.01 \pm 0.04 M_{\odot}$$

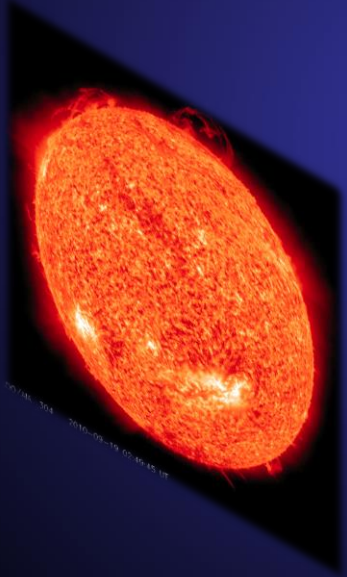
P. B. Demorest et al., Nature 467 (2010)
John Antoniadis et al., Science 340 (2013)

$$\frac{2GM}{c^2 R}$$

$\sim 10^{-10}$



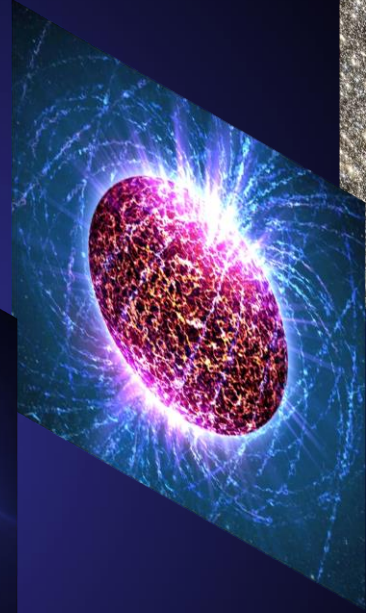
$\sim 10^{-7}$



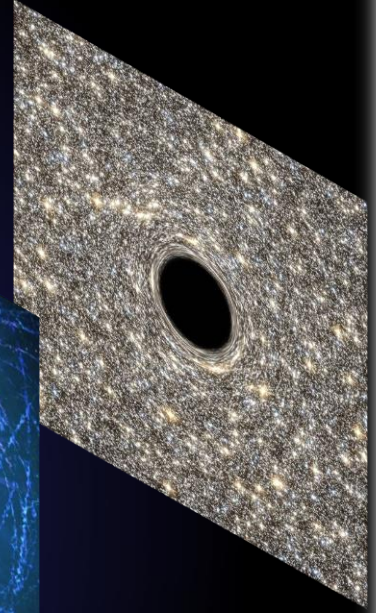
$\sim 10^{-4}$



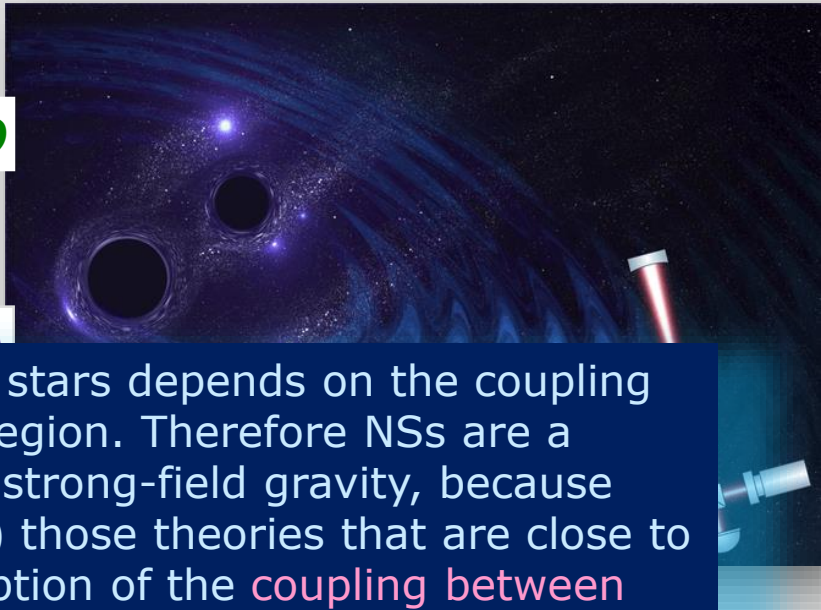
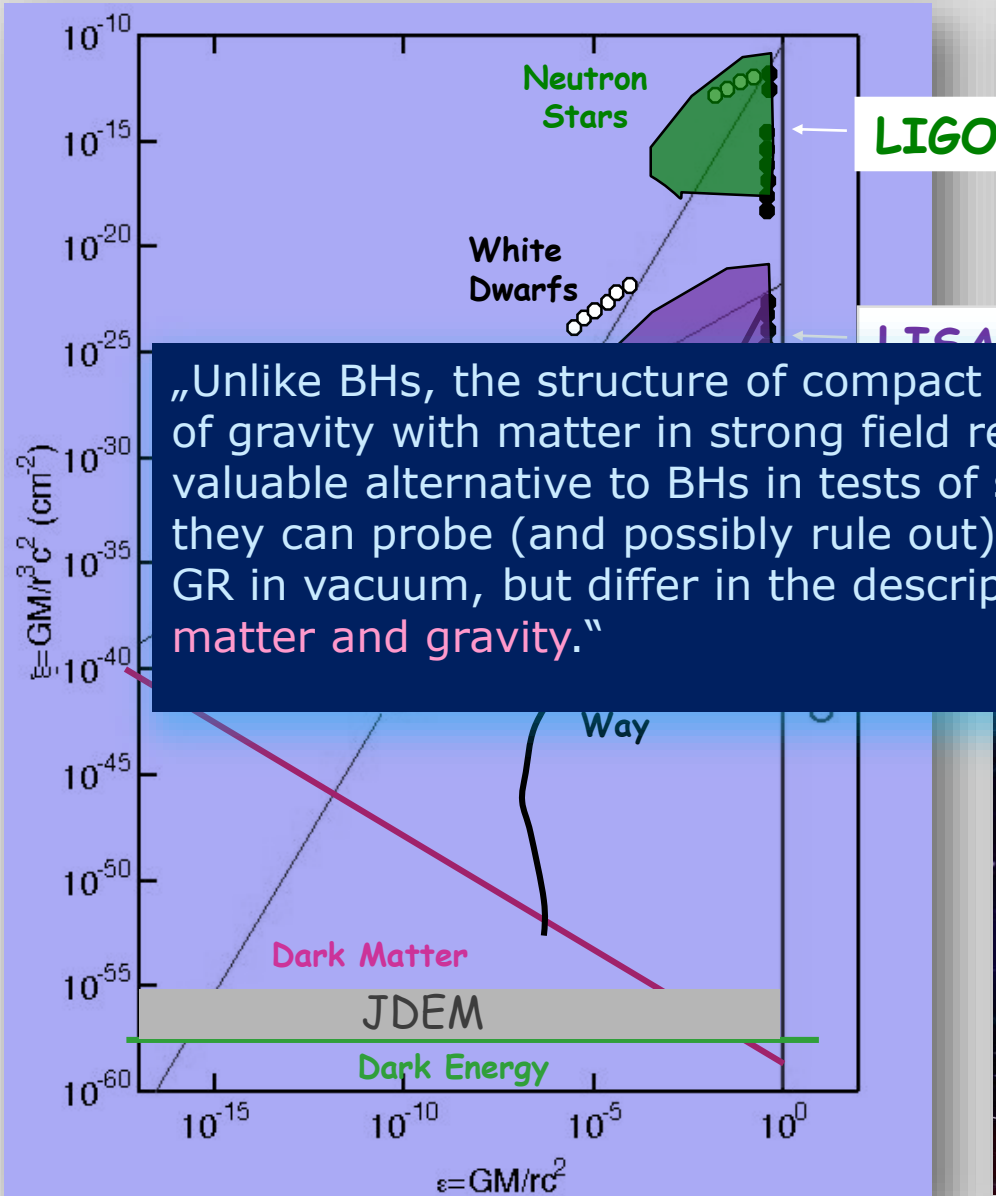
~ 0.3



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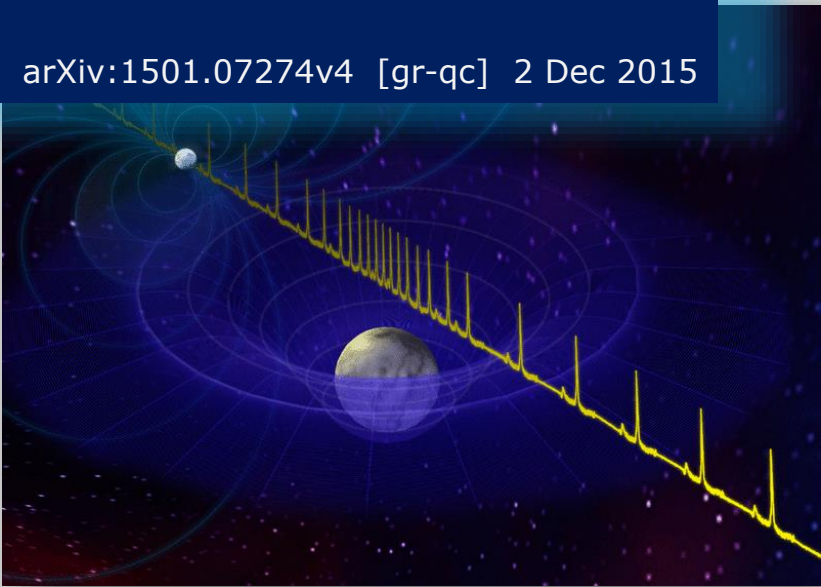


D. Psaltis, *Living Rev. Relativity* **11**, 9 (2008)



„Unlike BHs, the structure of compact stars depends on the coupling of gravity with matter in strong field region. Therefore NSs are a valuable alternative to BHs in tests of strong-field gravity, because they can probe (and possibly rule out) those theories that are close to GR in vacuum, but differ in the description of the **coupling between matter and gravity.**“

arXiv:1501.07274v4 [gr-qc] 2 Dec 2015



surface gravitational potential

Neutron Stars: Hyperon Puzzle

Available online at www.sciencedirect.com
ScienceDirect
Nuclear Physics A 00 (2013) 1-4

Three-body couplings in R...
and its effects on hyperonic star equatio...
Tsubakihara³, A. Ohnishi^b
^aLaboratory, Hokkaido University,
^bTheoretical Physics, Eyon University

Neutron star mass and the resolution of hyperon puzzle in modified gravity
Artyom V. Astashenok¹, Salvatore Capozziello^{2,3,4}, Sergei D. Odintsov^{5,6,7}
¹Kant Baltic Federal University, Institute of Physics and Technology, Nevel'skoje st. 14, 236041 Kaliningrad, Russia
²Dipartimento di Fisica, Università di Napoli "Federico II" and INFN Sez. di Napoli, Via Cintia, 9, I-80126, Napoli, Italy
³Istituto Nazionale di Fisica Nucleare (INFN), Viale F. Crispi, 7, I-67100, L'Aquila, Italy
⁴Compl. Univ. di Monte S. Angelo, Ed. G., Via Cintia, 9, I-80126, Napoli, Italy
⁵Grav. Sasso Science Institute (ISI), Via Sommarive, 14, I-38127, Trento, Italy
⁶Institució Catalana de Recerca i Estudis Avançats (ICREA), Torre C5-Par-2a pl. E-08193 Bellaterra, Barcelona, Spain
⁷Tomsk State Pedagogical University, National Research Tomsk State University, 634002 Tomsk, Russia

The influence of Strong Magnetic Field in Hyperonic Neutron Stars
Debora Peres Menezes
Santa Catarina

Hyperon mixing and universal many-body repulsion in neutron stars

Y. Yamamoto¹, T. Furumoto², N. Yasutake³, and Th.A. Rijken^{4,1}

¹Nishina Center for Accelerator-Based Science, Institute for Physics and Chemical Research (RIKEN), Wako, Saitama 351-0192, Japan

²National Institute of Technology, Ichinoseki College, Ichinoseki, Ibaraki 310-8501, Japan

³Department of Physics, Chiba Institute of Technology, 2-1-1 Shibazono Narashino, Chiba 275-8585, Japan

⁴IMAPP, University of Nijmegen, Nijmegen, The Netherlands

A multi-pomeron exchange potential (M_{pp}) and many-body repulsion in baryonic systems. The strength of M_{pp} is determined by G-matrix folding model. A binding energy calculation including the effect of M_{pp} is performed. PHYSICAL REVIEW C 89, 025803 (2014)

Effects of fermionic dark matter on properties of neutron stars
Jiang,^{1,2} Wei-Zhou Jiang,^{1,3,4,*} Dong-Rui Zhang,¹ and Rong-Yao Yang¹
¹Department of Physics, Southeast University, Nanjing 211189, China
²Department of Physics, Chinese Academy of Sciences, Beijing 100049, China
³Heavy Ion Accelerator, Lanzhou 730000, China
⁴Department of Physics, University of California, Santa Cruz, California 95064, USA
Received 16 November 2013; published 13 February 2014

Can very compact and very massive neutron stars both exist?
Alessandro Drago,¹ Andrea Lavagno,² and Giuseppe Pagliara¹
¹Dipartimento di Fisica e Scienze della Terra dell'Università di Ferrara and INFN Sezione di Ferrara, Via Saragat 1, I-44100 Ferrara, Italy
²Department of Science and Technology, Politecnico di Torino, I-10129 Torino, Italy
Received 13 October 2013; published 25 February 2014

Hyperon equations of state for supernovae and in density dependent hadron field theory
Sarmistha Banik
Department of Physics, University of Hyderabad, Hyderabad Campus, Hyderabad, India
Matthias Heiser
Astropartikelfysik, Universität Wien, Vienna, Austria

What does a measurement of mass and/or radius of a neutron star constrain? Equation of state or gravity?

Kazım Yavuz Ekşi,^{*} Can Güngör, and Murat Metehan Türkoğ
Istanbul Technical University, Faculty of Science and Letters, Department of Physics, 34469 Maslak, Istanbul, Turkey
(Received 29 November 2013; published 6 March 2014)

Eddington-inspired Born-Infeld gravity: Phenomenology of nonlinear gravity-matter coupling
Paolo Pani,¹ TERENCE DELSATE,^{1,2} and Vitor Cardoso^{1,3}
¹CENTRA, Departamento de Física, Instituto Superior Técnico, Universidade Técnica de Lisboa—UTL, Av. Rovisco Pais 1, 1049 Lisboa, Portugal
²Theoretical and Mathematical Physics Dept., Université de Mons, UMONS, 20, Place du Parc 7000 Mons, Belgium
³Department of Physics and Astronomy, The University of Mississippi, University, Mississippi 38677, USA
Received 17 January 2012; published 19 April 2012

of Nuclear Physics
India
Switzerland

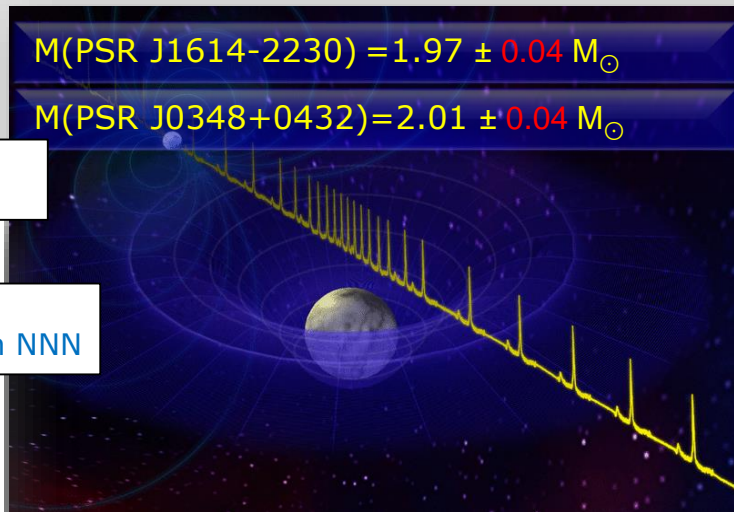
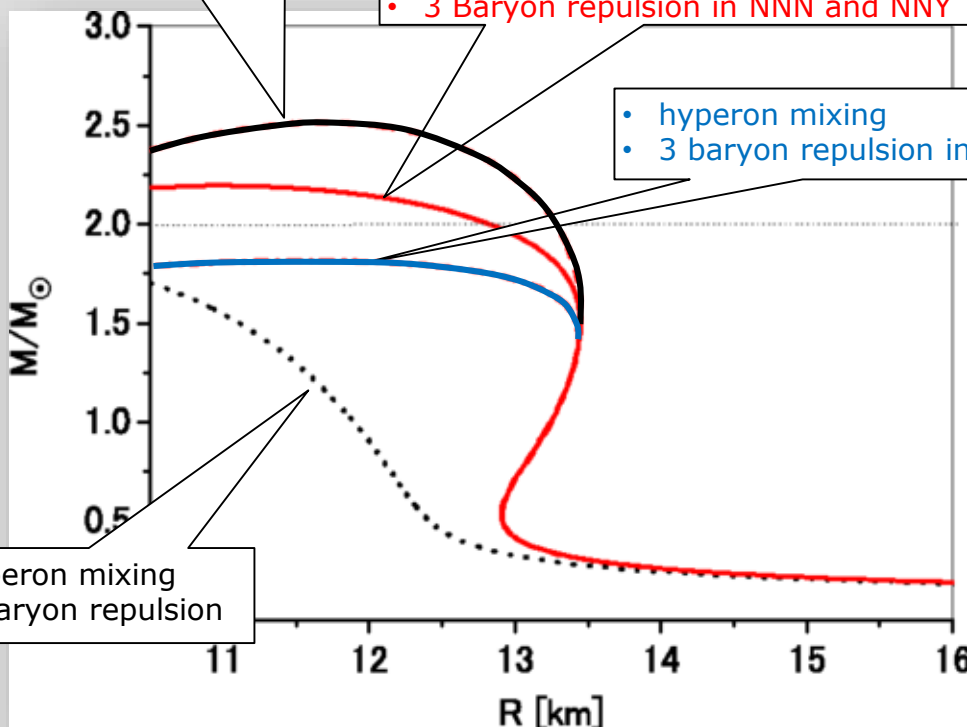
Y. Yamamoto, T. Furumoto, N. Yasutake, Th. A Rijken,
Phys. Rev. C 90, 045805 (2014)

- no hyperon mixing
- 3 Baryon repulsion

- hyperon mixing
- 3 Baryon repulsion in NNN and NNY

- hyperon mixing
- 3 baryon repulsion in NNN

- no hyperon mixing
- no 3 baryon repulsion



$M(\text{PSR J1614-2230}) = 1.97 \pm 0.04 M_{\odot}$
 $M(\text{PSR J0348+0432}) = 2.01 \pm 0.04 M_{\odot}$

- model constrained by terrestrial experiments
- universal many-body repulsion
- no ad hoc parameter to stiffen EOS

Yamamoto (HYP2015):

"Including 3- and 4-body repulsions leads to massive neutron stars with $2M_{\odot}$ in spite of significant softening of EOS by hyperon mixing"....

"Hyperon puzzle is a quantitative problem"



Tools

- Heavy Ion beams
- Electron beams
- Photon beams
- Meson beams
- Antiproton beams

Methods

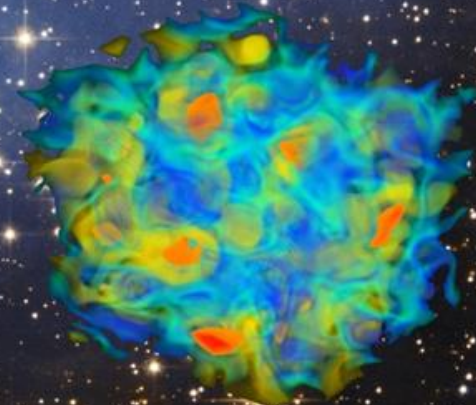
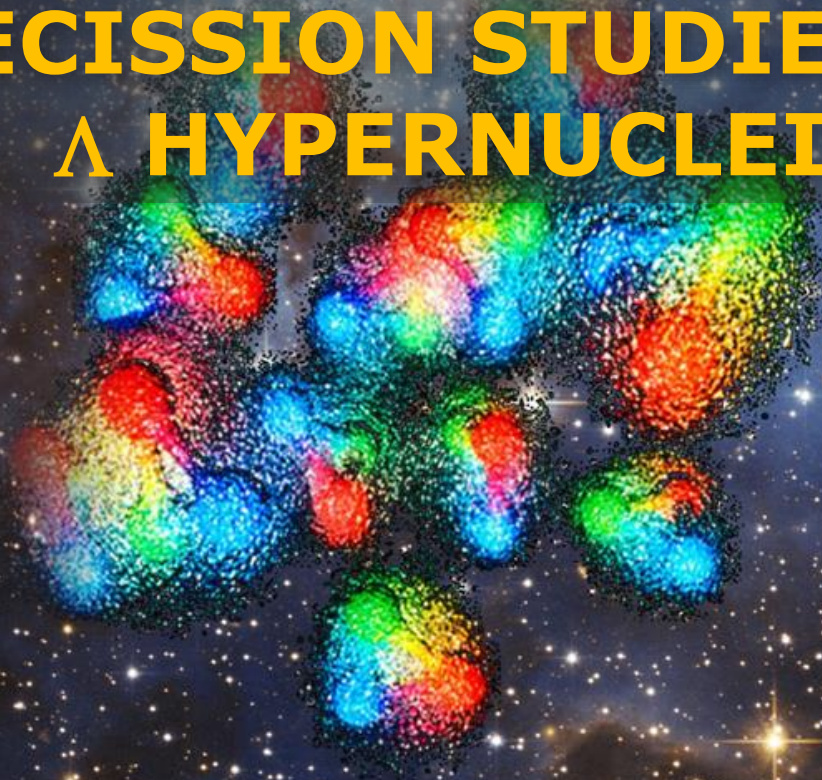
- Invariant mass studies
- Missing mass studies
- γ -spectroscopy
- π -spectroscopy

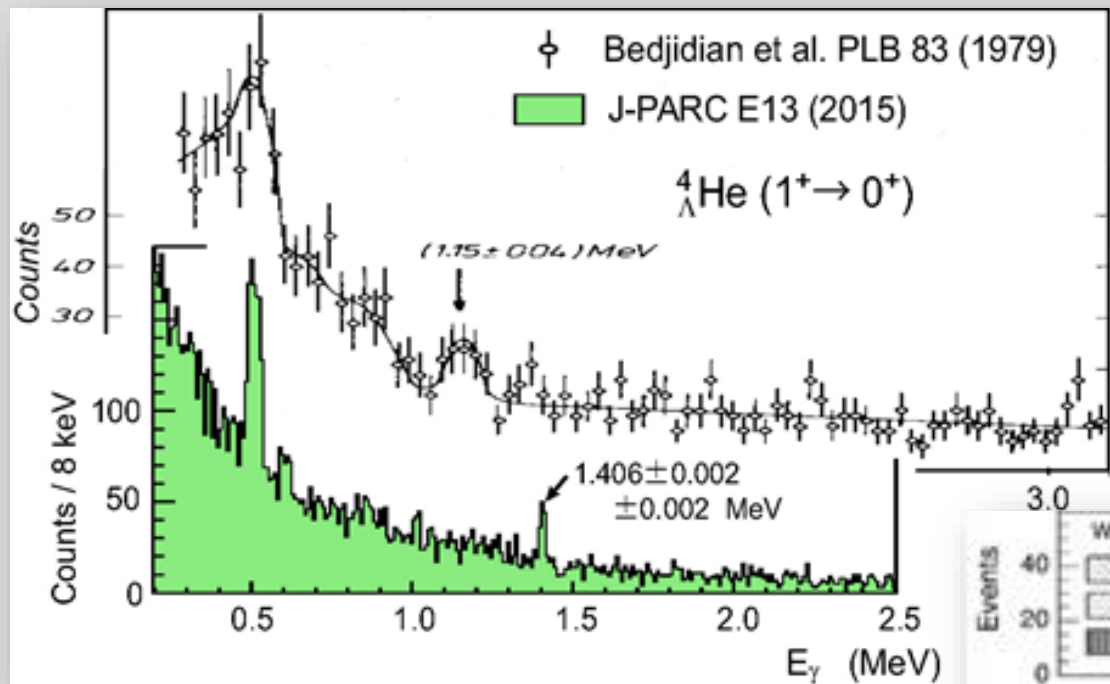
Observables

- Masses
- Excitation spectrum
- Lifetimes

PART 1

PRECISION STUDIES OF Λ HYPERNUCLEI

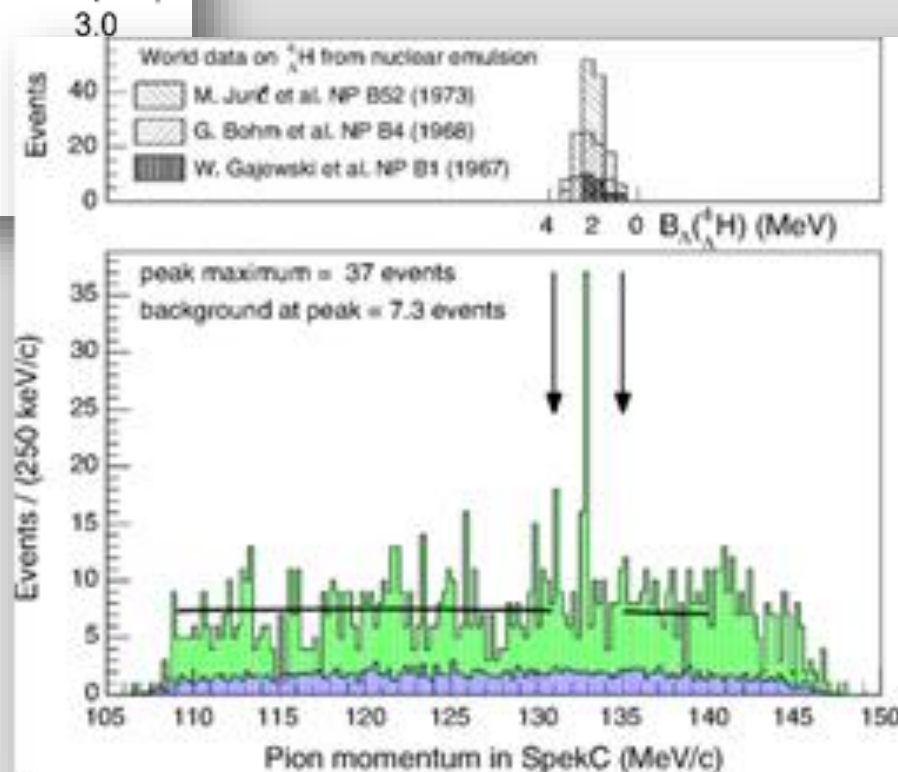




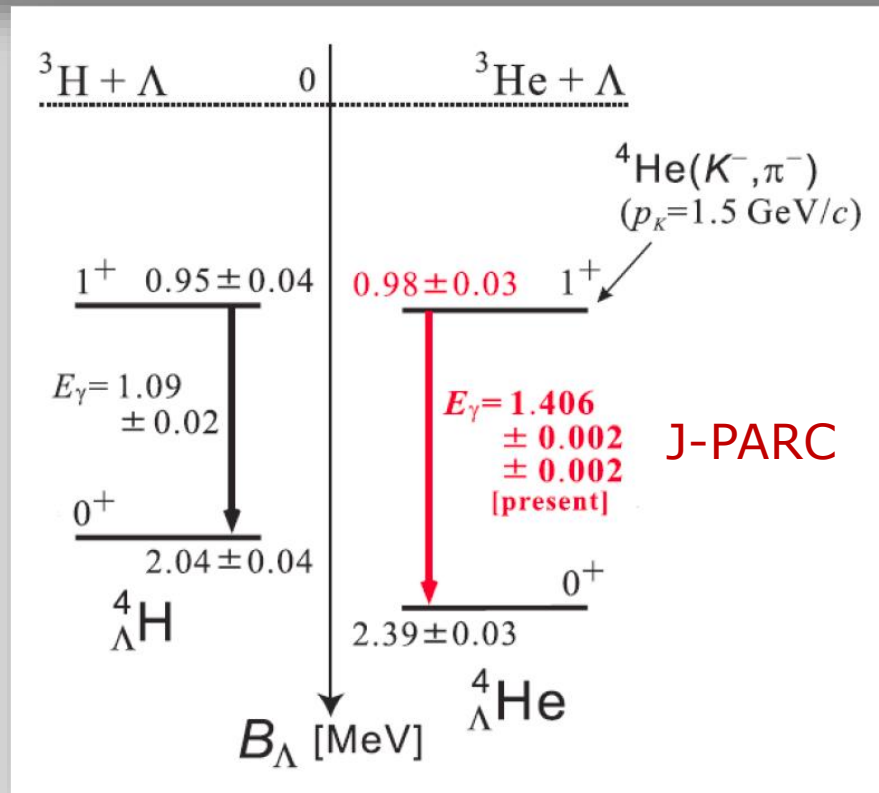
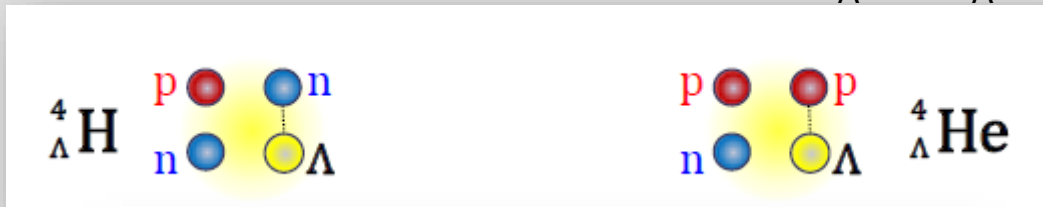
Phys. Rev. Lett. **114**,
232501 (2015)

Phys. Rev. Lett. **115**,
222501 (2015)

- Demonstrates the need for complementary experiments



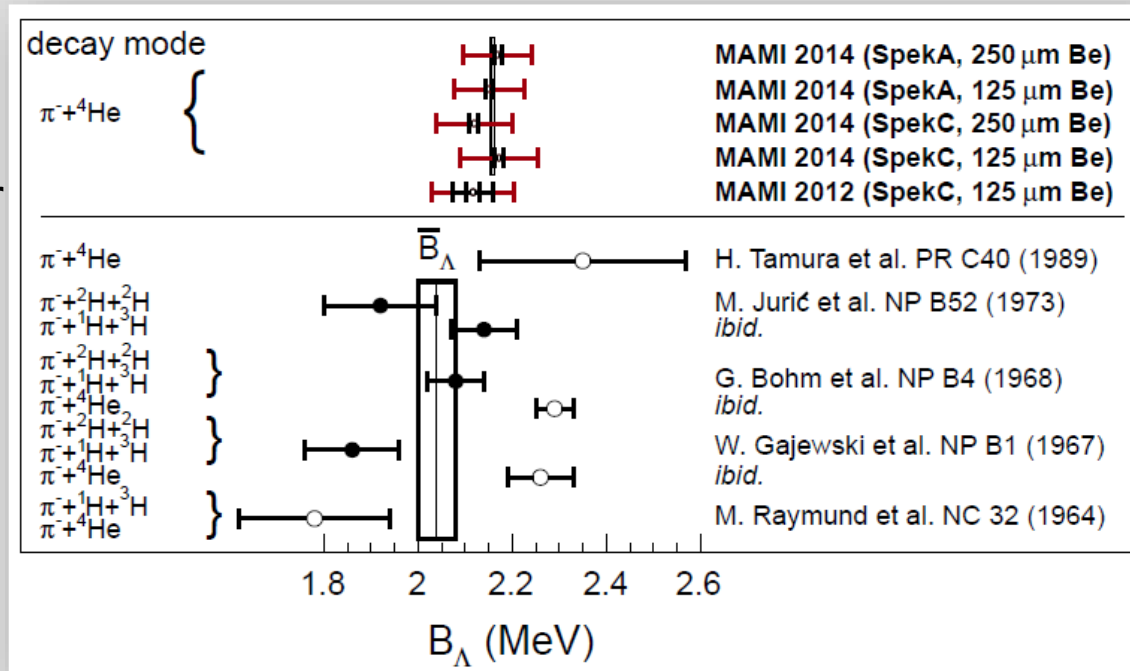
- ▶ CSB for NN interactions is 70 keV in the mirror nuclei ${}^3\text{H}$ and ${}^3\text{He}$
- ▶ Coulomb corrections are < 50 keV for the ${}^4_{\Lambda}\text{H} - {}^4_{\Lambda}\text{He}$ pair



- ▶ strong, **spin-dependent** charge symmetry breaking (CSB) in $A = 4$ mirror hypernuclei !

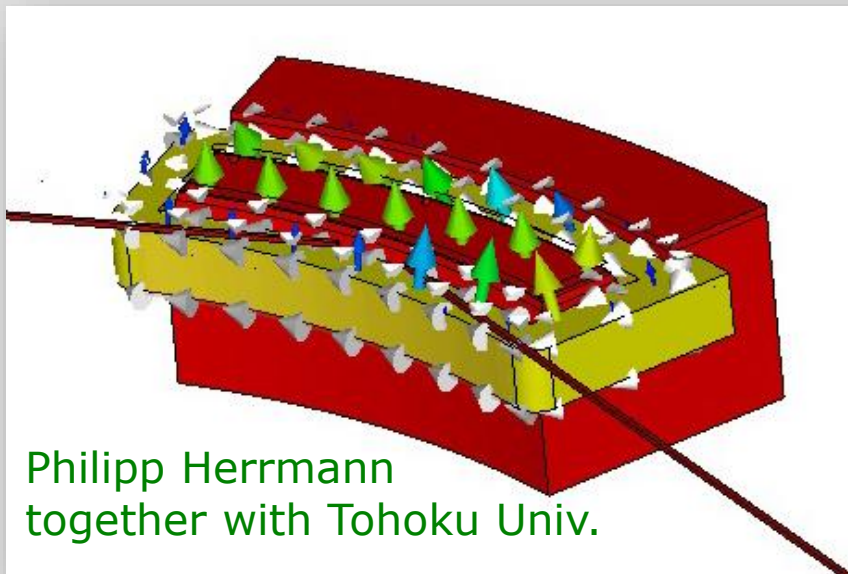
► Many improvements

- better pion rejection by improved aerogel
- suppression of background by improved shielding
- suppression of background by trigger upgrade in SpekA & C
- suppression of background by beam-line upgrade
- dedicated collimator for decay region
- better control of magnet field variations
- full overlap of SpekA and SpekC momentum acceptance

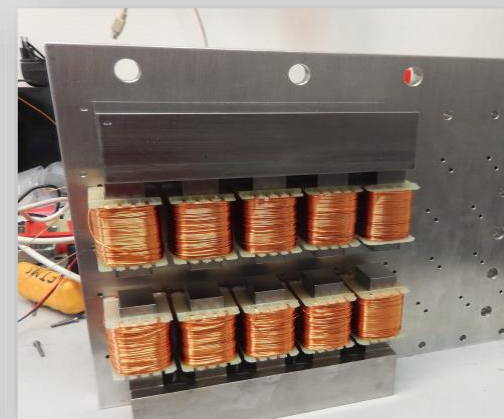
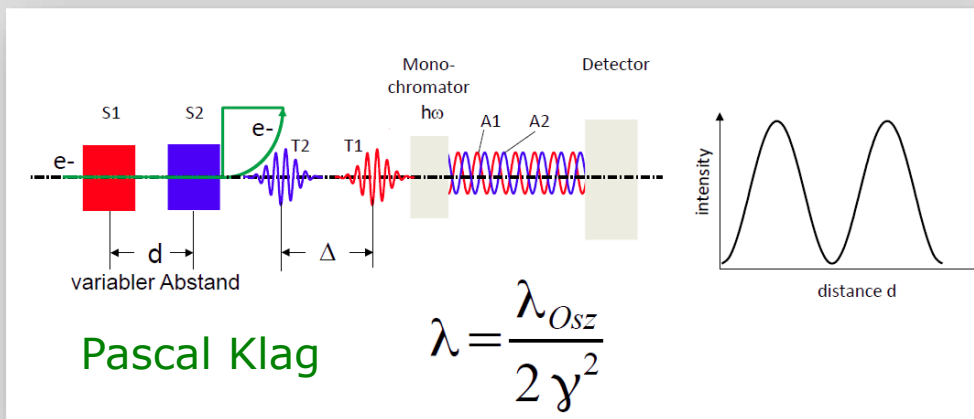


- independent measurement in two spectr., two targets, two beam-times
- consistent result for $B_\Lambda({}^4_\Lambda\text{H})$ from MAMI 2012 and MAMI 2014

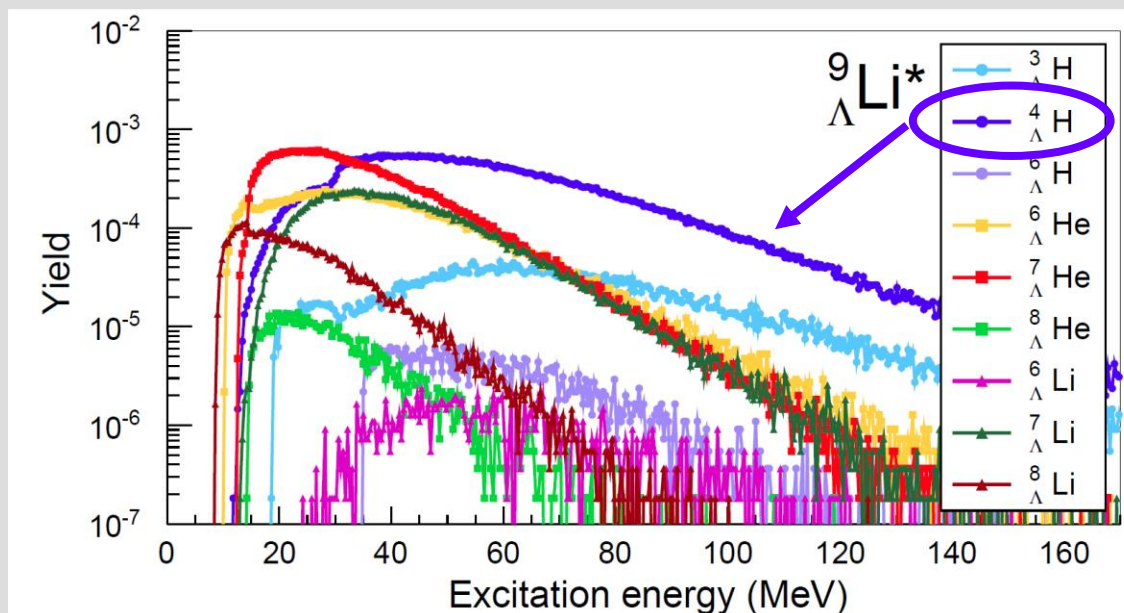
- ▶ two options to measure MAMI energy $O(10^{-4})$
 - ▶ measuring the absolute MAMI energy in a precisely calibrated magnet



- ▶ Interference of undulator radiation

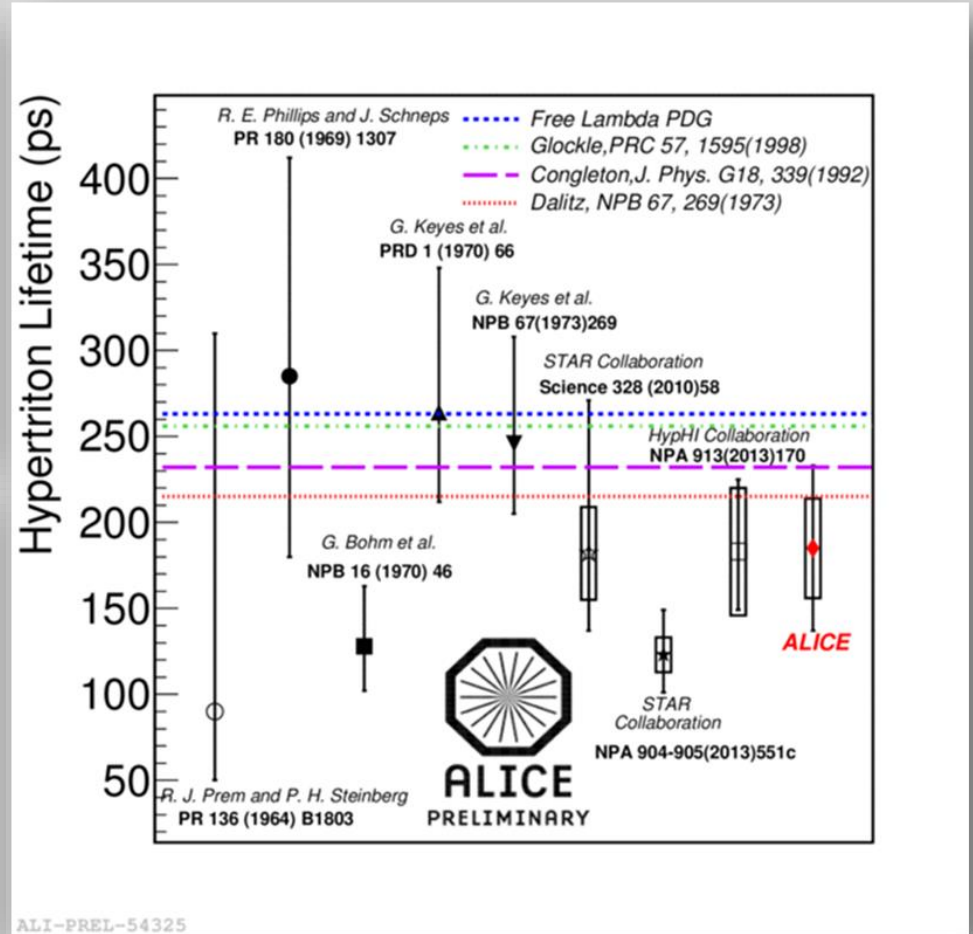
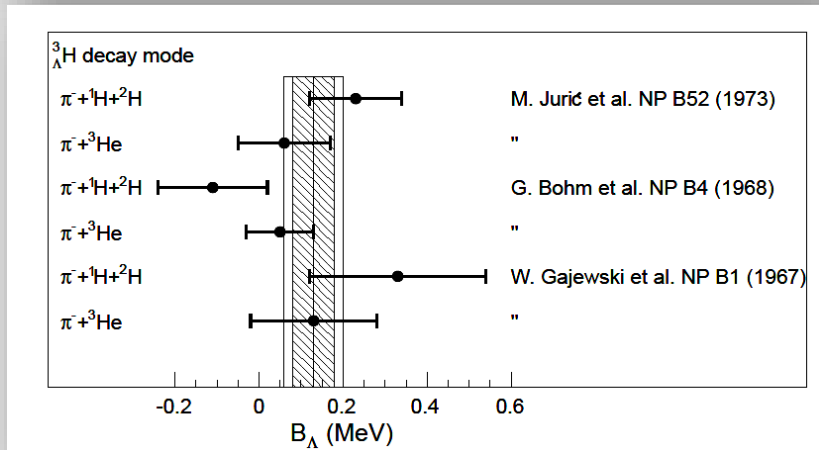
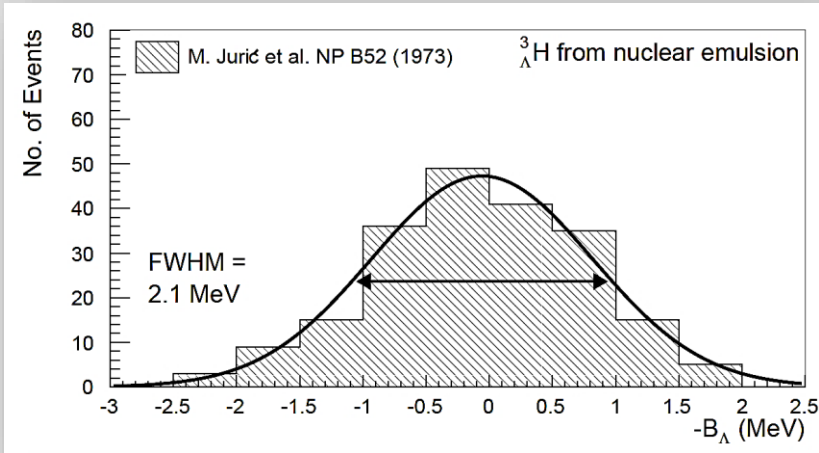


- ▶ Statistical decay calculations were performed
- ▶ **Scenario 1**: direct production of ${}^9_{\Lambda}\text{Li}^*$



- ▶ Expected excitation energy
 - ▶ convert proton into $\Lambda \Rightarrow$ proton hole state ~ 20 MeV
 - ▶ kinetic energy of captured Λ $p_{\text{FERMI}}^2/2M_{\Lambda}$ ~ 20 MeV
 - ▶ Binding energy of Λ ~ 10 MeV
- ▶ at $E_x \sim 50$ MeV ${}^4_{\Lambda}\text{H}$ most probable and other nuclei more than factor 3 less likely produced
- ▶ similar for nucleon knock-out scenario

► Small B_{Λ} (from about 200 analyzed events from emulsion)

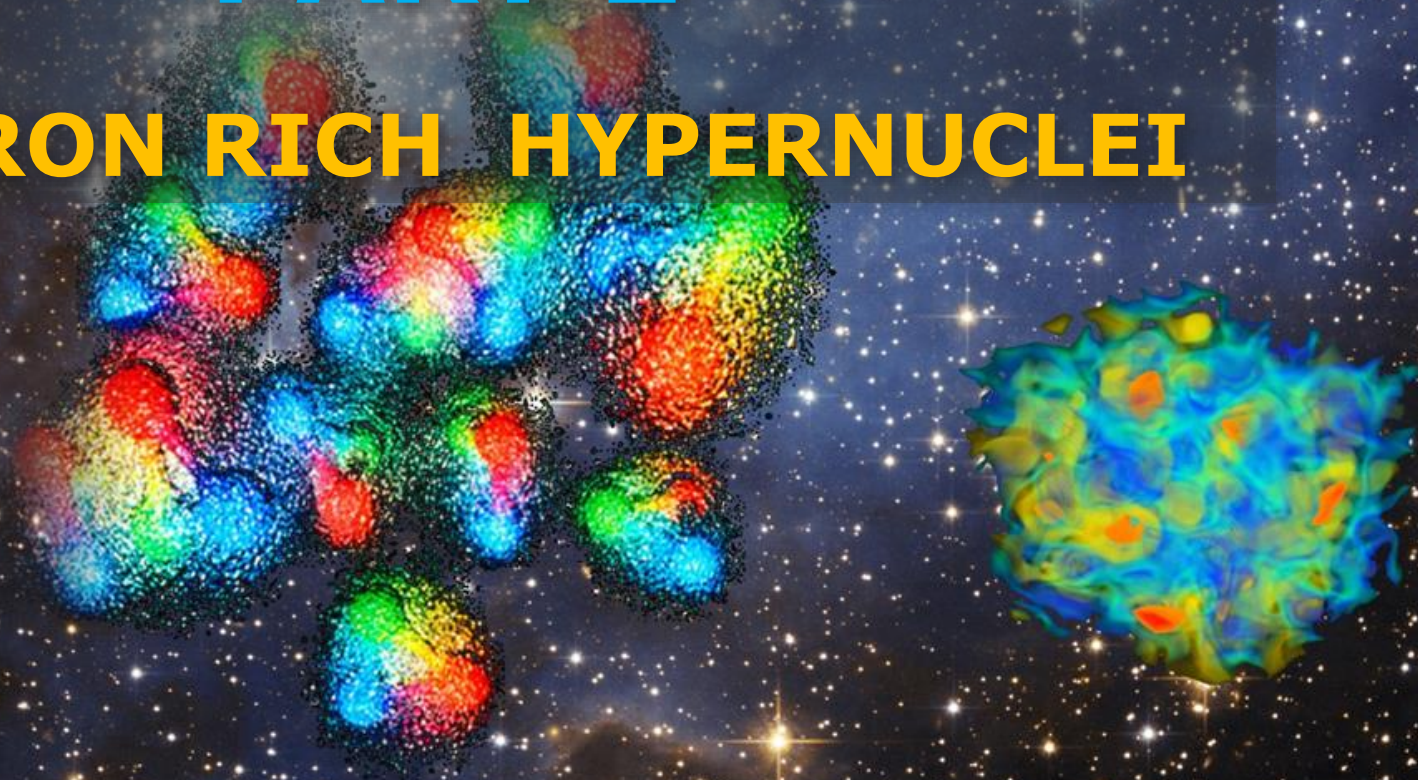


Achenbach, Pochodzalla, Schulz (PANIC 2014)

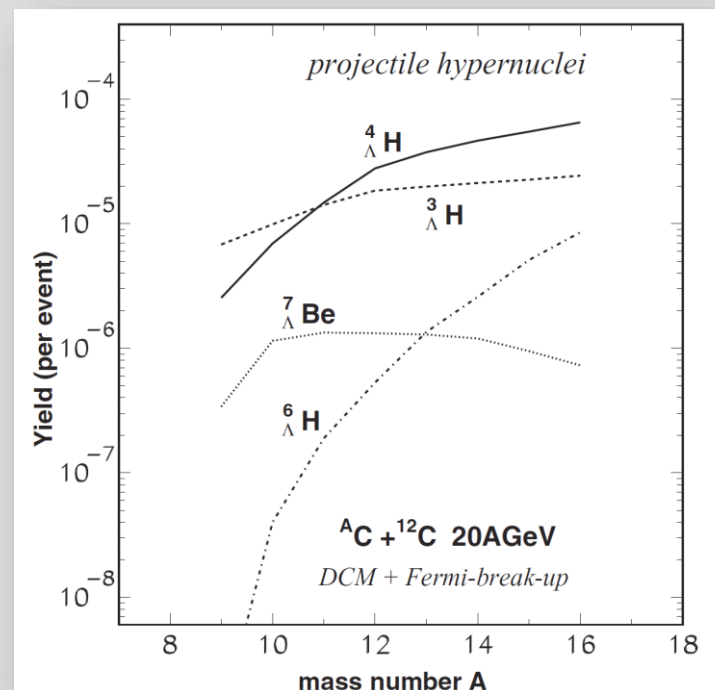
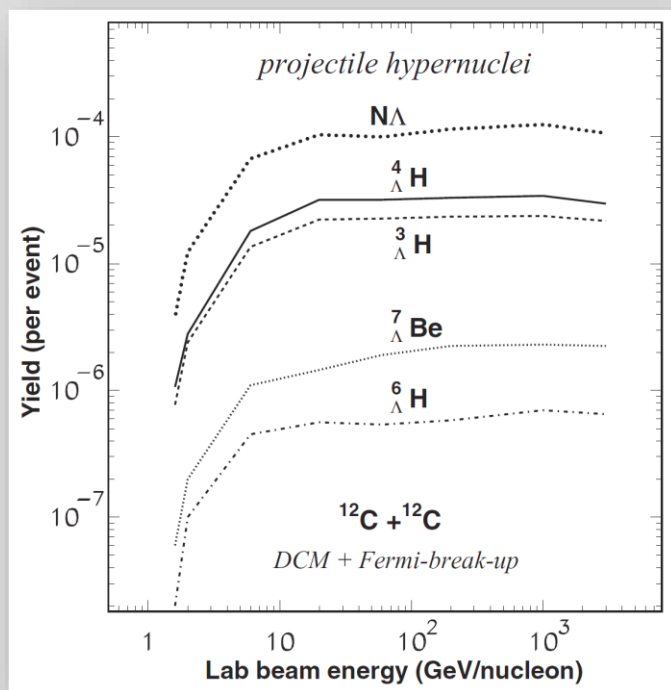
...and Lifetime surprisingly small

PART 2

NEUTRON RICH HYPERNUCLEI



- ▶ Light nuclei
 - ▶ meson induced reactions → spectroscopy
 - ▶ projectile fragmentation like
HYPHI, SuperFRS with exotic beams → lifetime



Botvina, Gudima, Pochodzalla PRC 88 (2013)

- ▶ Heavy neutron rich nuclei using exotic targets
 - ▶ electron beams and ${}^{40}\text{Ca}$, ${}^{48}\text{Ca}$ → spectroscopy
 - ▶ fragmentation- fission → lifetime

- ▶ First experiment done at J-Lab (L. Tang); analysis not finalized
- ▶ Observed: two fragment with $Z \geq 6$
- ▶ **decay time relative to beam** nano time structure

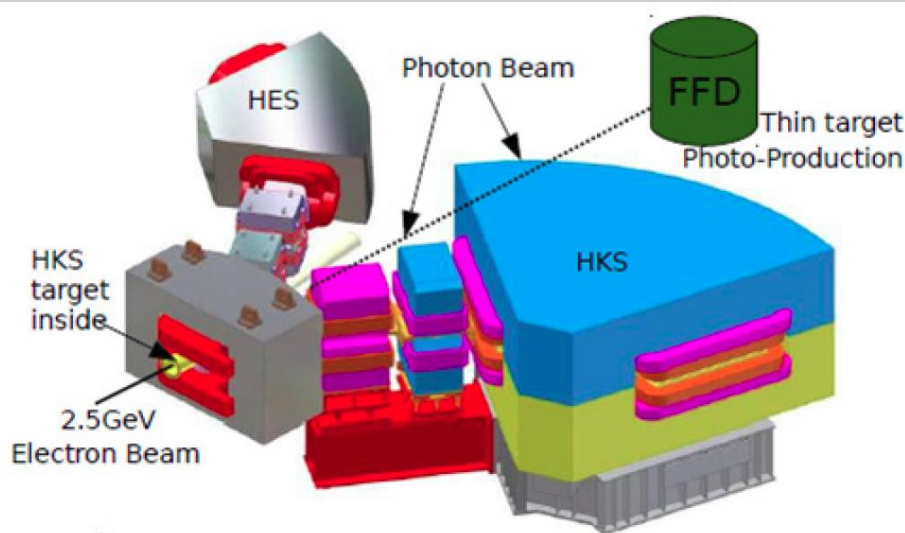


FIG. 2. The experimental setup of E05-115 and E02-017 in Hall C at JLab.

- ▶ Targets Fe, Cu, Ag, Bi, (Au, U) simultaneously
 ⇒ some systematic errors cancel

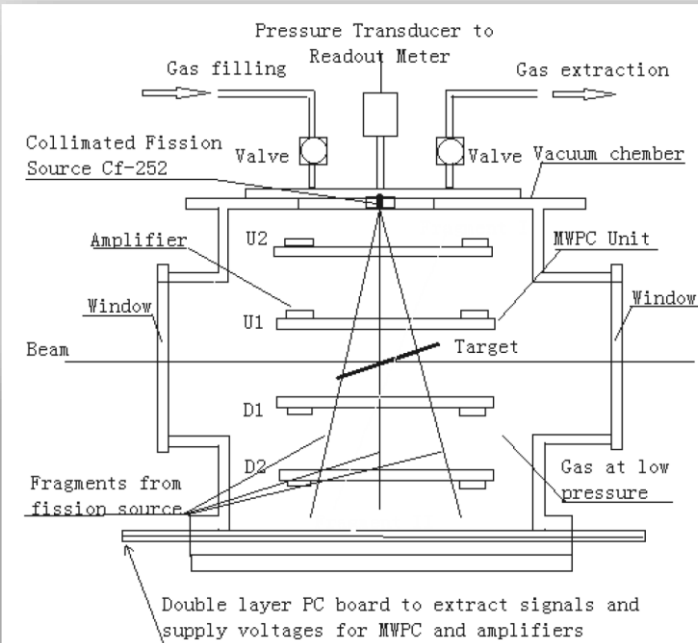
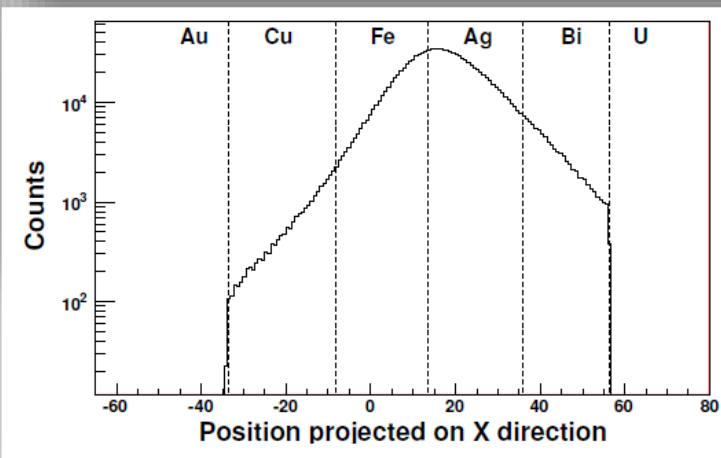
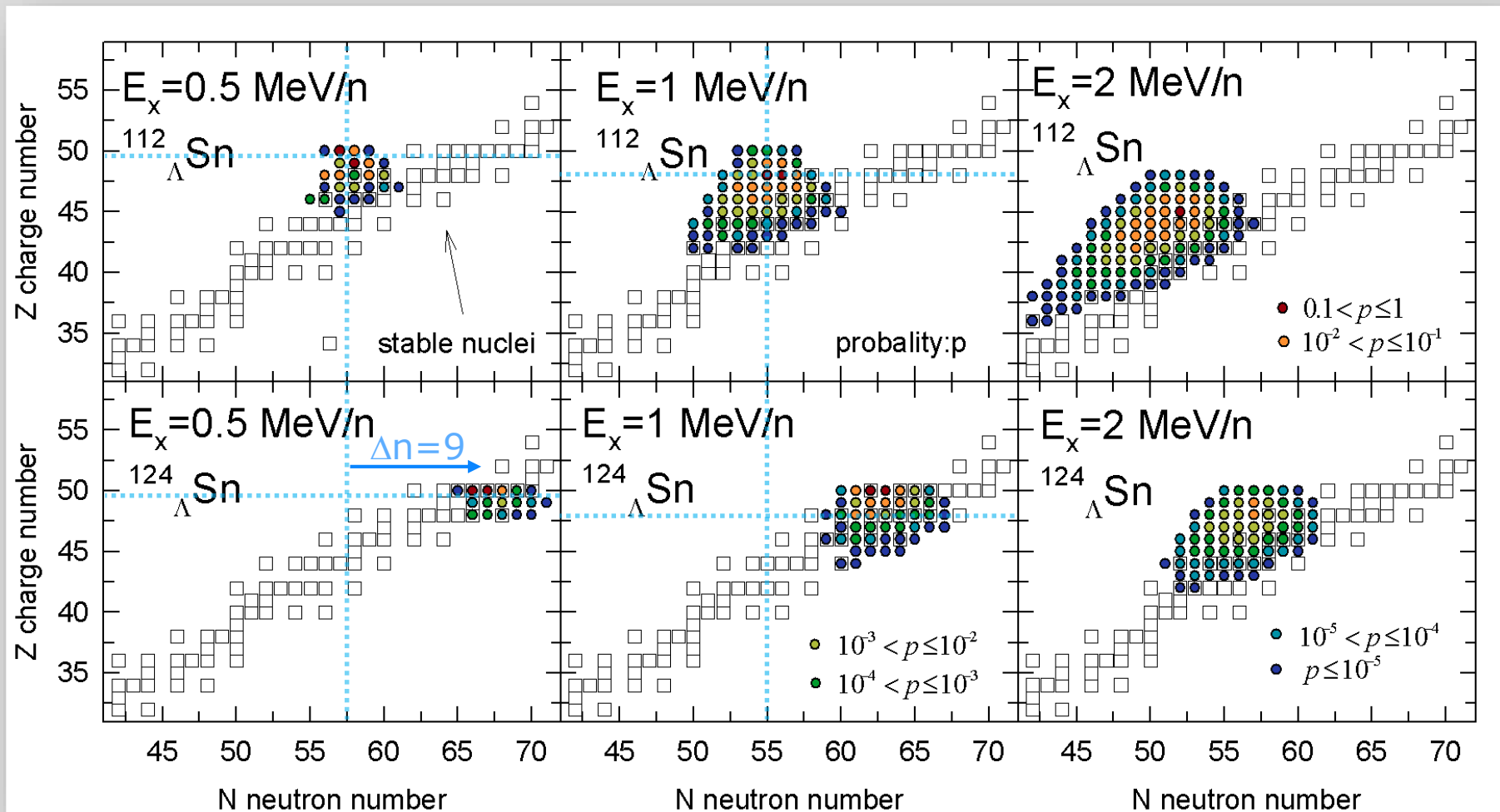
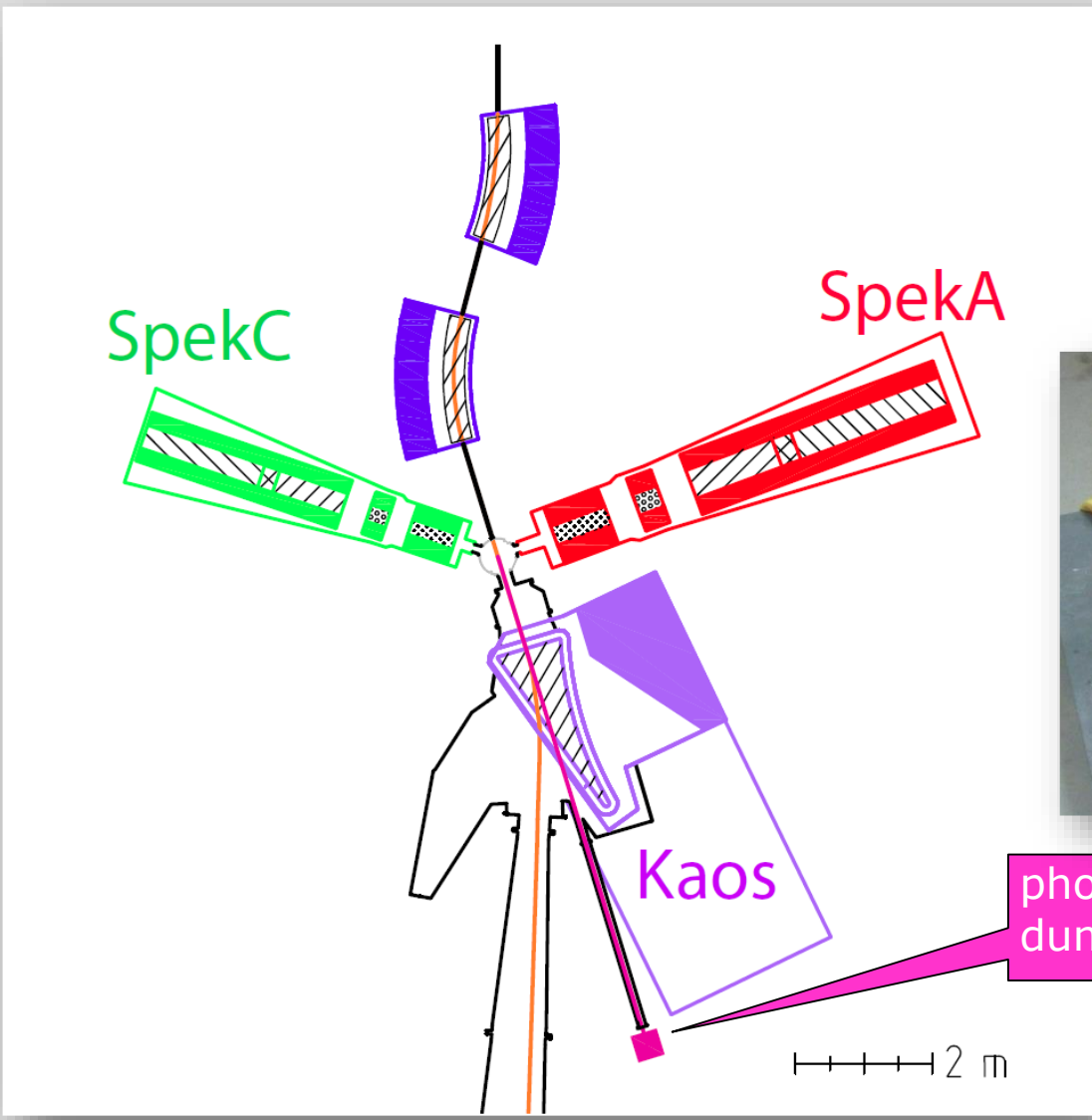


FIG. 1. Schematic sketch of the fission fragment detector (FFD) [20].



- ▶ expected excitation energy as before $\sim 1\text{MeV/nucleon}$
- ▶ produced hyperresidues reflect N/Z of target



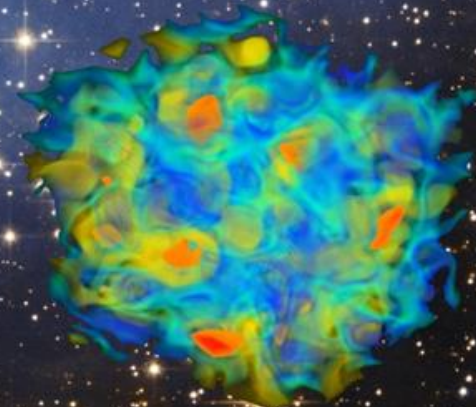


e⁻-beam

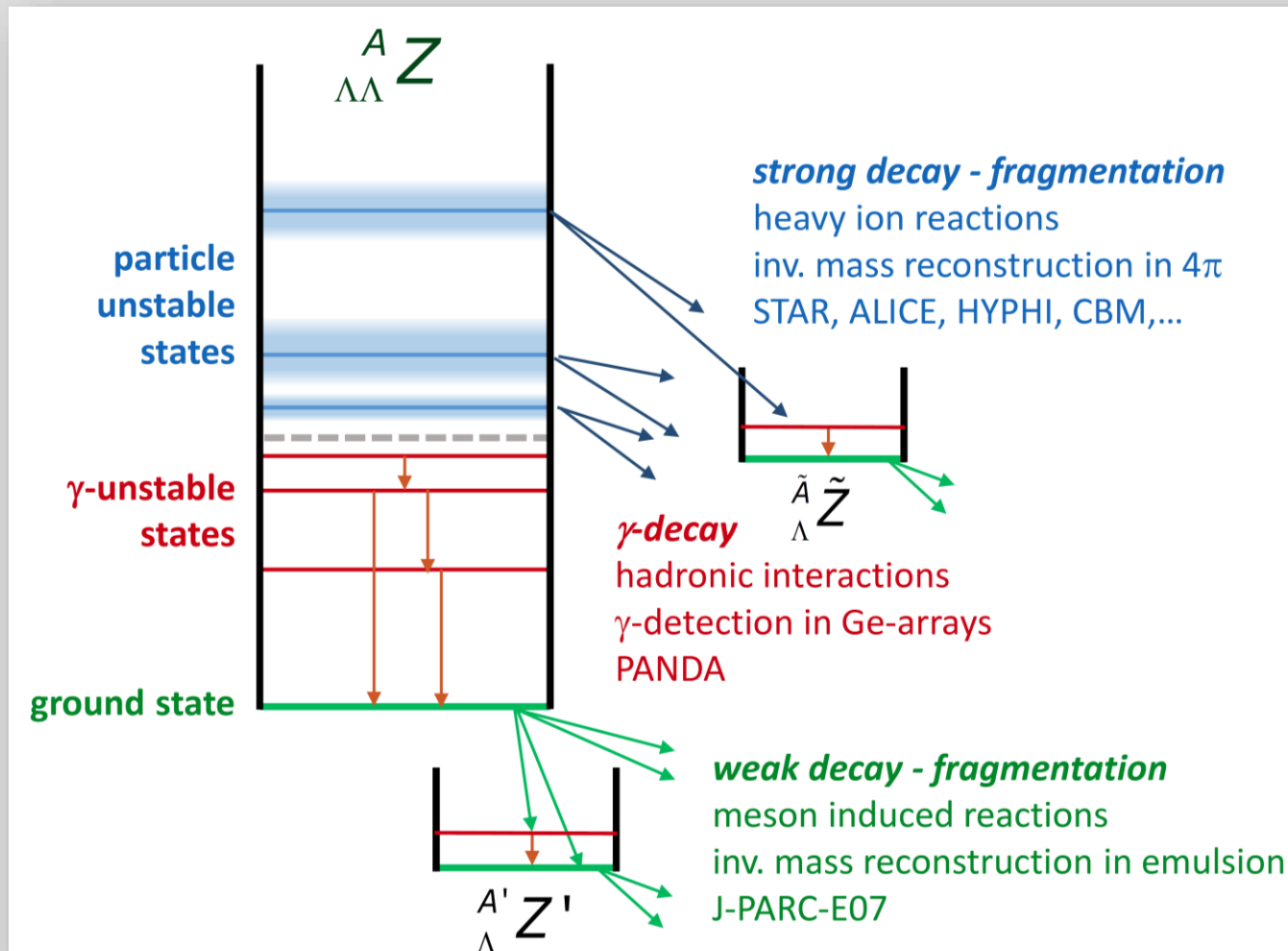
photon beam dump

PART 3

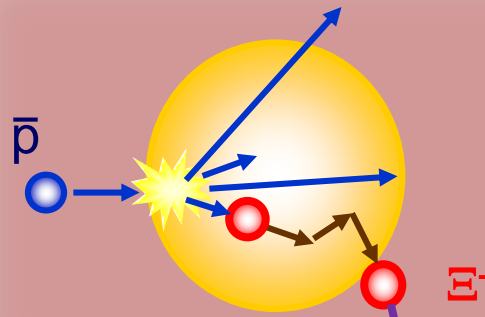
S=-2 SYSTEMS



- ▶ missing mass (K^-, K^+) reactions \Rightarrow Ξ bound state J-PARC
- ▶ Ξ capture \Rightarrow Ξ atoms J-PARC, FAIR
- ▶ Ξ capture and $\Xi^- p \rightarrow \Lambda \Lambda$ \Rightarrow $\Lambda \Lambda$ hypernuclei J-PARC, FAIR, HI



Ξ^- production
 $\bar{p}N \rightarrow \Xi^- \bar{\Xi}$



rescattering in
primary target nucleus

deceleration in
secondary target

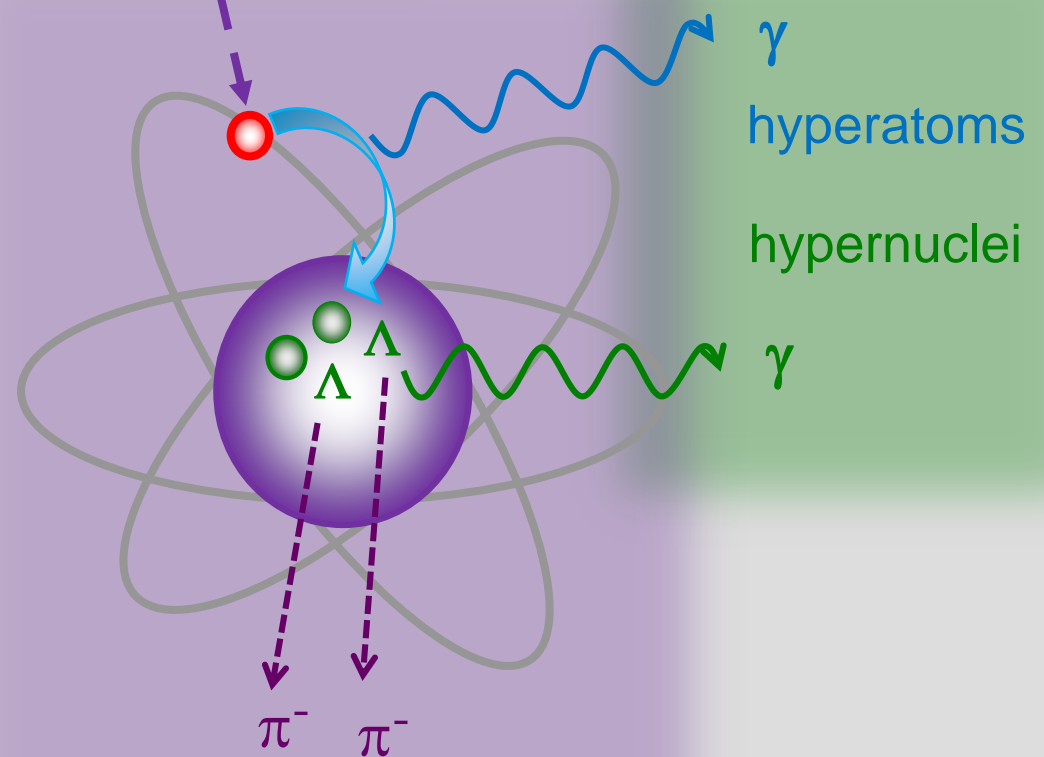
capture of Ξ

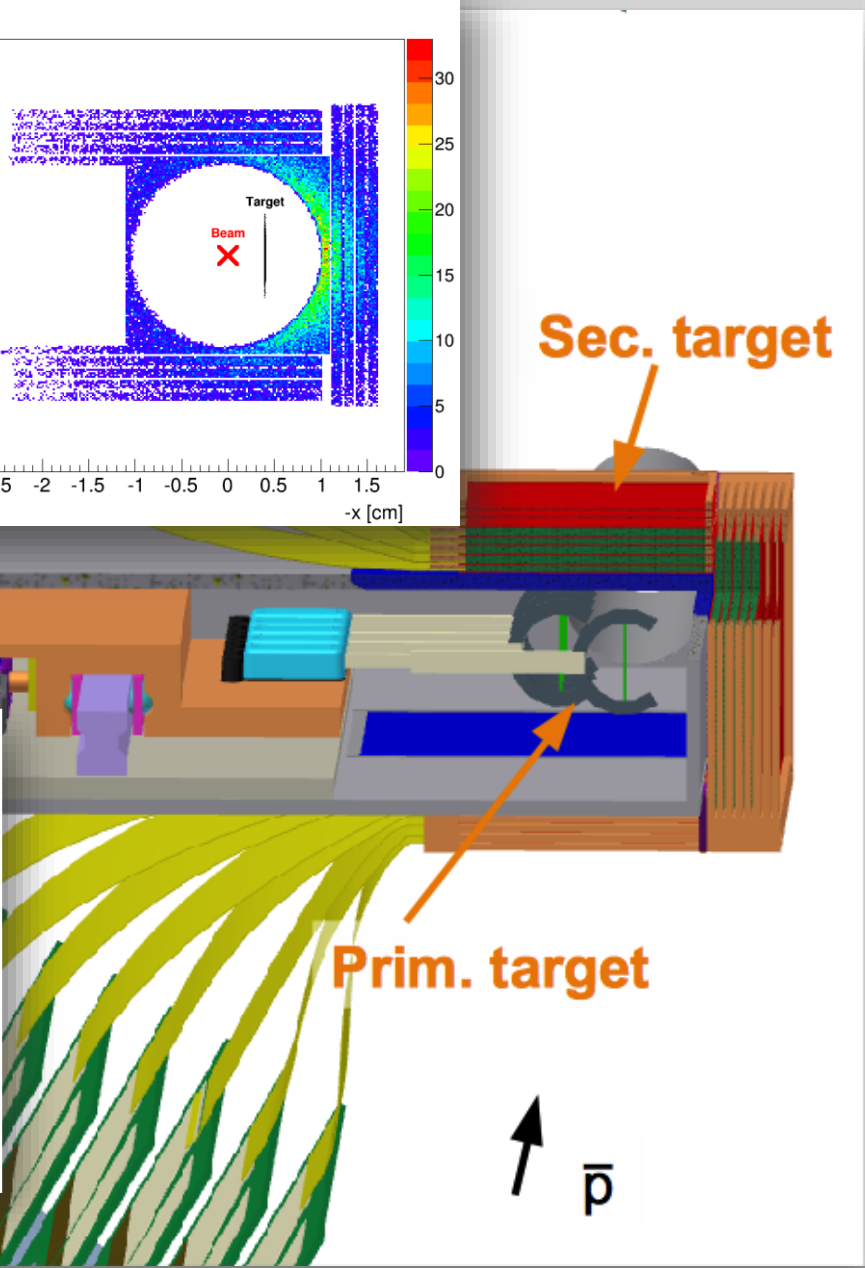
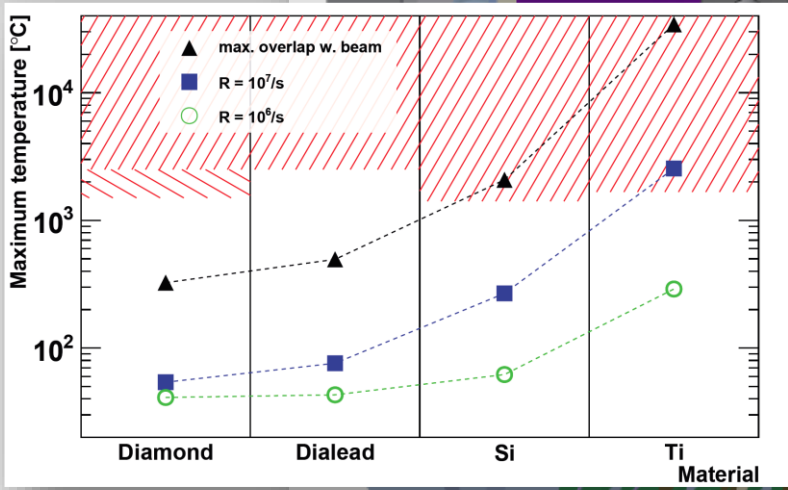
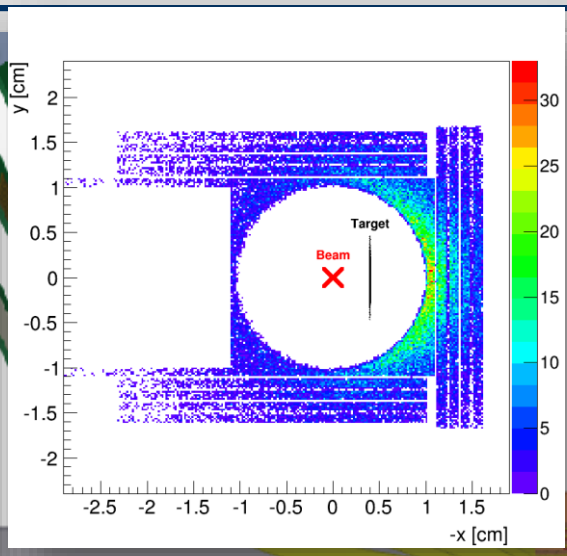
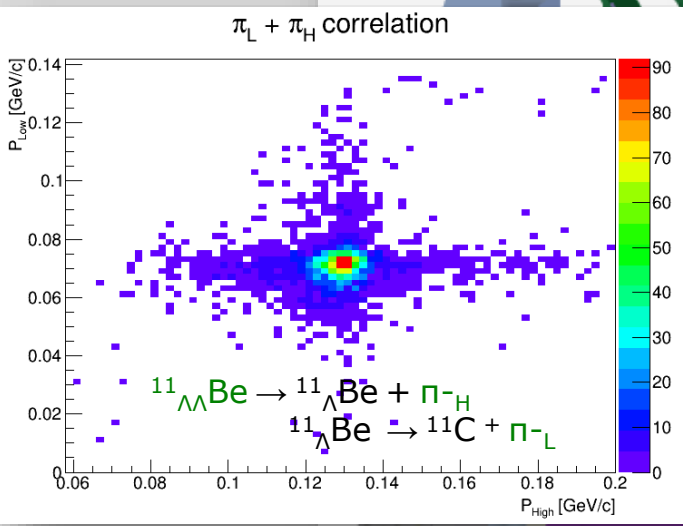
atomic cascade of Ξ^-

$\Xi^- p \rightarrow \Lambda\Lambda$ conversion
fragmentation
 \rightarrow excited $\Lambda\Lambda$ -nucleus

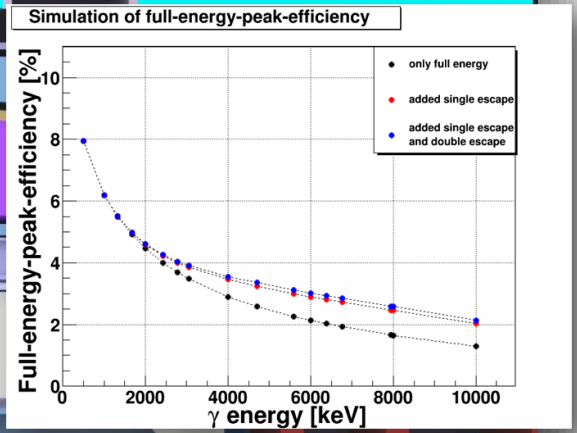
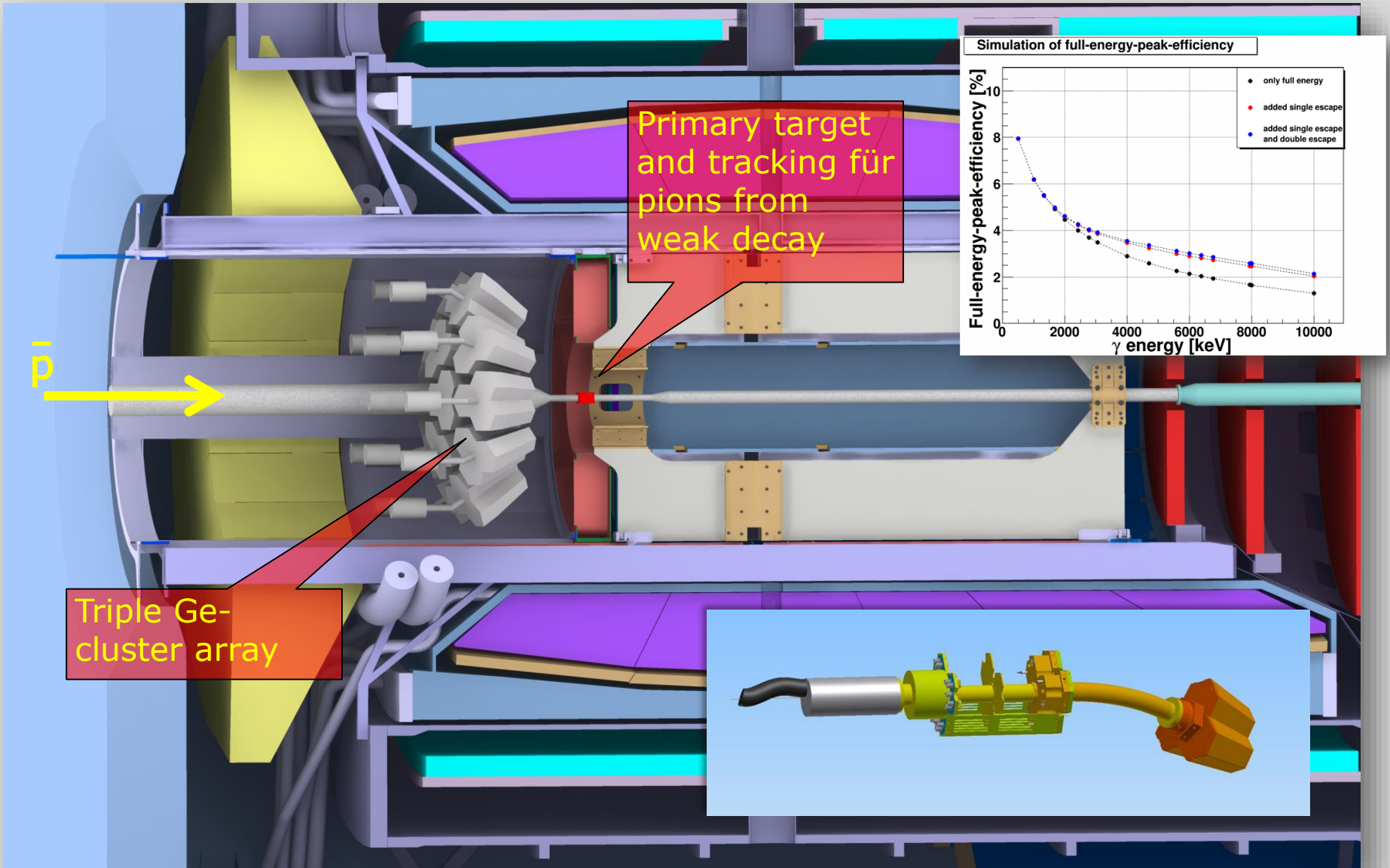
γ -decay of $\Lambda\Lambda$ hypernuclei

weak pionic decay



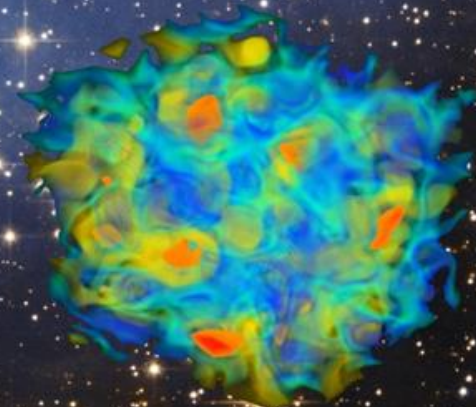


The HYP setup at PANDA



PART 4

S=-3 SYSTEMS



PHYSICAL REVIEW D

VOLUME 8, NUMBER 3

1 AUGUST 1973

Certification of Three Old Cosmic-Ray Emulsion Events as Ω^- Decays and Interactions

Luis W. Alvarez

Lawrence Berkeley Laboratory, University of California, Berkeley, California 94720

(Received 10 April 1972; revised manuscript received 3 May 1973)

In the "pre-accelerator years," when large stacks of emulsion were exposed to cosmic rays at high altitude, three events were found in which K^- mesons were emitted from slowly moving particles. The Ω^- is the only presently known particle that can give rise to a K^- when moving at nonrelativistic speed, but none of the three events has until now been clearly identified as an Ω^- . One of the cosmic-ray events (Eisenberg, 1954) has been incorrectly interpreted as an Ω^- decaying in flight; it is now shown to be an interaction in flight of an Ω^- with a silver nucleus. The second event is a clear-cut example of an Ω^- decaying in orbit, bound to an emulsion nucleus. The third event is quite complicated, but can be unambiguously attributed to the decay of an Ω^- atomically bound to an N^{14} nucleus, followed by a collision of the daughter Λ with the N^{14} , in which the compound system then fragments into ${}_{\Lambda}C^{13} + p + n$. The mass of the Ω^- as determined by each of the last two events (Fry *et al.*, 1955) agrees closely with the mean of all bubble-chamber events.

- Note: in nuclei secondary processes possible



...seen in emulsions ~10 years prior to the „discovery“ at Brookhaven

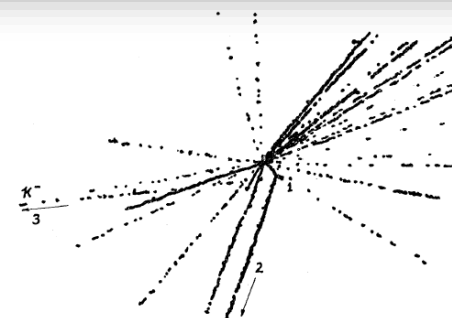
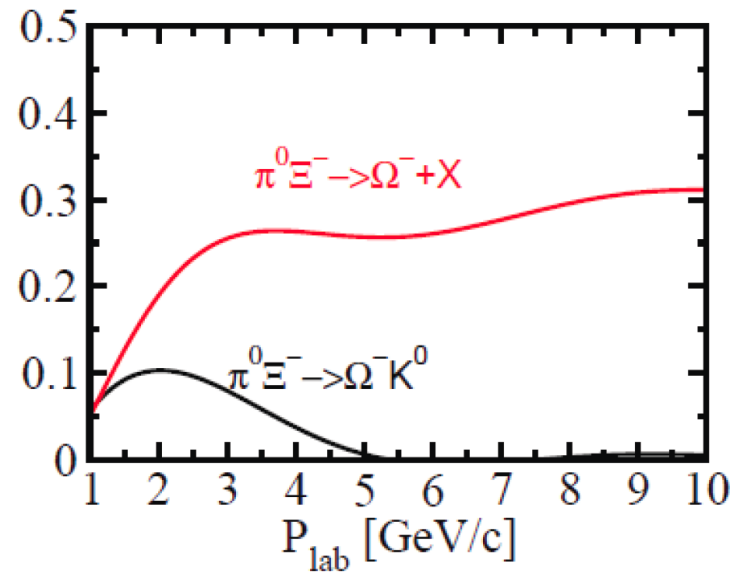
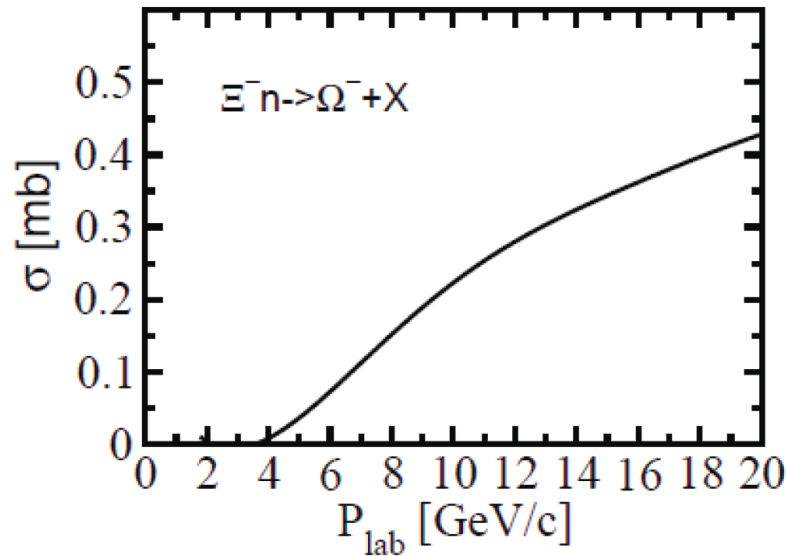
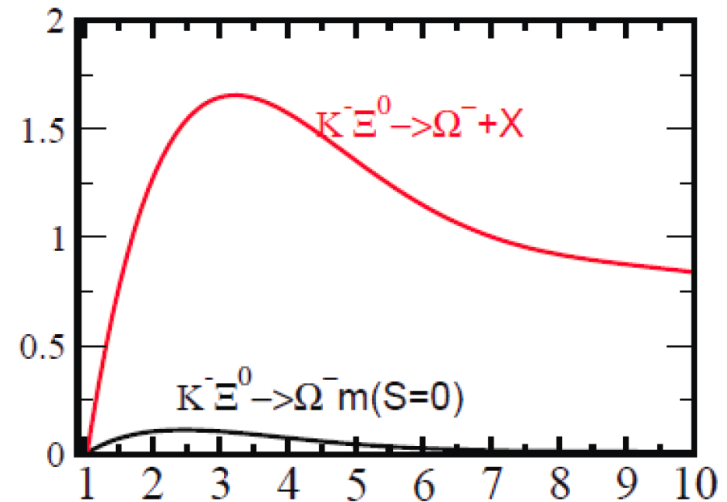
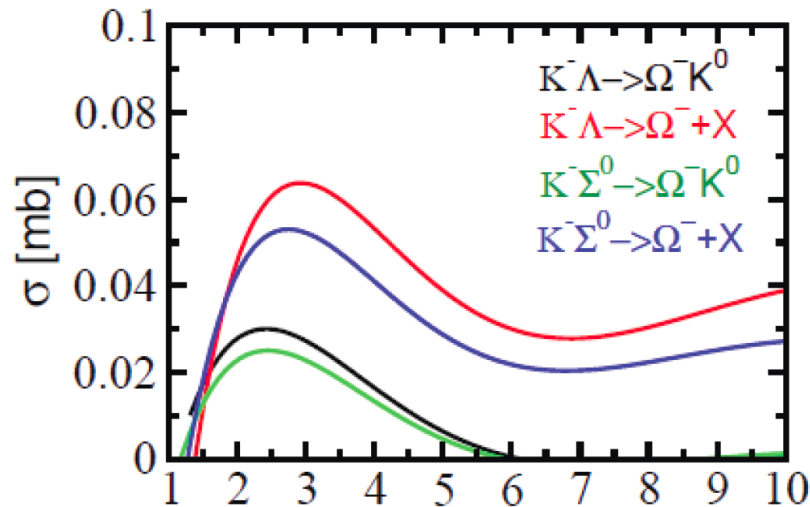


Fig. 1. A projection drawing of the K -mesonic decay of a slow particle is shown above. Track 1 is a short recoil. Track 2 was produced by a particle of $Z=1$. Track 3 was produced by a negative K -meson. A few tracks of particles from the primary star which are in the same direction as the connecting track, but at a different depth, were omitted from the drawing for the sake of clarity.

...in-medium Ω^- production at J-PARC ?



D. Aston *et al.*, Phys. Lett. B **194**, 579 (1987)

| | | |
|-------------|----------------------|------|
| Ξ^0 | $1/2^+$ | **** |
| Ξ^- | $1/2^+$ | **** |
| $\Xi(1530)$ | $3/2^+$ | **** |
| $\Xi(1620)$ | | * |
| $\Xi(1690)$ | | *** |
| $\Xi(1820)$ | $3/2^-$ | *** |
| $\Xi(1950)$ | | *** |
| $\Xi(2030)$ | $\geq \frac{5}{2}^?$ | *** |
| $\Xi(2120)$ | | * |
| $\Xi(2250)$ | | ** |
| $\Xi(2370)$ | | ** |
| $\Xi(2500)$ | | * |

| | | |
|------------------|---------|------|
| Ω^- | $3/2^+$ | **** |
| $\Omega(2250)^-$ | | *** |
| $\Omega(2380)^-$ | | ** |
| $\Omega(2470)^-$ | | ** |

$\Omega(2250)^-$ $I(J^P) = 0(?)^?$ Status: ***

$\Omega(2250)^-$ MASS

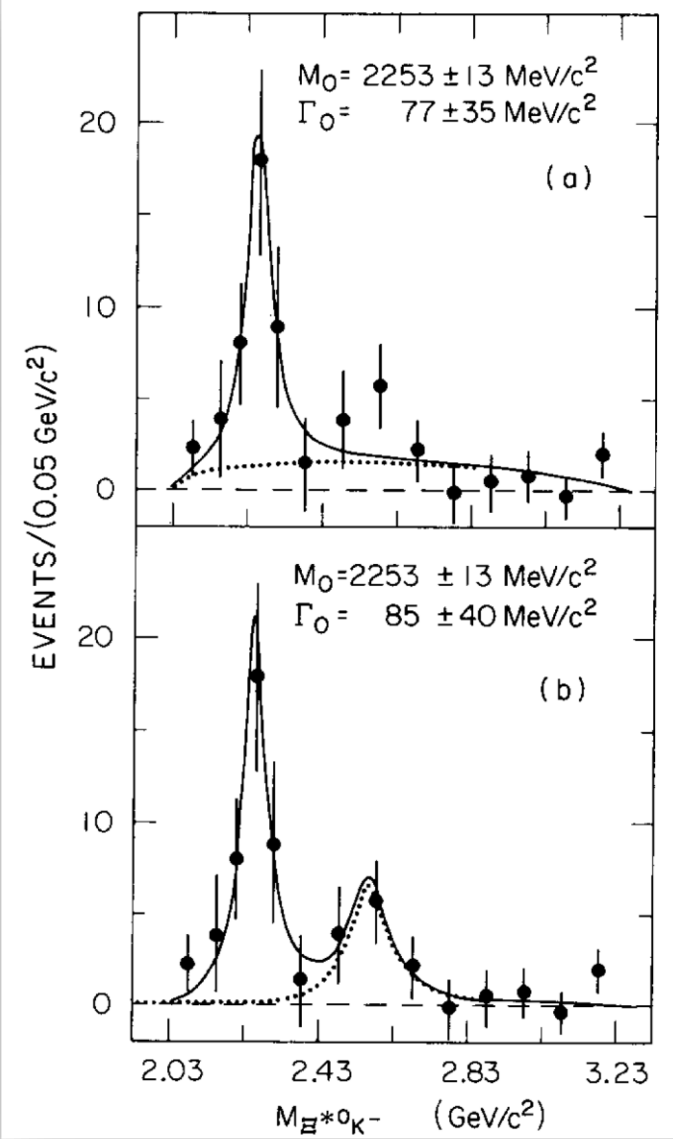
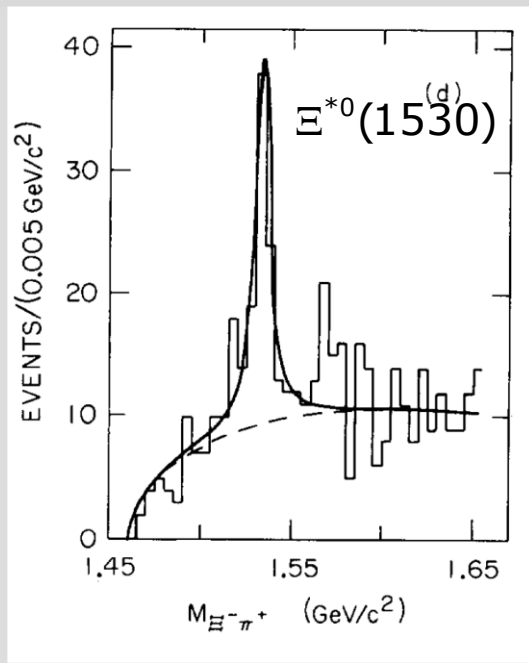
| VALUE (MeV) | EVTS | DOCUMENT ID | TECN | COMMENT |
|--------------------------------------------|------|-------------|----------|------------------|
| 2252 ± 9 OUR AVERAGE | | | | |
| 2253 \pm 13 | 44 | ASTON | 87B LASS | $K^- p$ 11 GeV/c |
| 2251 \pm 9 \pm 8 | 78 | BIAGI | 86B SPEC | SPS Ξ^- beam |

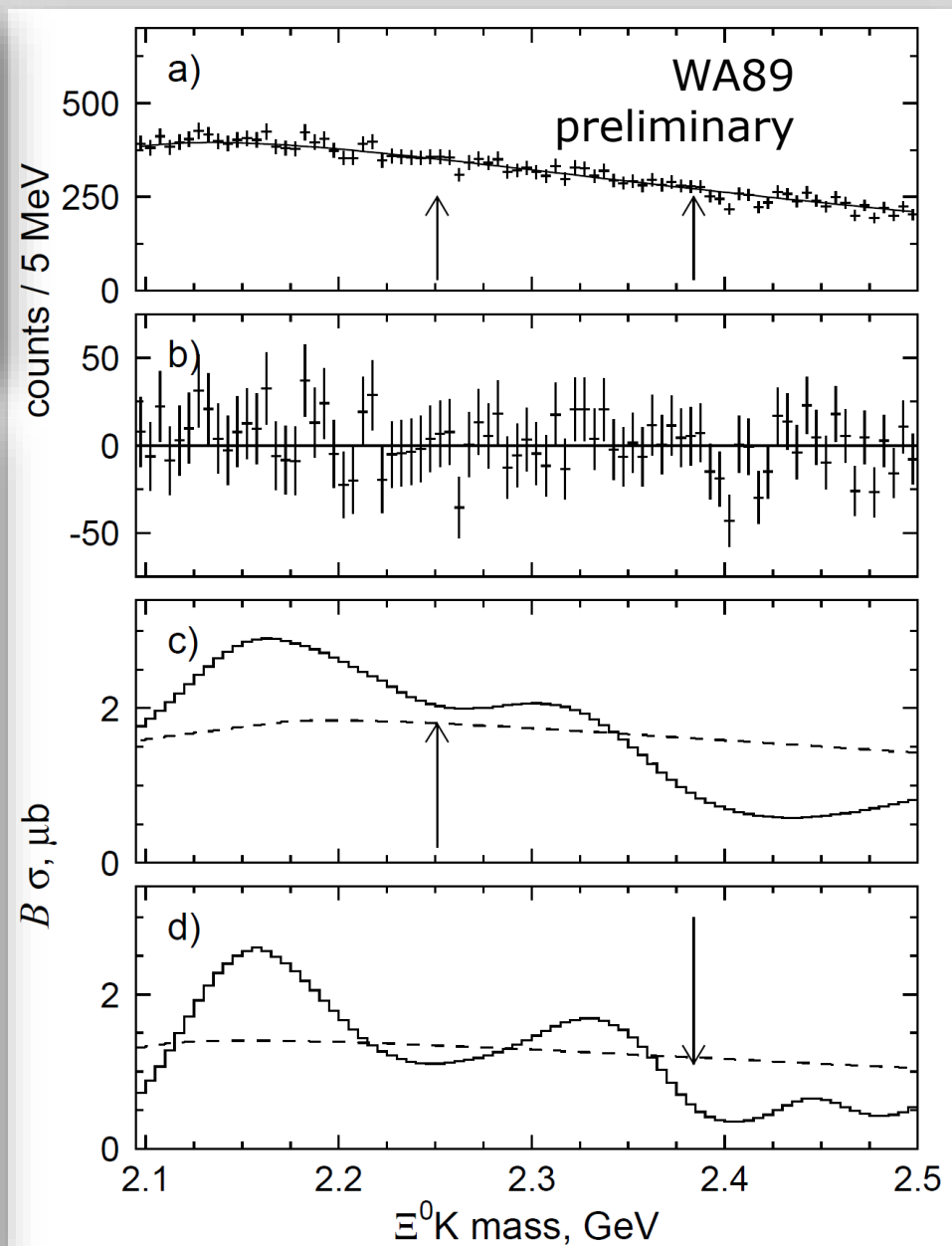
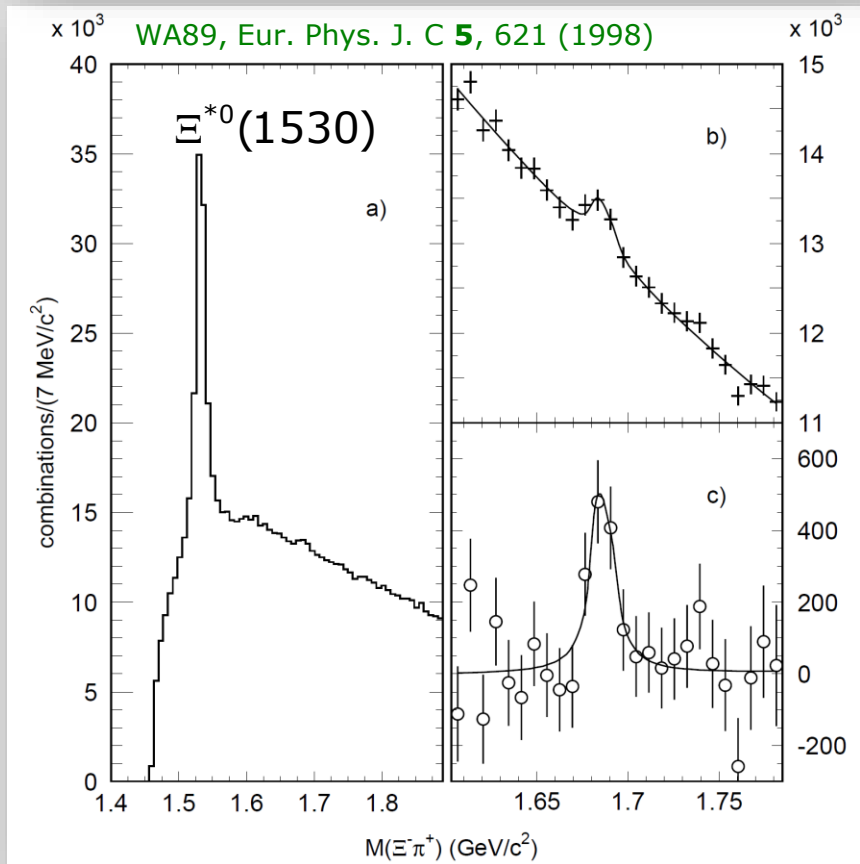
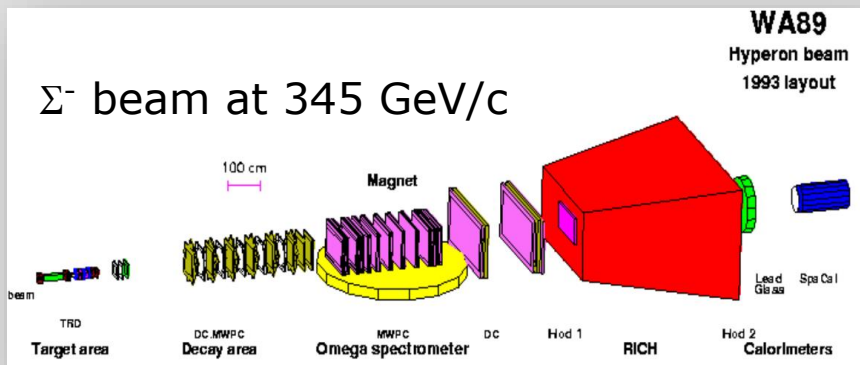
$\Omega(2380)^-$ Status: **

OMITTED FROM SUMMARY TABLE

$\Omega(2380)^-$ MASS

| VALUE (MeV) | EVTS | DOCUMENT ID | TECN | COMMENT |
|-----------------------------------------------|------|-------------|----------|------------------|
| ≈ 2380 OUR ESTIMATE | | | | |
| 2384 \pm 9 \pm 8 | 45 | BIAGI | 86B SPEC | SPS Ξ^- beam |





Spin Effect

WA89

J. Pochodzalla



$$\Sigma^- = |d\uparrow d\uparrow s\downarrow\rangle$$

octet

$$\Sigma^- = |d\uparrow d\uparrow s\downarrow\rangle$$

$$\Sigma^+ = |u\uparrow u\uparrow s\downarrow\rangle$$

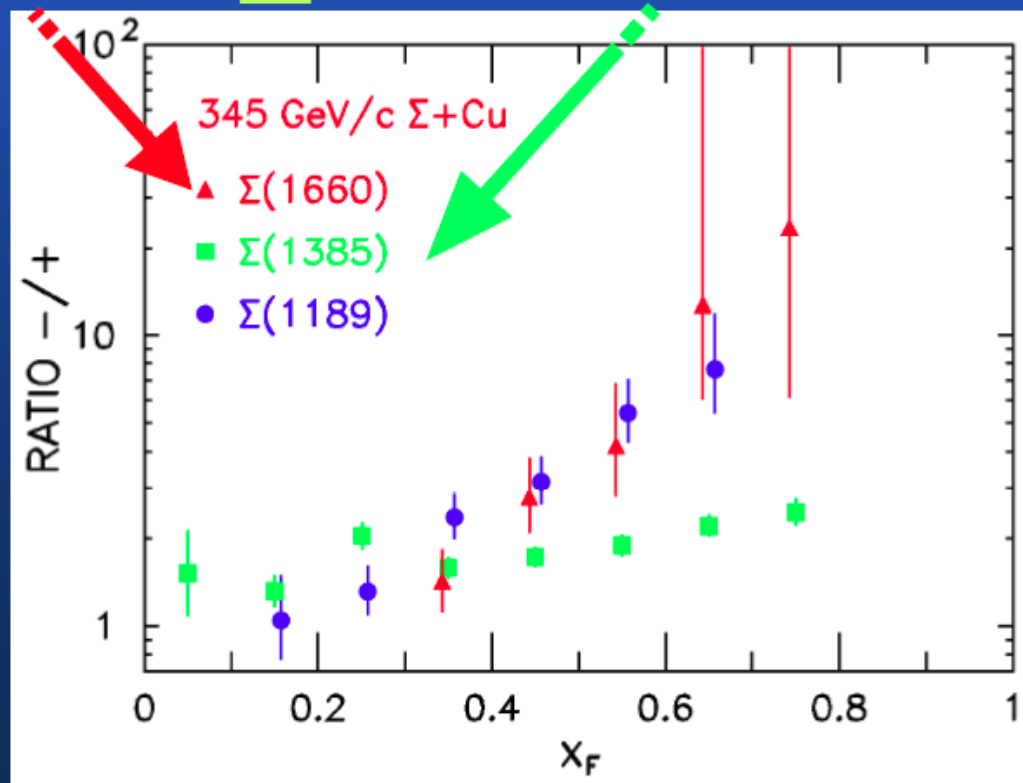
decuplet

$$\Sigma^- = |d\uparrow d\uparrow s\uparrow\rangle$$

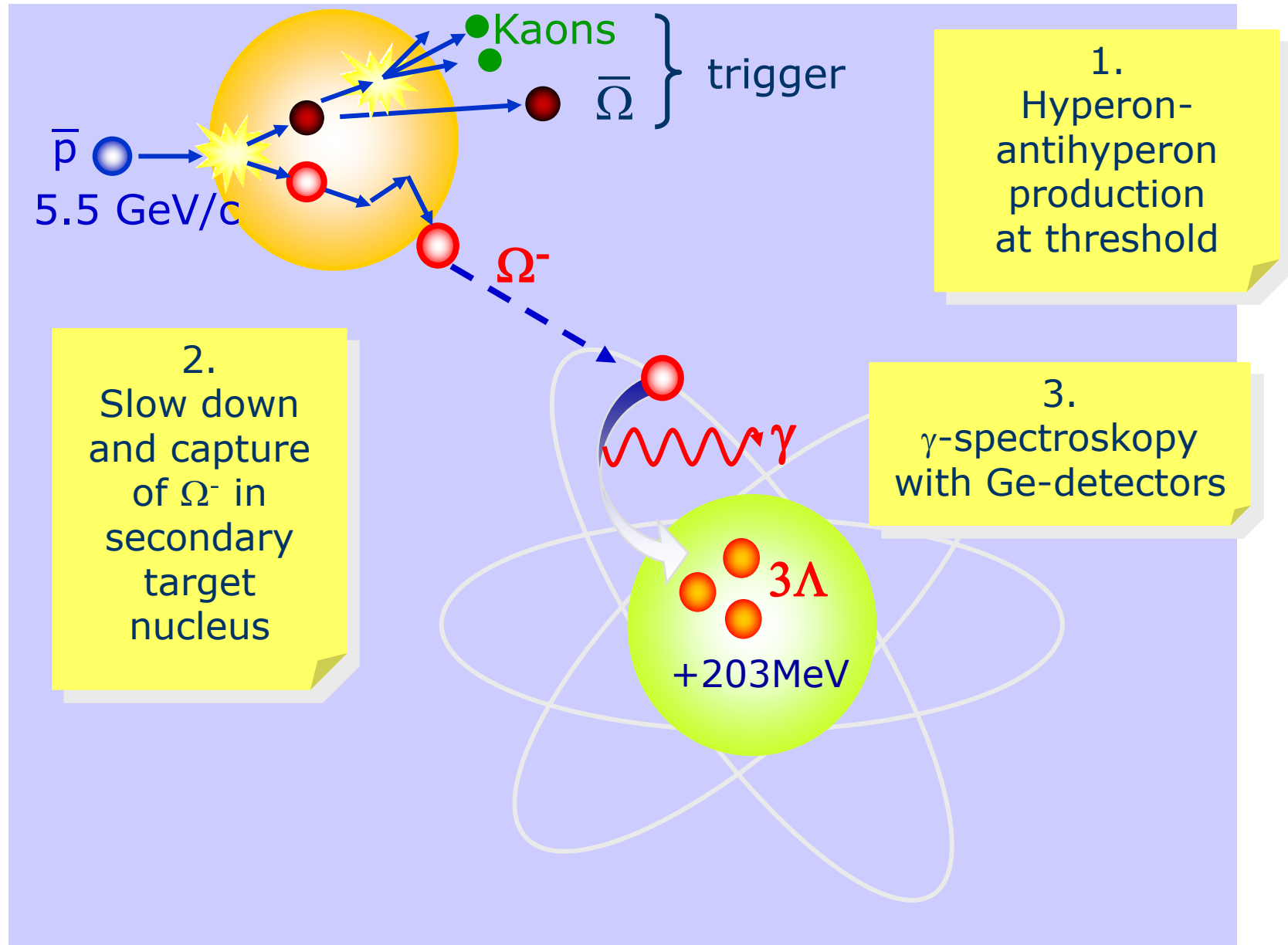
$$\Sigma^+ = |u\uparrow u\uparrow s\uparrow\rangle$$

Leading effect...

- depends on diquark (spin) structure of projectile



Production of Ω -Atoms



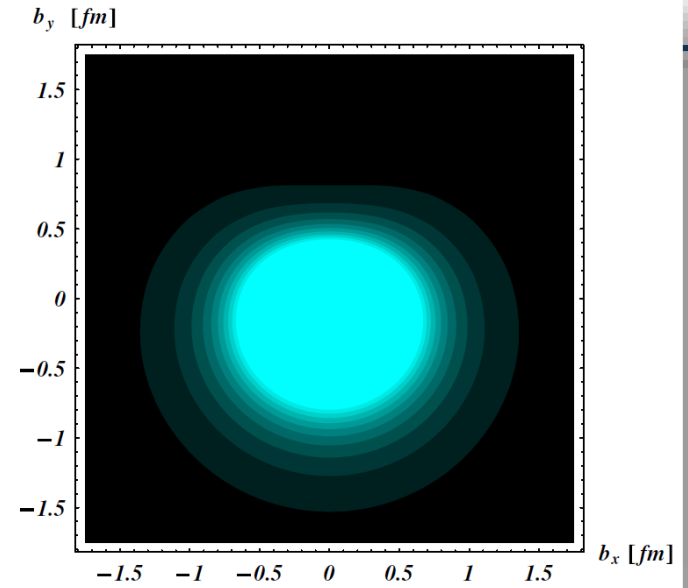
Proton vs. Omega

PHYSICAL REVIEW D 83, 054011 (2011)

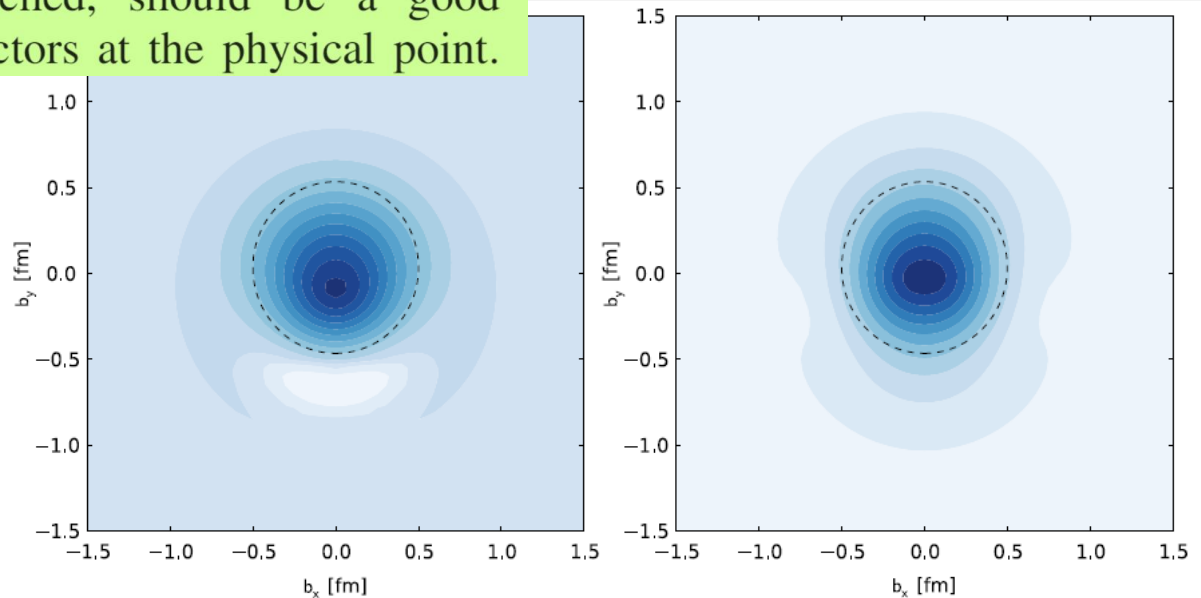
Extracting the Ω^- electric quadrupole moment from lattice QCD data

G. Ramalho¹ and M. T. Peña^{1,2}

Another important issue is that in sea quark effects for the Ω^- only at most one single light quark participates, and therefore the pion has no role in this case. As in chiral perturbation theory loops involving mesons heavier than the pion are suppressed, the Ω^- becomes then a special case where meson cloud corrections to the valence quark core are expected to be small. A consequence of the smallness of the meson cloud effects is that lattice QCD simulations, quenched or unquenched, should be a good approximation to Ω^- form factors at the physical point.



the x axis. Left: $\rho_{T3/2}^{\Omega}(\vec{b})$. Right: $\rho_{T1/2}^{\Omega}(\vec{b})$. A
evaluation of the densities we used the dipole



- ▶ $J=1/2$ baryons have no *spectroscopic* quadrupole moment

$$Q_i = \int d^3r \rho(r) (3z^2 - r^2)$$

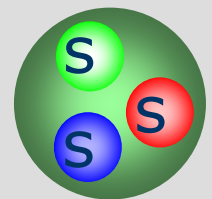
$$Q_s \propto (3J_z^2 - J(J+1)) \xrightarrow{J=1/2, J_z=1/2} 0$$

- ▶ The Ω^- Baryon is the only „elementary“ particle whose quadrupole moment can be measured
 - ▶ $J=3/2$
 - ▶ long mean lifetime $0.82 \cdot 10^{-10}$ s

- ▶ Contributions to *intrinsic* quadrupole moment of baryons
 - ▶ General: One-gluon exchange and meson exchange
 - ▶ Ω : only one-gluon contributions to quadrupole moment

A.J. Buchmann Z. Naturforsch. **52** (1997) 877-940

- ▶ sensitive to SU(3) symmetry e.g. within SU(3) limit $m_u/m_s=1$



$$Q_\Omega = Q_\Delta(\text{gluon})$$

A very strange Atom

- ▶ hyperfine splitting in Ω -atom
 \Rightarrow electric quadrupole moment of Ω

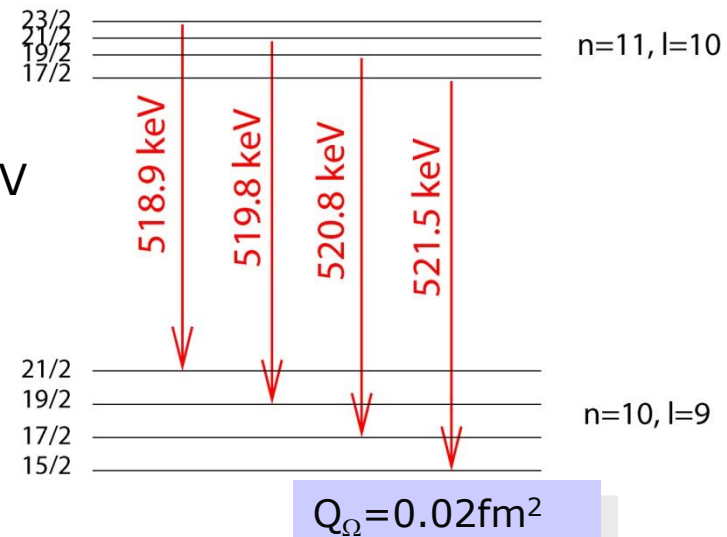
spin-orbit $\Delta E_{/s} \sim (aZ)^4 l \cdot m_{\Omega}$

quadrupole $\Delta E_{\Theta} \sim (aZ)^4 Q m^3_{\Omega}$

R.M. Sternheimer, M. Goldhaber, Phys. Rev. A 8, 2207 (1973)

M.M. Giannini, M.I. Krivoruchenko, Phys. Lett. B 291, 329 (1992)

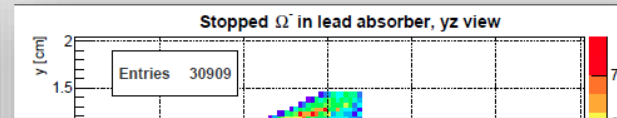
- ▶ prediction $Q_{\Omega} = (0 - 3.1) 10^{-2} \text{ fm}^2$
 - ▶ $E(n=11, l=10 \rightarrow n=10, l=9) \sim 520 \text{ keV}$
 - ▷ calibration with 511keV line!
 - ▶ $\Delta E_{\Theta} \sim \text{few tenth of keV for Pb}$





Ω^- Quadrupole Moment

| Model | Q [fm ²] | Reference |
|------------------|----------------------|-----------------------------------------------------------------------------------|
| NRQM | 0.018 | S.S. Gershtein, Yu.M., Zinoviev Sov. J. Nucl. Phys. 33, 772 (1981) |
| NRQM | 0.004 | J.-M. Richard, Z. Phys. C 12, 369 (1982) |
| NRQM | 0.031 | N. Isgur, G. Karl, R. Koniuk, Phys. Rev. D 25, 2395 (1982) |
| SU(3) Bag model | 0.052 | M.I. Krivoruchenko, Sov. J. Nucl. Phys. 45, 109 (1987) |
| QCD-SR | 0.1 | K. Azizi, Eur. Phys. J C 61, 311 (2009); T.M. Aliev, et al., arxiv: 0904.2485 |
| NRQM with mesons | 0.0057 | W.J. Leonard, W.J. Gerace, Phys. Rev. D 41, 924 (1990) |
| NQM | 0.028 | M.I. Krivoruchenko, M.M. Giannini, Phys. Rev. D 43, 3763 (1991) |
| Lattice QCD | 0.005 | D.B. Leinweber, T. Draper, R.M. Woloshyn, Phys. Rev. D 46, 3067 (1992) |
| HB χ PT | 0.009 | M.N. Butler, M.J. Savage, R.P. Springer, Phys. Rev. D 49, 3459 (1994) |
| Skyrme | 0.024 | J. Kroll, B. Schwesinger, Phys. Lett. B 334, 287 (1994) |
| Skyrme | 0.0 | Yoongseok Oh, ep-ph/9506308 |
| QM | 0.022 | A.J. Buchmann, Z. Naturforschung 52a, 877 (1997) |
| χ QM | 0.026 | G. Wagner, A.J. Buchmann, A. Faessler, J. Phys. G 26, 267 (2000) |
| GP QCD | 0.024 | A.J. Buchmann, E.M. Henley, Phys. Rev. D 65, 073017 (2002) |
| χ PT+q QCD | 0.0086 | L.S. Geng, J. Martin Camalich, M.J. Vicente Vacas, Phys. Rev. D 80, 034027 (2009) |
| Lattice QCD | 0.0096 \pm 0.0002 | G. Ramalho, M.T. Pena, Phys.Rev.D83:054011 (2011), arxiv:1012.2168 |



For Ξ^- atoms, low luminosity and Fe absorber:

Single X-ray lines $(6,5) \rightarrow (5,4)$: $\sim 3400/\text{month}$

Cascade events $(7,6) \rightarrow (6,5) \wedge (6,5) \rightarrow (5,4)$ $\sim 100/\text{month}$

for Ta target $\sim 25\%$ less

\Rightarrow ideal for commissioning phase of hypernucleus setup

20mm

x [cm]

Estimate of $\Omega\bar{\Omega}$ compared to $\Xi\bar{\Xi}$ (full luminosity)

Production yield: $\times 1/100$ (A.B. Kaidalov, Volkovitsky, Z. Phys. C63, 517 (1994))

Stopping probability $\times 1/100$

Ω^- atoms $\sim 5/\text{day}$

Single X-rays $\sim 4/\text{month}$

\Rightarrow For the first time this textbook experiment is within reach

but: stil many open questions: trigger, background,...

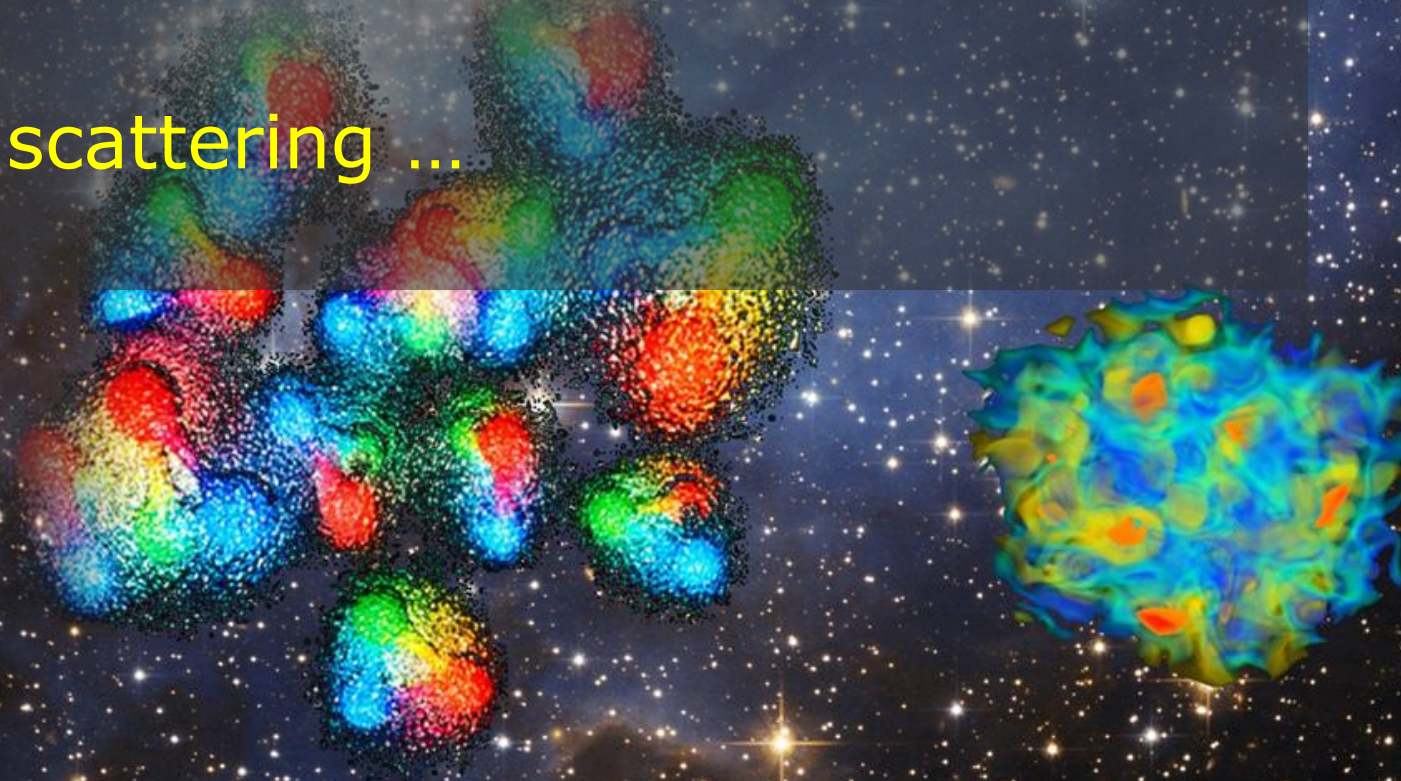
energy [mev]

Many more things...

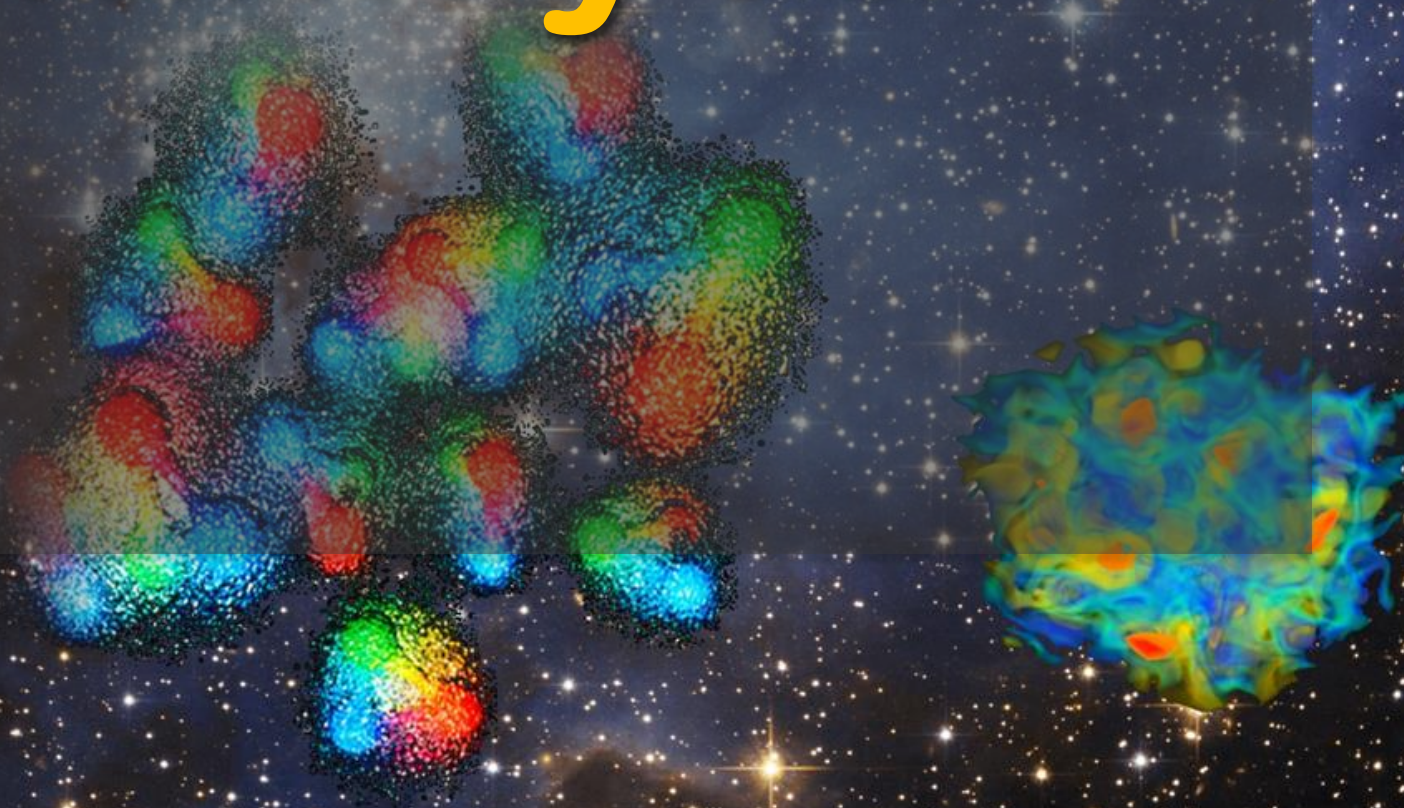
...kaonic atoms...

...antihyperons in nuclei...

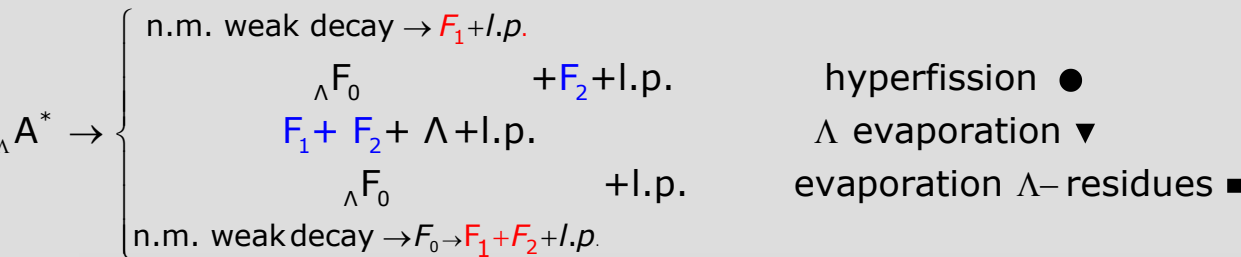
...hyperon scattering ...



Thank you

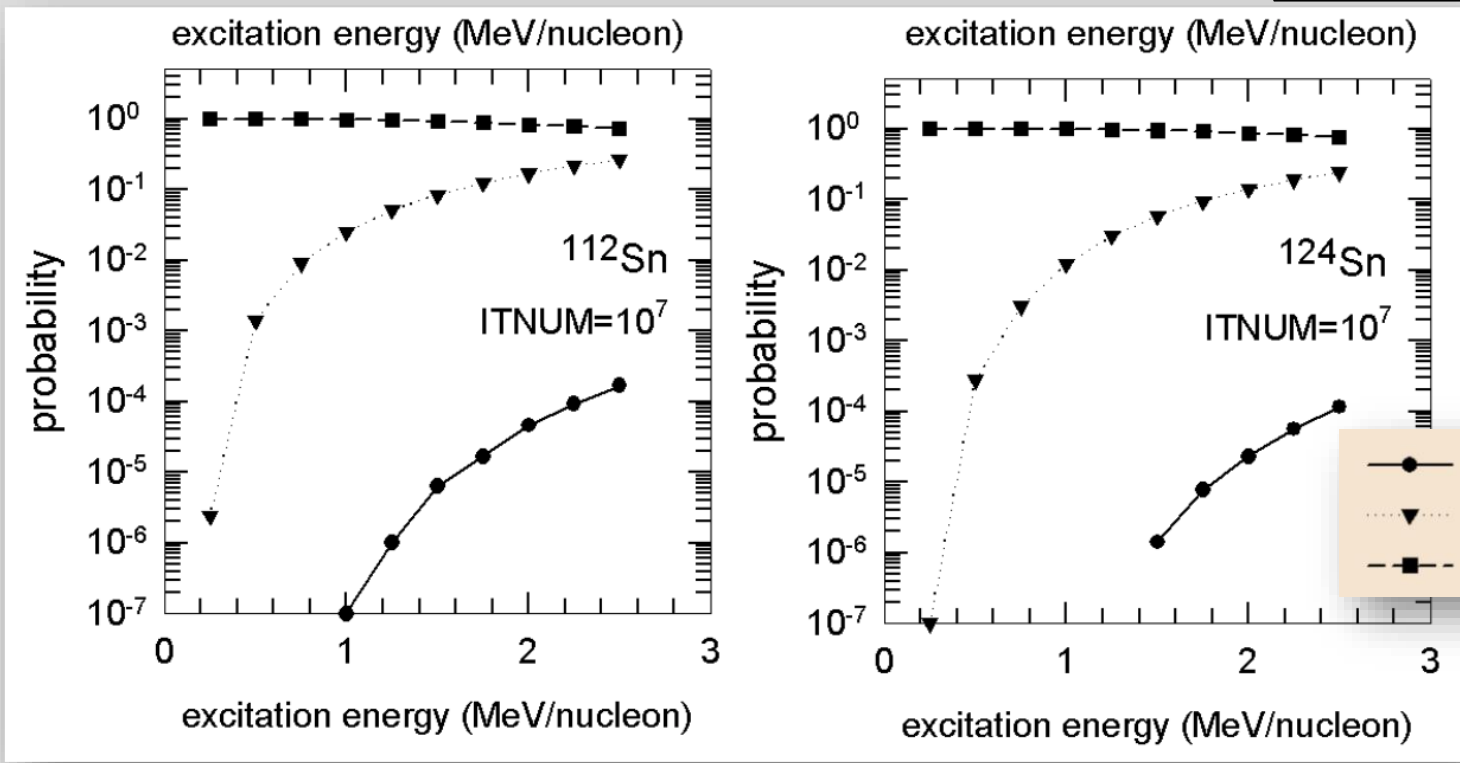


- ▶ expected excitation energy of initial excited hyperfragment $\sim 50\text{-}100\text{MeV}$ i.e. $\sim 1\text{MeV/nucleon}$



Tang: „delayed“
 \Rightarrow 1 prompt + 1 delayed fragment

\Rightarrow two prompt fragments
 \Rightarrow two delayed fragments
Tang: „prompt“



- ▶ use different isotopes as target